# ALMOST MINIMIZERS FOR THE THIN OBSTACLE PROBLEMS 

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To my family

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#### Abstract

In this dissertation, we consider almost minimizers for the thin obstacle problems in different settings: Laplacian, fractional Laplacian and equation with variable coefficients.

In Chapter 1, we consider Anzellotti-type almost minimizers for the thin obstacle (or Signorini) problem with zero thin obstacle and establish their $C^{1, \beta}$ regularity on the either side of the thin manifold, the optimal growth away from the free boundary, the $C^{1, \gamma}$ regularity of the regular part of the free boundary, as well as a structural theorem for the singular set. The analysis of the free boundary is based on a successful adaptation of energy methods such as a one-parameter family of Weiss-type monotonicity formulas, Almgren-type frequency formula, and the epiperimetric and logarithmic epiperimetric inequalities for the solutions of the thin obstacle problem. This chapter is based on a joint work with Arshak Petrosyan [1].

In Chapter 2, we study almost minimizers for the thin obstacle problem with variable Hölder continuous coefficients and zero thin obstacle and establish their $C^{1, \beta}$ regularity on the either side of the thin space. Under an additional assumption of quasisymmetry, we establish the optimal growth of almost minimizers as well as the regularity of the regular set and a structural theorem on the singular set. The proofs are based on the generalization of Weiss- and Almgren-type monotonicity formulas for almost minimizers established earlier in the case of constant coefficients (Chapter 1). This chapter is based on recent joint work with Arshak Petrosyan and Mariana Smit Vega Garcia [2].

In Chapter 3, we introduce a notion of almost minimizers for certain variational problems governed by the fractional Laplacian, with the help of the Caffarelli-Silvestre extension. In particular, we study almost fractional harmonic functions and almost minimizers for the fractional obstacle problem with zero obstacle. We show that for a certain range of parameters, almost minimizers are almost Lipschitz or $C^{1, \beta}$-regular. This is based on a work in collaboration with Arshak Petrosyan [3].


## 1. ALMOST MINIMIZERS FOR THE THIN OBSTACLE PROBLEM

### 1.1 Introduction and main results

### 1.1.1 The thin obstacle (or Signorini) problem

Let $D \subset \mathbb{R}^{n}$ be an open set and $\mathcal{M} \subset \mathbb{R}^{n}$ a smooth ( $n-1$ )-dimensional manifold (the thin space) and consider the problem of minimizing the Dirichlet energy

$$
\begin{equation*}
J_{D}(u):=\int_{D}|\nabla u(x)|^{2} d x \tag{1.1.1}
\end{equation*}
$$

among all functions $u \in W^{1,2}(D)$ satisfying

$$
u=g \text { on } \partial D, \quad u \geq \psi \text { on } \mathcal{M} \cap D,
$$

where $\psi: \mathcal{M} \rightarrow \mathbb{R}$ is the so-called thin obstacle and $g: \partial D \rightarrow \mathbb{R}$ is the prescribed boundary data with $g \geq \psi$ on $\mathcal{M} \cap \partial D$. This problem is known as the thin obstacle problem. In other words, it is a constrained minimization problem for the energy functional $J_{D}$ on a closed convex set

$$
\mathfrak{K}_{\psi, g}(D, \mathcal{M}):=\left\{u \in W^{1,2}(D): u=g \text { on } \partial D, u \geq \psi \text { on } \mathcal{M} \cap D\right\}
$$

This problem can be viewed as a scalar version of the Signorini problem with unilateral constraint from elastostatics [4] and is often referred to as the Signorini problem. It goes back to the origins of variational inequalities and is considered as one of the prototypical examples of such problems, see [5]. An equivalent formulation is given in the form

$$
\begin{aligned}
\Delta u=0 & \text { on } D \backslash \mathcal{M}, \\
u=g & \text { on } \partial D \\
u \geq \psi, \quad \partial_{\nu^{+}} u+\partial_{\nu^{-}} u \geq 0, \quad\left(\partial_{\nu^{+}} u+\partial_{\nu^{-}} u\right)(u-\psi)=0 & \text { on } \mathcal{M} \cap D,
\end{aligned}
$$

where the conditions on $\mathcal{M} \cap D$ are known as the Signorini complementarity (or ambiguous) conditions. Here, $\partial_{\nu^{ \pm}}$are the exterior normal derivatives from the either side of $\mathcal{M}$. In particular, at points on $\mathcal{M} \cap D$ we must have one of the two boundary conditions satisfied: either $u=\psi$ or $\partial_{\nu^{+}} u+\partial_{\nu^{-}} u=0$. The set

$$
\begin{equation*}
\Gamma(u):=\partial_{\mathcal{M}}\{x \in \mathcal{M} \cap D: u(x)=\psi(x)\}, \tag{1.1.2}
\end{equation*}
$$

which separates the regions where different boundary conditions are satisfied, is known as the free boundary and plays a central role in the analysis of the problem.

Because of the presence of the thin obstacle, it is not hard to realize that the solutions $u$ of the Signorini problem are at most Lipschitz across $\mathcal{M}$, even if both $\mathcal{M}$ and $\psi$ are smooth, as we may have $\partial_{\nu^{+}} u+\partial_{\nu^{-}} u>0$ at some points on $\mathcal{M}^{1}$. However, it has been known since the works [6]-[8] that the solutions of the thin obstacle problem are $C^{1, \beta}$ on $\mathcal{M}$ and consequently on the either side of $\mathcal{M}$, up to $\mathcal{M}$. In recent years, there has been a renewed interest in this problem, following the breakthrough result of Athanasopoulos and Caffarelli [9] on the optimal $C^{1,1 / 2}$ regularity of the minimizers (on the either side of $\mathcal{M}$ ) as well as its relation to the obstacle-type problems for the fractional Laplacian through the Caffarelli-Silvestre extension [10]. There has also been a significant effort in understanding the structure and the regularity of the free boundary. The results have been obtained in many settings, such as for the equations with variable coefficients, time-dependent versions, problems for fractional Laplacian and other nonlocal equations, both regarding the regularity of minimizers, as well as the properties of the free boundary; see e.g., [11]-[30] and many others.

### 1.1.2 Almost minimizers

In [31], Anzellotti introduced the notion of almost minimizers for energy functionals. Given $r_{0}>0$, we say that $\omega:\left(0, r_{0}\right) \rightarrow[0, \infty)$ is a modulus of continuity or a gauge function, if $\omega(r)$ is monotone nondecreasing in $r$ and $\omega(0+)=0$.

[^0]Definition 1.1.1 (Almost minimizers). Given $r_{0}>0$ and a gauge function $\omega(r)$ on ( $0, r_{0}$ ), we say that $u \in W_{\mathrm{loc}}^{1,2}(D)$ is an almost minimizer (or $\omega$-minimizer) for the functional $J_{D}$, if, for any ball $B_{r}\left(x_{0}\right) \Subset D$ with $0<r<r_{0}$, we have

$$
\begin{equation*}
J_{B_{r}\left(x_{0}\right)}(u) \leq(1+\omega(r)) J_{B_{r}\left(x_{0}\right)}(v) \quad \text { for any } v \in u+W_{0}^{1,2}\left(B_{r}\left(x_{0}\right)\right) \tag{1.1.3}
\end{equation*}
$$

The idea is that the Dirichlet energy of $u$ on the ball $B_{r}\left(x_{0}\right)$ is not necessarily minimal among all competitors $v \in u+W_{0}^{1,2}\left(B_{r}\left(x_{0}\right)\right)$ but almost minimal in the sense that it cannot decrease more than by a factor of $1+\omega(r)$. In the specific case of the energy functional $J_{D}$ in (1.1.1), i.e., the Dirichlet energy, we refer to the almost minimizers of $J_{D}$ as almost harmonic functions in $D$.

Results on almost minimizers for more general energy functionals can be found in [32][35]. Similar notions were considered earlier in the context of the geometric measure theory [36], [37], see also [38]. Almost minimizers are also related to quasiminimizers, introduced in [39], [40], see also [41]. For energy functionals exhibiting free boundaries, almost minimizers have been considered only recently in [42]-[46].

Almost minimizers can be viewed as perturbations of minimizers of various nature, but their study is motivated also by the observation that the minimizers with certain constrains, such as the ones with fixed volume or solutions of the obstacle problem, are realized as almost minimizers of unconstrained problems, see e.g. [31]. Yet another motivation is that the study of almost minimizers reveals a unique perspective on the problem and leads to the development of methods relying on less technical assumptions, thus allowing further generalization.

In this chapter we extend the notion of almost minimizers to the thin obstacle problem. Essentially, in (1.1.3), we restrict the function $u$ and its competitors $v$ to stay above the thin obstacle $\psi$ on $\mathcal{M}$.

Definition 1.1.2 (Almost minimizer for the thin obstacle (or Signorini) problem). Given $r_{0}>0$ and a gauge function $\omega(r)$ on $\left(0, r_{0}\right)$, we say that $u \in W_{\mathrm{loc}}^{1,2}(D)$ is an almost minimizer
for the thin obstacle (or Signorini) problem, if $u \geq \psi$ on $\mathcal{M} \cap D$ and, for any ball $B_{r}\left(x_{0}\right) \Subset D$ with $0<r<r_{0}$, we have

$$
\begin{equation*}
J_{B_{r}\left(x_{0}\right)}(u) \leq(1+\omega(r)) J_{B_{r}\left(x_{0}\right)}(v), \quad \text { for any } v \in \mathfrak{K}_{\psi, u}\left(B_{r}\left(x_{0}\right), \mathcal{M}\right) \tag{1.1.4}
\end{equation*}
$$

Note that in the case when $\mathcal{M} \cap B_{r}\left(x_{0}\right)=\emptyset$, the condition (1.1.4) is the same as (1.1.3) and thus almost minimizers of the Signorini problem are almost harmonic in $D \backslash \mathcal{M}$. As in the case of the solutions of the Signorini problem, we are interested in the regularity properties of almost minimizers as well as the structure and the regularity of the free boundary $\Gamma(u) \subset \mathcal{M}$ as defined in (1.1.2).

Some examples of almost minimizers are given in Appendix 1.A. We would also like to mention here that a related notion of almost minimizers for the fractional obstacle problem has been considered by the authors in [3].

### 1.1.3 Main results

Because of the technical nature of the problem, in this chapter we restrict ourselves only to the case when $\omega(r)=r^{\alpha}$ for some $\alpha>0, \mathcal{M}$ is flat, specifically $\mathcal{M}=\mathbb{R}^{n-1} \times\{0\}$, and the thin obstacle $\psi=0$. As we are mainly interested in local properties of almost minimizers and their free boundaries, we assume that $D$ is the unit ball $B_{1}, u \in W^{1,2}\left(B_{1}\right)$, and the constant $r_{0}=1$ in Definition 1.1.2. We also assume that $u$ is even in $x_{n}$-variable:

$$
u\left(x^{\prime}, x_{n}\right)=u\left(x^{\prime},-x_{n}\right), \quad \text { for any } x=\left(x^{\prime}, x_{n}\right) \in B_{1} .
$$

Our first main result is then as follows.
Theorem A ( $C^{1, \beta}$-regularity of almost minimizers). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$, under the assumptions above. Then, $u \in C_{\mathrm{loc}}^{1, \beta}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)$ for $\beta=\beta(\alpha, n)$ and

$$
\|u\|_{C^{1, \beta}(K)} \leq C\|u\|_{W^{1,2}\left(B_{1}\right)},
$$

for any $K \Subset B_{1}^{ \pm} \cup B_{1}^{\prime}$ and $C=C(n, \alpha, K)$.

The proof is obtained by using Morrey and Campanato space estimates, following the original idea of Anzellotti [31]. However, in our case the proof is much more elaborate and, in a sense, based on the idea that the solutions of the Signorini problem are 2-valued harmonic functions, as we have to work with both even and odd extensions of $u$ and $\nabla u$ from $B_{1}^{+}$to $B_{1}$.

While the optimal regularity for the minimizer (or solutions) of the Signorini problem is $C^{1,1 / 2}$, we do not expect such regularity for almost minimizers. However, we are able to establish the optimal growth for almost minimizers, which then allows to study the local properties of the free boundary

$$
\Gamma(u)=\partial\{u(\cdot, 0)=0\} \cap B_{1}^{\prime} .
$$

Theorem B (Optimal growth near free boundary). Let $u$ be as in Theorem A. Then,

$$
\int_{\partial B_{r}\left(x_{0}\right)} u^{2} \leq C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+2}
$$

for $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(u), 0<r<r_{0}(n, \alpha)$.
One of the ingredients in the proof is an Almgren-type monotonicity formula, which we describe below. For an almost minimizer $u$, Almgren's frequency [47] is defined by

$$
N\left(r, u, x_{0}\right):=\frac{r \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2}}{\int_{\partial B_{r}\left(x_{0}\right)} u^{2}}, \quad x_{0} \in \Gamma(u) .
$$

It is one of the most important monotone quantities in the analysis of the free boundary for the Signorini problem, see e.g. Chapter 9 in [48]. We show that for almost minimizers a small modification of $N$ is monotone.

Theorem C (Monotonicity of the truncated frequency). Let $u$ be as in Theorem A. Then for any $\kappa_{0} \geq 2$, there is $b=b\left(n, \alpha, \kappa_{0}\right)$ such that

$$
r \mapsto \widehat{N}_{\kappa_{0}}\left(r, u, x_{0}\right):=\min \left\{\frac{1}{1-b r^{\alpha}} N\left(r, u, x_{0}\right), \kappa_{0}\right\}
$$

is monotone for $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(u)$, and $0<r<r_{0}\left(n, \alpha, \kappa_{0}\right)$. Moreover, we have that either

$$
\widehat{N}_{\kappa_{0}}\left(0+, u, x_{0}\right)=3 / 2 \quad \text { or } \quad \widehat{N}_{\kappa_{0}}\left(0+, u, x_{0}\right) \geq 2
$$

We give an indirect proof of this fact, based on an one-parametric family of Weiss-type energy functionals $\left\{W_{\kappa}\right\}_{0<\kappa<\kappa_{0}}$, see Theorem 1.5.1, that go back to the work [14] for the solutions of the Signorini problem and Weiss [49] for the classical obstacle problem. The fact that $\widehat{N} \geq 3 / 2$ at free boundary points is crucial for the proof of the optimal growth (Theorem B), however, the proof of Theorem B requires also an application of so-called epiperimetric inequality for Weiss's energy functional $W_{3 / 2}$ (see [20]), to remove a remaining logarithmic term.

Our next result concerns the subset of the free boundary

$$
\mathcal{R}(u):=\left\{x_{0} \in \Gamma(u): \widehat{N}\left(0+, u, x_{0}\right)=3 / 2\right\}
$$

where Almgren's frequency is minimal, known as the regular set of $u$.

Theorem D (Regularity of the regular set). Let $u$ be as in Theorem A. Then $\mathcal{R}(u)$ is $a$ relatively open subset of the free boundary $\Gamma(u)$ and is a $(n-2)$-dimensional manifold of class $C^{1, \gamma}$.

Our proof of this theorem is based on the use of the epiperimetric inequality and is similar to the one for the solutions of the Signorini problem in [20].

Finally, we state our main result for the so-called singular set. A free boundary point $x_{0} \in \Gamma(u)$ is called singular if the coincidence set $\Lambda(u):=\{u(\cdot, 0)=0\}$ has $H^{n-1}$-density zero at $x_{0}$, i.e.,

$$
\lim _{r \rightarrow 0+} \frac{H^{n-1}\left(\Lambda(u) \cap B_{r}^{\prime}\left(x_{0}\right)\right)}{H^{n-1}\left(B_{r}^{\prime}\right)}=0
$$

If $\widehat{N}_{\kappa_{0}}\left(0+, u, x_{0}\right)=\kappa<\kappa_{0}$, then $x_{0}$ is singular if and only if $\kappa=2 m, m \in \mathbb{N}$ (see Proposition 1.10.1). For such $\kappa$, we then define

$$
\Sigma_{\kappa}(u):=\left\{x_{0} \in \Gamma(u): \widehat{N}_{\kappa_{0}}\left(0+, u, x_{0}\right)=\kappa\right\} .
$$

Theorem E (Structure of the singular set). Let $u$ be as in Theorem A. Then, for any $\kappa=2 m<\kappa_{0}, m \in \mathbb{N}, \Sigma_{\kappa}(u)$ is contained in a countable union of $(n-2)$-dimensional manifolds of class $C^{1, l o g}$.

A more refined version of this theorem is given in Theorem 1.10.10. The proof is based on the logarithmic epiperimetric inequality of Colombo-Spolaor-Velichkov [27] for Weiss's energy functional $W_{\kappa}$, with $\kappa=2 m, m \in \mathbb{N}$. We also point out that this inequality is instrumental in the proof of the optimal growth at singular points, which is rather immediate for the solutions of the Signorini problem, but far more complicated for the almost minimizers (see Lemmas 1.10.3-1.10.6).

## Proofs of Theorems A-E

While we don't give formal proofs of Theorems A-E, in the main body of the chapter, they follow from the combination of results there. More specifically,

- Theorem A follows by combining Theorems 1.3.1 and 1.4.1.
- The statement of Theorem B is contained in that of Lemma 1.7.4.
- Theorem C follows by combining Theorem 1.5.4 and Corollary 1.9.2.
- The statement of Theorem D is contained in that of Theorem 1.9.5.
- The statement of Theorem E is contained in that of Theorem 1.10.10.


### 1.1.4 Notation

Throughout the thesis we use the following notation. $\mathbb{R}^{n}$ stands for the $n$-dimensional Euclidean space. We denote the points of $\mathbb{R}^{n}$ by $x=\left(x^{\prime}, x_{n}\right)$, where $x^{\prime}=\left(x_{1}, \ldots, x_{n-1}\right) \in$ $\mathbb{R}^{n-1}$. We routinely identify $x^{\prime} \in \mathbb{R}^{n-1}$ with $\left(x^{\prime}, 0\right) \in \mathbb{R}^{n-1} \times\{0\}$. $\mathbb{R}_{ \pm}^{n}$ stand for open halfspaces $\left\{x \in \mathbb{R}^{n}: \pm x_{n}>0\right\}$.

For $x \in \mathbb{R}^{n}, r>0$, we use the following notations for balls of radius $r$, centered at $x$.

$$
\begin{aligned}
B_{r}(x) & =\left\{y \in \mathbb{R}^{n}:|x-y|<r\right\}, & & \text { ball in } \mathbb{R}^{n}, \\
B_{r}^{ \pm}\left(x^{\prime}\right) & =B_{r}\left(x^{\prime}, 0\right) \cap\left\{ \pm x_{n}>0\right\}, & & \text { half-ball in } \mathbb{R}^{n}, \\
B_{r}^{\prime}\left(x^{\prime}\right) & =B_{r}\left(x^{\prime}, 0\right) \cap\left\{x_{n}=0\right\}, & & \text { ball in } \mathbb{R}^{n-1}, \text { or thin ball. }
\end{aligned}
$$

We typically drop the center from the notation if it is the origin. Thus, $B_{r}=B_{r}(0)$, $B_{r}^{\prime}=B_{r}^{\prime}(0)$, etc.

Next, for a direction $\mathrm{e} \in \mathbb{R}^{n}$, we denote

$$
\partial_{\mathrm{e}} u=\nabla u \cdot \mathrm{e},
$$

the directional derivative of $u$ in the direction e. For the standard coordinate directions $\mathrm{e}_{i}$, $i=1, \ldots, n$, we simply write

$$
u_{x_{i}}=\partial_{x_{i}} u=\partial_{\mathrm{e}_{i}} u
$$

Moreover, by $\partial_{x_{n}}^{ \pm} u\left(x^{\prime}, 0\right)$ we mean the limit of $\partial_{x_{n}} u$ from within $B_{r}^{ \pm}$, specifically,

$$
\begin{aligned}
& \partial_{x_{n}}^{+} u\left(x^{\prime}, 0\right)=\lim _{\substack{y \rightarrow\left(x^{\prime}, 0\right) \\
y \in B_{r}^{+}}} \partial_{x_{n}} u(y)=-\partial_{\nu^{+}} u\left(x^{\prime}, 0\right), \\
& \partial_{x_{n}}^{-} u\left(x^{\prime}, 0\right)=\lim _{\substack{y \rightarrow\left(x^{\prime}, 0\right) \\
y \in B_{r}^{-}}} \partial_{x_{n}} u(y)=\partial_{\nu^{-}} u\left(x^{\prime}, 0\right),
\end{aligned}
$$

where $\nu^{ \pm}=\mp \mathrm{e}_{n}$ are unit outward normal vectors for $B_{r}^{ \pm}$on $B_{r}^{\prime}$.
In integrals, we often drop the variable and the measure of integration if it is with respect to the Lebesgue measure or the surface measure. Thus,

$$
\int_{B_{r}} u=\int_{B_{r}} u(x) d x, \quad \int_{\partial B_{r}} u=\int_{\partial B_{r}} u(x) d S_{x},
$$

where $S_{x}$ stands for the surface measure.
If $E$ is a set of positive and finite Lebesgue measure, we indicate by $\langle u\rangle_{E}$ the integral mean value of a function $u$ over $E$. That is,

$$
\langle u\rangle_{E}:=f_{E} u=\frac{1}{|E|} \int_{E} u .
$$

In particular, we indicate by $\langle u\rangle_{x, r}$ the integral mean value of a function $u$ over $B_{r}(x)$. That is,

$$
\langle u\rangle_{x, r}:=f_{B_{r}(x)} u=\frac{1}{\omega_{n} r^{n}} \int_{B_{r}(x)} u
$$

where $\omega_{n}=\left|B_{1}\right|$ is the volume of unit ball in $\mathbb{R}^{n}$. Similarly to the other notations, we drop the origin if it is 0 and write $\langle u\rangle_{r}$ for $\langle u\rangle_{0, r}$.

### 1.2 Almost harmonic functions

In this section we recall some results of Anzellotti [31] on almost harmonic functions, i.e., almost minimizers of the Dirichlet integral $J_{D}(v)=\int_{D}|\nabla v|^{2}$. We also state here some of the relevant auxiliary results that we will need also in the treatment of almost minimizers for the Signorini problem.

Theorem 1.2.1. Let $u$ be an almost harmonic function in an open set $D$ with a gauge function $\omega$. Then
(i) $u$ is locally almost Lipschitz, i.e., $u \in C_{\operatorname{loc}}^{0, \sigma}(D)$ for all $\sigma \in(0,1)$.
(ii) If $\omega(r) \leq C r^{\alpha}$ for some $\alpha \in(0,2)$, then $u \in C_{\mathrm{loc}}^{1, \alpha / 2}(D)$.

While we refer to [31] for the full proof of this theorem, we would like to outline the key steps in Anzellotti's argument. The idea to prove $C^{0, \sigma}$ and $C^{1, \alpha / 2}$ regularity of $u$ is through the Morrey and Campanato space estimates, namely, by establishing that

$$
\begin{align*}
& \int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} \leq C \rho^{n-2+2 \sigma}  \tag{1.2.1}\\
& \int_{B_{\rho}\left(x_{0}\right)}\left|\nabla u-\langle\nabla u\rangle_{x_{0}, \rho}\right|^{2} \leq C \rho^{n+\alpha} \tag{1.2.2}
\end{align*}
$$

for $x_{0} \in K \Subset D$, and $0<\rho<\rho_{0}$, with $C$ and $\rho_{0}$ depending on $n, r_{0}, d=\operatorname{dist}(K, \partial D)$, the gauge function $\omega$, and $\|u\|_{W^{1,2}(D)}$.

To obtain the estimates above, one starts by choosing a special competitor $v$ in (1.1.3). Namely, we take $v=h$ which solves the Dirichlet problem

$$
\Delta h=0 \quad \text { in } B_{r}\left(x_{0}\right), \quad h=u \quad \text { on } \partial B_{r}\left(x_{0}\right) .
$$

Equivalently, $h$ is the minimizer of the Dirichlet energy $\int_{B_{r}\left(x_{0}\right)}|\nabla v|^{2}$ among all functions in $u+W_{0}^{1,2}\left(B_{r}\left(x_{0}\right)\right)$. We call this $h$ the harmonic replacement of $u$ in $B_{r}\left(x_{0}\right)$. We then have the following concentric ball estimates for $h$.

Proposition 1.2.1. Let $h$ be harmonic in $B_{r}\left(x_{0}\right)$ and $0<\rho<r$. Then

$$
\begin{align*}
& \int_{B_{\rho}\left(x_{0}\right)}|\nabla h|^{2} \leq\left(\frac{\rho}{r}\right)^{n} \int_{B_{r}\left(x_{0}\right)}|\nabla h|^{2}  \tag{1.2.3}\\
& \int_{B_{\rho}\left(x_{0}\right)}\left|\nabla h-\langle\nabla h\rangle_{x_{0}, \rho}\right|^{2} \leq\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}\left(x_{0}\right)}\left|\nabla h-\langle\nabla h\rangle_{x_{0}, r}\right|^{2} . \tag{1.2.4}
\end{align*}
$$

Proof. The estimates above follow from the monotonicity in $\rho$ of the quantities

$$
\frac{1}{\rho^{n}} \int_{B_{\rho}\left(x_{0}\right)}|\nabla h|^{2}, \quad \frac{1}{\rho^{n+2}} \int_{B_{\rho}\left(x_{0}\right)}\left|\nabla h-\langle\nabla h\rangle_{x_{0}, \rho}\right|^{2} .
$$

Noticing that $\langle\nabla h\rangle_{x_{0}, \rho}=\nabla h\left(x_{0}\right)$, an easy proof is obtained by decomposing $h$ into the sum of the series of homogeneous harmonic polynomials.

We next use the almost minimizing property of $u$ to deduce perturbed versions of the estimates above.

Proposition 1.2.2. Let $u$ be an almost harmonic function in $D$. Then for any ball $B_{r}\left(x_{0}\right) \Subset$ $D$ with $r<r_{0}$ and $0<\rho<r$ we have

$$
\begin{align*}
& \int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} \leq 2\left[\left(\frac{\rho}{r}\right)^{n}+\omega(r)\right] \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2}  \tag{1.2.5}\\
& \int_{B_{\rho}\left(x_{0}\right)}\left|\nabla u-\langle\nabla u\rangle_{x_{0}, \rho}\right|^{2} \leq 9\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}\left(x_{0}\right)}\left|\nabla u-\langle\nabla u\rangle_{x_{0}, r}\right|^{2}  \tag{1.2.6}\\
& \quad+24 \omega(r) \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2} .
\end{align*}
$$

Proof. If $h$ is a harmonic replacement of $u$ in $B_{r}\left(x_{0}\right)$, we first note that

$$
\begin{aligned}
\int_{B_{r}\left(x_{0}\right)}|\nabla(u-h)|^{2} & =\int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2}-|\nabla h|^{2}-2 \int_{B_{r}\left(x_{0}\right)} \nabla h \nabla(u-h) \\
& =\int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2}-|\nabla h|^{2} \leq \omega(r) \int_{B_{r}\left(x_{0}\right)}|\nabla h|^{2} \leq \omega(r) \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2} .
\end{aligned}
$$

Then, combined with (1.2.3), we estimate

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} \leq 2 \int_{B_{\rho}\left(x_{0}\right)}|\nabla h|^{2}+2 \int_{B_{\rho}\left(x_{0}\right)}|\nabla(u-h)|^{2}
$$

$$
\leq 2\left[\left(\frac{\rho}{r}\right)^{n}+\omega(r)\right] \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2},
$$

which gives (1.2.5). To obtain (1.2.6), we argue very similarly by using additionally that by Jensen's inequality

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\langle\nabla u\rangle_{x_{0}, \rho}-\langle\nabla h\rangle_{x_{0}, \rho}\right|^{2} \leq \int_{B_{\rho}\left(x_{0}\right)}|\nabla u-\nabla h|^{2} .
$$

For more details we refer to the proof of Theorem 1.4.1, Case 1.1.

From here, one deduces the estimates (1.2.1)-(1.2.2) with the help of the following useful lemma. The proof can be found e.g. in [50].

Lemma 1.2.2. Let $r_{0}>0$ be a positive number and let $\varphi:\left(0, r_{0}\right) \rightarrow(0, \infty)$ be a nondecreasing function. Let $a$, $\beta$, and $\gamma$ be such that $a>0, \gamma>\beta>0$. There exist two positive numbers $\varepsilon=\varepsilon(a, \gamma, \beta), c=c(a, \gamma, \beta)$ such that, if

$$
\varphi(\rho) \leq a\left[\left(\frac{\rho}{r}\right)^{\gamma}+\varepsilon\right] \varphi(r)+b r^{\beta}
$$

for all $\rho$, $r$ with $0<\rho \leq r<r_{0}$, where $b \geq 0$, then one also has, still for $0<\rho<r<r_{0}$,

$$
\varphi(\rho) \leq c\left[\left(\frac{\rho}{r}\right)^{\beta} \varphi(r)+b \rho^{\beta}\right] .
$$

We can now give a formal proof of Theorem 1.2.1.

Proof of Theorem 1.2.1. (i) Taking $r_{0}$ small enough so that $\omega\left(r_{0}\right)<\varepsilon$, a direct application of Lemma 1.2.2 to (1.2.5) produces the estimate (1.2.1), which in turn implies that $u \in C_{\mathrm{loc}}^{0, \sigma}(D)$, by the Morrey space embedding theorem.
(ii) Using that $\omega(r) \leq C r^{\alpha}$, combined with the estimate (1.2.1), we first obtain

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\nabla u-\langle\nabla u\rangle_{x_{0}, \rho}\right|^{2} \leq 9\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}\left(x_{0}\right)}\left|\nabla u-\langle\nabla u\rangle_{x_{0}, r}\right|^{2}+C r^{n-2+2 \sigma+\alpha} .
$$

If $\sigma$ is so that $\alpha^{\prime}=-2+2 \sigma+\alpha>0$, Lemma 1.2.2 implies that

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\nabla u-\langle\nabla u\rangle_{x_{0}, \rho}\right|^{2} \leq C \rho^{n+\alpha^{\prime}}
$$

By the Campanato space embedding, we therefore obtain that $\nabla u \in C_{\mathrm{loc}}^{0, \alpha^{\prime} / 2}(D)$. However, it is easy to bootstrap the regularity up to $C_{\mathrm{loc}}^{0, \alpha / 2}$ by noticing that we now know that $\nabla u$ is locally bounded in $D$ and thus $\int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2} \leq C r^{n}$. Plugging that in the last term of (1.2.6), we obtain that

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\nabla u-\langle\nabla u\rangle_{x_{0}, \rho}\right|^{2} \leq 9\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}\left(x_{0}\right)}\left|\nabla u-\langle\nabla u\rangle_{x_{0}, r}\right|^{2}+C r^{n+\alpha}
$$

and repeating the arguments above conclude that $u \in C_{\mathrm{loc}}^{1, \alpha / 2}$.

### 1.3 Almost Lipschitz regularity of almost minimizers

In this section we prove the first regularity results for the almost minimizers for the Signorini problem, see Definition 1.1.2. Recall that we assume $D=B_{1}, \mathcal{M}=\mathbb{R}^{n-1} \times\{0\}$, $\psi=0, r_{0}=1$, and $\omega(r)=r^{\alpha}$ for some $\alpha>0$. Furthermore we assume that $u$ is even symmetric in $x_{n}$-variable.

Theorem 1.3.1. Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$. Then $u \in C^{0, \sigma}\left(B_{1}\right)$ for all $0<\sigma<1$. Moreover, for any $K \Subset B_{1}$,

$$
\begin{equation*}
\|u\|_{C^{0, \sigma}(K)} \leq C\|u\|_{W^{1,2}\left(B_{1}\right)} \tag{1.3.1}
\end{equation*}
$$

with $C=C(n, \alpha, \sigma, K)$.
The idea of the proof is to follow that of Anzellotti [31] that we outlined in Section 1.2 and to prove an estimate similar to (1.2.5). The proof of the latter estimate followed by a perturbation argument from a similar estimate for the harmonic replacement of $u$. However, in the case of the Signorini problem, the harmonic replacements are not necessarily admissible competitors. Instead, for $B_{r}\left(x_{0}\right) \Subset B_{1}$, we consider the Signorini replacements $h$ of $u$ in $B_{r}\left(x_{0}\right)$, which solve the Signorini problem in $B_{r}\left(x_{0}\right)$ with the thin obstacle 0 on $\mathcal{M}$
and boundary values $h=u$ on $\partial B_{r}\left(x_{0}\right)$. Equivalently, Signorini replacements are the minimizers of $J_{B_{r}\left(x_{0}\right)}$ on the constraint set $\mathfrak{K}_{0, u}\left(B_{r}\left(x_{0}\right), \mathcal{M}\right)$ and they also satisfy the variational inequality ${ }^{2}$

$$
\begin{equation*}
\int_{B_{r}\left(x_{0}\right)} \nabla h \cdot \nabla(v-h) \geq 0 \quad \text { for any } v \in \mathfrak{K}_{0, u}\left(B_{r}\left(x_{0}\right), \mathcal{M}\right) \tag{1.3.2}
\end{equation*}
$$

We then have the following concentric ball estimates for Signorini replacements similar to the one for harmonic replacements, at least when the center of the balls is on $\mathcal{M}=\mathbb{R}^{n-1} \times\{0\}$.

Proposition 1.3.1. Let $x_{0} \in \mathcal{M}$ and let $h$ be a solution of the Signorini problem in $B_{r}\left(x_{0}\right)$ with zero obstacle on $\mathcal{M}$, even in $x_{n}$-variable. Then,

$$
\begin{equation*}
\int_{B_{\rho}\left(x_{0}\right)}|\nabla h|^{2} \leq\left(\frac{\rho}{r}\right)^{n} \int_{B_{r}\left(x_{0}\right)}|\nabla h|^{2}, \quad 0<\rho<r . \tag{1.3.3}
\end{equation*}
$$

Proof. We claim that $|\nabla h|^{2}$ is subharmonic in $B_{r}\left(x_{0}\right)$. This follows from the fact that $h_{x_{\mathrm{i}}}^{ \pm}$, $\mathrm{i}=1, \ldots, n-1$, are subharmonic in $B_{r}\left(x_{0}\right)$, see [48], and similarly that the even extensions $\widetilde{h}_{x_{n}}^{ \pm}$of $h_{x_{n}}^{ \pm}$in $x_{n}$-variable from $B_{R}^{+}\left(x_{0}\right)$ to all of $B_{R}\left(x_{0}\right)$ are also subharmonic. These are all consequences of the fact that a continuous nonnegative function, subharmonic in its positivity set is subharmonic, see Ex. 2.6 in [48].

The subharmonicity of $|\nabla h|^{2}$ in $B_{r}\left(x_{0}\right)$ then implies, by the sub-mean value property, that the function

$$
\rho \mapsto \frac{1}{\rho^{n}} \int_{B_{\rho}\left(x_{0}\right)}|\nabla h|^{2}
$$

is monotone nondecreasing. This readily implies (1.3.3).

We next have the perturbed version of Proposition 1.3.1.

Proposition 1.3.2. Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$, and $B_{r}\left(x_{0}\right) \subset B_{1}$. Then, there is $C_{1}=C_{1}(n)>1$ such that

$$
\begin{equation*}
\int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} \leq C_{1}\left[\left(\frac{\rho}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2}, \quad 0<\rho<r . \tag{1.3.4}
\end{equation*}
$$

${ }^{2} \uparrow$ which follows from the inequality $\int_{B_{r}\left(x_{0}\right)}|\nabla h|^{2} \leq \int_{B_{r}\left(x_{0}\right)}|\nabla((1-\varepsilon) h+\varepsilon v)|^{2}, \varepsilon \in(0,1)$ by a first variation argument.

Proof. By using the continuity argument, we may assume that $B_{r}\left(x_{0}\right) \Subset B_{1}$. We first prove the estimate when $x_{0}$ is in the thin space, i.e., $x_{0} \in B_{1}^{\prime}$ and then extend it to arbitrary $x_{0} \in B_{1}$.

Case 1. Suppose $x_{0} \in B_{1}^{\prime}$ and let $h$ be the Signorini replacement of $u$ in $B_{r}\left(x_{0}\right)$. Recall that $h$ satisfies (1.3.2). Then, plugging $v=u$, we obtain

$$
\begin{equation*}
\int_{B_{r}\left(x_{0}\right)} \nabla h \cdot \nabla u-|\nabla h|^{2} \geq 0 \tag{1.3.5}
\end{equation*}
$$

Using this, we can estimate

$$
\begin{align*}
\int_{B_{r}\left(x_{0}\right)}|\nabla(u-h)|^{2} & =\int_{B_{r}\left(x_{0}\right)}\left(|\nabla u|^{2}+|\nabla h|^{2}-2 \nabla u \cdot \nabla h\right) \\
& \leq \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2}-\int_{B_{r}\left(x_{0}\right)}|\nabla h|^{2}  \tag{1.3.6}\\
& \leq\left(1+r^{\alpha}\right) \int_{B_{r}\left(x_{0}\right)}|\nabla h|^{2}-\int_{B_{r}\left(x_{0}\right)}|\nabla h|^{2} \\
& =r^{\alpha} \int_{B_{r}\left(x_{0}\right)}|\nabla h|^{2} \leq r^{\alpha} \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2},
\end{align*}
$$

where in the very last step we have used that $h$ minimizes the Dirichlet integral among all functions in $\mathfrak{K}_{0, u}\left(B_{r}\left(x_{0}\right), \mathcal{M}\right)$.

Next, we use the same perturbation argument as in the proof of (1.2.5). By using (1.3.3) and (1.3.6), we estimate

$$
\begin{aligned}
\int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} & \leq 2 \int_{B_{\rho}\left(x_{0}\right)}|\nabla h|^{2}+2 \int_{B_{\rho}\left(x_{0}\right)}|\nabla(u-h)|^{2} \\
& \leq 2\left(\frac{\rho}{r}\right)^{n} \int_{B_{r}\left(x_{0}\right)}|\nabla h|^{2}+2 r^{\alpha} \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2} \\
& \leq 2\left[\left(\frac{\rho}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2} .
\end{aligned}
$$

Thus, (1.3.4) follows in this case.
Case 2. Consider now the case $x_{0} \in B_{1}^{+}$. If $\rho \geq r / 4$, then we simply have

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} \leq 4^{n}\left(\frac{\rho}{r}\right)^{n} \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2} .
$$

Thus, we may assume $\rho<r / 4$. Then, let $d:=\operatorname{dist}\left(x_{0}, B_{1}^{\prime}\right)>0$ and choose $x_{1} \in \partial B_{d}\left(x_{0}\right) \cap B_{1}^{\prime}$. Case 2.1. If $\rho \geq d$, then we use $B_{\rho}\left(x_{0}\right) \subset B_{2 \rho}\left(x_{1}\right) \subset B_{r / 2}\left(x_{1}\right) \subset B_{r}\left(x_{0}\right)$ and the result of Case 1 to write

$$
\begin{aligned}
\int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} & \leq \int_{B_{2 \rho}\left(x_{1}\right)}|\nabla u|^{2} \leq C\left[\left(\frac{2 \rho}{r / 2}\right)^{n}+(r / 2)^{\alpha}\right] \int_{B_{r / 2}\left(x_{1}\right)}|\nabla u|^{2} \\
& \leq C\left[\left(\frac{\rho}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2} .
\end{aligned}
$$

Case 2.2. Suppose now $d>\rho$. If $d>r$, then $B_{r}\left(x_{0}\right) \Subset B_{1}^{+}$. Since $u$ is almost harmonic in $B_{1}^{+}$, we can apply Proposition 1.2.2 to obtain

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} \leq 2\left[\left(\frac{\rho}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2} .
$$

Thus, we may assume $d \leq r$. Then we note that $B_{d}\left(x_{0}\right) \subset B_{1}^{+}$and by a limiting argument from the previous estimate, we obtain

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} \leq 2\left[\left(\frac{\rho}{d}\right)^{n}+r^{\alpha}\right] \int_{B_{d}\left(x_{0}\right)}|\nabla u|^{2} .
$$

Case 2.2.1. If $r / 4 \leq d$, then

$$
\int_{B_{d}\left(x_{0}\right)}|\nabla u|^{2} \leq 4^{n}\left(\frac{d}{r}\right)^{n} \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2},
$$

which immediately implies (1.3.4).
Case 2.2.2. It remains to consider the case $\rho<d<r / 4$. Using Case 1 again, we have

$$
\begin{aligned}
\int_{B_{d}\left(x_{0}\right)}|\nabla u|^{2} & \leq \int_{B_{2 d}\left(x_{1}\right)}|\nabla u|^{2} \leq C\left[\left(\frac{2 d}{r / 2}\right)^{n}+(r / 2)^{\alpha}\right] \int_{B_{r / 2}\left(x_{1}\right)}|\nabla u|^{2} \\
& \leq C\left[\left(\frac{d}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2},
\end{aligned}
$$

which also implies (1.3.4). This concludes the proof of the proposition.
We can now give the proof of the almost Lipschitz regularity of almost minimizers.

Proof of Theorem 1.3.1. Let $K \Subset B_{1}$ and $x_{0} \in K$. Take $\delta=\delta(n, \alpha, \sigma, K)>0$ such that $\delta<\operatorname{dist}\left(K, \partial B_{1}\right)$ and $\delta^{\alpha} \leq \varepsilon\left(C_{1}, n, n+2 \sigma-2\right)$, where $\varepsilon=\varepsilon\left(C_{1}, n, n+2 \sigma-2\right)$ is as in Lemma 1.2.2. Then for all $0<\rho<r<\delta$, by (1.3.4),

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} \leq C_{1}\left[\left(\frac{\rho}{r}\right)^{n}+\varepsilon\right] \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2} .
$$

By applying Lemma 1.2.2, we obtain

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} \leq C(n, \sigma)\left(\frac{\rho}{r}\right)^{n+2 \sigma-2} \int_{B_{r}\left(x_{0}\right)}|\nabla u|^{2}
$$

Taking $r \nearrow \delta$, we can therefore conclude

$$
\begin{equation*}
\int_{B_{\rho}\left(x_{0}\right)}|\nabla u|^{2} \leq C(n, \alpha, \sigma, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} \rho^{n+2 \sigma-2} \tag{1.3.7}
\end{equation*}
$$

From here, we use the Morrey space embedding to obtain $u \in C^{0, \sigma}(K)$ with the norm estimate

$$
\|u\|_{C^{0, \sigma}(K)} \leq C(n, \alpha, \sigma, K)\|u\|_{W^{1,2}\left(B_{1}\right)}
$$

as required.

## $1.4 C^{1, \beta}$ regularity of almost minimizers

In this section we establish the $C^{1, \beta}$ regularity of almost minimizers for some $\beta>0$. The idea is again to use Signorini replacements and an appropriate version of the concentric ball estimate (1.2.4) for solutions of the Signorini problem.

As we saw in the proof of the almost Lipschitz regularity of almost minimizers, it is enough to obtain such estimates when balls are centered at $x_{0}$ on the thin space $\mathcal{M}=\mathbb{R}^{n-1} \times\{0\}$. It turns out that to prove a proper version of (1.2.4), we have to work with both even and odd extensions in $x_{n}$-variable of Signorini replacements $h$ from $B_{r}^{+}\left(x_{0}\right)$ to $B_{r}\left(x_{0}\right)$. The reason is that even extensions are harmonic across the positivity set $\{h(\cdot, 0)>0\}$, while the odd extensions are harmonic across the interior of the coincidence set $\{h(\cdot, 0)=0\}$.

Proposition 1.4.1. Let $h$ be a solution of the Signorini problem in $B_{r}\left(x_{0}\right)$ with $x_{0} \in \mathcal{M}$, even in $x_{n}$-variable. Define

$$
\widehat{\nabla h}:= \begin{cases}\nabla h\left(x^{\prime}, x_{n}\right), & x_{n} \geq 0 \\ \nabla h\left(x^{\prime},-x_{n}\right), & x_{n}<0\end{cases}
$$

the even extension of $\nabla h$ from $B_{r}^{+}\left(x_{0}\right)$ to $B_{r}\left(x_{0}\right)$. Then for $0<\alpha<1$, there is $C=C(n)$ such that for all $0<\rho \leq(3 / 4) r$,

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{x_{0}, \rho}\right|^{2} \leq C \frac{\rho^{n+1}}{r^{n+3}} \int_{B_{r}\left(x_{0}\right)} h^{2} .
$$

Proof. This is an immediate corollary of the estimate

$$
\begin{equation*}
\|\nabla h\|_{C^{0,1 / 2}\left(B_{(3 / 4) r}^{ \pm}\left(x_{0}\right) \cup B_{(3 / 4) r}\left(x_{0}\right)\right)} \leq C(n) r^{-\frac{n+3}{2}}\|h\|_{L^{2}\left(B_{r}^{+}\left(x_{0}\right)\right)} \tag{1.4.1}
\end{equation*}
$$

see e.g. Theorem 9.13 in [48]. Indeed, for $0<\rho \leq(3 / 4) r$, we have

$$
\begin{aligned}
\int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{x_{0}, \rho}\right|^{2} & =2 \int_{B_{\rho}^{+}\left(x_{0}\right)}\left|\nabla h-\langle\nabla h\rangle_{B_{r}^{+}\left(x_{0}\right)}\right|^{2} \\
& \leq C(n) \rho^{n+1}\|\nabla h\|_{C^{0,1 / 2}\left(B_{(3 / 4) r}^{+}\left(x_{0}\right)\right)}^{2} \\
& \leq C(n) \frac{\rho^{n+1}}{r^{n+3}} \int_{B_{r}\left(x_{0}\right)} h^{2} .
\end{aligned}
$$

We now prove the $C^{1, \beta}$ regularity of almost minimizers.
Theorem 1.4.1. Let $u$ be an almost minimizer of the Signorini problem in $B_{1}$. Define

$$
\widehat{\nabla u}\left(x^{\prime}, x_{n}\right):= \begin{cases}\nabla u\left(x^{\prime}, x_{n}\right), & x_{n} \geq 0 \\ \nabla u\left(x^{\prime},-x_{n}\right), & x_{n}<0\end{cases}
$$

Then

$$
\widehat{\nabla u} \in C^{0, \beta}\left(B_{1}\right) \quad \text { with } \beta=\frac{\alpha}{4(2 n+\alpha)} .
$$

Moreover, for any $K \Subset B_{1}$ there holds

$$
\begin{equation*}
\|\widehat{\nabla u}\|_{C^{0, \beta}(K)} \leq C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)} . \tag{1.4.2}
\end{equation*}
$$

Proof. Without loss of generality, we may assume that $K$ is a ball centered at 0 . Fix a small $r_{0}=r_{0}(n, \alpha, K)>0$ to be chosen later. Particularly, we will ask $R_{0}:=r_{0}^{\frac{2 n}{2 n+\alpha}} \leq$ $(1 / 2) \operatorname{dist}\left(K, \partial B_{1}\right)$, which will imply that

$$
\widetilde{K}:=\left\{x \in B_{1}: \operatorname{dist}(x, K) \leq R_{0}\right\} \Subset B_{1} .
$$

Our goal now is to show that for $x_{0} \in K, 0<\rho<r<r_{0}$,

$$
\begin{align*}
& \int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, \rho}\right|^{2} \leq C(n, \alpha)\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{B_{r}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, r}\right|^{2} \\
&+C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+2 \beta} \tag{1.4.3}
\end{align*}
$$

which readily gives the estimate (1.4.2) by applying Lemma 1.2.2 and using the Campanato space embedding.

We first prove (1.4.3) for $x_{0} \in K \cap B_{1}^{\prime}$, by taking the advantage of the symmetry of $\widehat{\nabla u}$, and then we argue as in the proof of Proposition 1.3.2 to extend it to all $x_{0} \in K$.

Case 1. Suppose $x_{0} \in K \cap B_{1}^{\prime}$. For notational simplicity, we assume $x_{0}=0$ (by shifting the center of the domain $D=B_{1}$ to $-x_{0}$ ) and let $0<r<r_{0}$ be given. Let us also denote

$$
\alpha^{\prime}:=1-\frac{\alpha}{8 n} \in(0,1), \quad R:=r^{\frac{2 n}{2 n+\alpha}} .
$$

We then split our proof into two cases:

$$
\sup _{\partial B_{R}}|u| \leq C_{3} R^{\alpha^{\prime}} \quad \text { and } \quad \sup _{\partial B_{R}}|u|>C_{3} R^{\alpha^{\prime}}
$$

with $C_{3}=2[u]_{0, \alpha^{\prime}, \widetilde{K}}=2 \sup _{\substack{x, y \in \widetilde{K} \\ x \neq y}} \frac{|u(x)-u(y)|}{|x-y|^{\alpha^{\prime}}}$.

Case 1.1. Assume first that $\sup _{\partial B_{R}}|u| \leq C_{3} R^{\alpha^{\prime}}$. Let $h$ be the Signorini replacement of $u$ on $B_{R}$. Then, for any $0<\rho<r$, we have

$$
\int_{B_{\rho}}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{\rho}\right|^{2} \leq 3 \int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{\rho}\right|^{2}+3 \int_{B_{\rho}}|\widehat{\nabla u}-\widehat{\nabla h}|^{2}+3 \int_{B_{\rho}}\left|\langle\widehat{\nabla u}\rangle_{\rho}-\langle\widehat{\nabla h}\rangle_{\rho}\right|^{2} .
$$

Besides, by Jensen's inequality, we have

$$
\int_{B_{\rho}}\left|\langle\widehat{\nabla u}\rangle_{\rho}-\langle\widehat{\nabla h}\rangle_{\rho}\right|^{2} \leq \int_{B_{\rho}}|\widehat{\nabla u}-\widehat{\nabla h}|^{2} .
$$

Hence, combining the estimates above, we obtain

$$
\begin{equation*}
\int_{B_{\rho}}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{\rho}\right|^{2} \leq 3 \int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{\rho}\right|^{2}+6 \int_{B_{\rho}}|\widehat{\nabla u}-\widehat{\nabla h}|^{2} . \tag{1.4.4}
\end{equation*}
$$

Similarly

$$
\begin{equation*}
\int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{r}\right|^{2} \leq 3 \int_{B_{r}}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{r}\right|^{2}+6 \int_{B_{r}}|\widehat{\nabla u}-\widehat{\nabla h}|^{2} . \tag{1.4.5}
\end{equation*}
$$

Next, note that if $r_{0} \leq(3 / 4)^{\frac{2 n+\alpha}{\alpha}}$, then $r \leq(3 / 4) R$, and thus by Proposition 1.4.1,

$$
\begin{equation*}
\int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{\rho}\right|^{2} \leq C(n, \alpha) \frac{\rho^{n+1}}{R^{n+3}} \int_{B_{R}} h^{2} \tag{1.4.6}
\end{equation*}
$$

Then, using (1.4.4), (1.4.5), and (1.4.6), we obtain

$$
\begin{align*}
\int_{B_{\rho}}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{\rho}\right|^{2} & \leq 3 \int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{\rho}\right|^{2}+6 \int_{B_{\rho}}|\widehat{\nabla u}-\widehat{\nabla h}|^{2} \\
& \leq C(n, \alpha) \frac{\rho^{n+1}}{R^{n+3}} \int_{B_{R}} h^{2}+6 \int_{B_{\rho}}|\widehat{\nabla u}-\widehat{\nabla h}|^{2} \tag{1.4.7}
\end{align*}
$$

Now take $\delta=\delta(n, \alpha, K)>0$ such that $\delta<\operatorname{dist}\left(K, \partial B_{1}\right)$ and $\delta^{\alpha} \leq \varepsilon=\varepsilon\left(C_{1}, n, n+2 \alpha^{\prime}-2\right)$, where $C_{1}$ is as in Theorem 1.3.1 and $\varepsilon$ is as in Lemma 1.2.2. If $r_{0} \leq \delta^{\frac{2 n+\alpha}{2 n}}$, then $R<\delta$, and therefore by (1.3.7),

$$
\int_{B_{R}}|\widehat{\nabla u}|^{2} \leq C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} R^{n+2 \alpha^{\prime}-2}
$$

Thus, using the above inequality, combined with (1.3.5), we obtain

$$
\begin{align*}
\int_{B_{R}}|\widehat{\nabla u}-\widehat{\nabla h}|^{2} & \leq \int_{B_{R}}|\widehat{\nabla u}|^{2}-\int_{B_{R}}|\widehat{\nabla h}|^{2} \\
& \leq R^{\alpha} \int_{B_{R}}|\widehat{\nabla h}|^{2} \leq R^{\alpha} \int_{B_{R}}|\widehat{\nabla u}|^{2}  \tag{1.4.8}\\
& \leq C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} R^{n+\alpha+2 \alpha^{\prime}-2} \\
& =C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2 n+\alpha}\left(n-\frac{1}{2}\right)} .
\end{align*}
$$

We next use that $h^{2}$ is subharmonic in $B_{R}$. (This can be seen for instance by a direct computation $\Delta\left(h^{2}\right)=2\left(|\nabla h|^{2}+h \Delta h\right)=2|\nabla h|^{2} \geq 0$, or by using the fact that $h^{ \pm}$are subharmonic.) Then,

$$
\begin{equation*}
\left\langle h^{2}\right\rangle_{R} \leq \sup _{B_{R}} h^{2}=\sup _{\partial B_{R}} h^{2}=\sup _{\partial B_{R}} u^{2} \leq C_{3}^{2} R^{2 \alpha^{\prime}} . \tag{1.4.9}
\end{equation*}
$$

Also note that by (1.3.1), $C_{3} \leq C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}$. Hence,

$$
\begin{equation*}
\frac{r^{n+1}}{R^{n+3}} \int_{B_{R}} h^{2}=C(n) \frac{r^{n+1}}{R^{3}}\left\langle h^{2}\right\rangle_{R} \leq C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} . \tag{1.4.10}
\end{equation*}
$$

Now (1.4.7), (1.4.8), (1.4.10) give

$$
\begin{align*}
\int_{B_{\rho}}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{\rho}\right|^{2} & \leq C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2}\left(r^{n+\frac{\alpha}{2(2 n+\alpha)}}+r^{n+\frac{\alpha}{2 n+\alpha}\left(n-\frac{1}{2}\right)}\right)  \tag{1.4.11}\\
& \leq C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} .
\end{align*}
$$

Case 1.2. Now suppose $\sup _{\partial B_{R}}|u|>C_{3} R^{\alpha^{\prime}}$. By the choice of $C_{3}=2[u]_{0, \alpha, \widetilde{K}}$, we have either $u \geq\left(C_{3} / 2\right) R^{\alpha^{\prime}}$ in all of $B_{R}$ or $u \leq-\left(C_{3} / 2\right) R^{\alpha^{\prime}}$ in all of $B_{R}$. However, from the inequality $u(0) \geq 0$, the only possibility is

$$
u \geq \frac{C_{3}}{2} R^{\alpha^{\prime}} \quad \text { in } B_{R}
$$

Let $h$ again be the Signorini replacement of $u$ in $B_{R}$. Then from positivity of $h=u>0$ on $\partial B_{R}$ and superharmonicity of $h$ in $B_{R}$, it follows that $h>0$ in $B_{R}$ and is therefore harmonic there. Thus,

$$
\int_{B_{\rho}}\left|\nabla h-\langle\nabla h\rangle_{\rho}\right|^{2} \leq\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\nabla h-\langle\nabla h\rangle_{r}\right|^{2}, \quad 0<\rho<r .
$$

We now want to obtain a version of this estimate for $\widehat{\nabla h}$. We start by observing that $\nabla h$ and $\widehat{\nabla h}$ differ only in the $n$-th component. The $n$-th component of $\nabla h, h_{x_{n}}$, is odd in $x_{n}$. On the other hand, the $n$-th component of $\widehat{\nabla h}$, is even in $x_{n}$ and is given by

$$
\widehat{h_{x_{n}}}(x)=\left\{\begin{array}{ll}
h_{x_{n}}\left(x^{\prime}, x_{n}\right), & x_{n} \geq 0 \\
h_{x_{n}}\left(x^{\prime},-x_{n}\right), & x_{n}<0
\end{array} .\right.
$$

Then we have

$$
\int_{B_{\rho}}\left|\widehat{h_{x_{n}}}-\left\langle\widehat{h_{x_{n}}}\right\rangle_{\rho}\right|^{2}=\int_{B_{\rho}}{\widehat{h_{x_{n}}}}^{2}-\frac{1}{\left|B_{\rho}\right|}\left(\int_{B_{\rho}} \widehat{h_{x_{n}}}\right)^{2}=\int_{B_{\rho}}\left|h_{x_{n}}-\left\langle h_{x_{n}}\right\rangle_{\rho}\right|^{2}-\frac{1}{\left|B_{\rho}\right|}\left(\int_{B_{\rho}} \widehat{h_{x_{n}}}\right)^{2},
$$

where we have used that $\left\langle h_{x_{n}}\right\rangle_{\rho}=0$. Hence, we arrive at

$$
\begin{equation*}
\int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{\rho}\right|^{2}=\int_{B_{\rho}}\left|\nabla h-\langle\nabla h\rangle_{\rho}\right|^{2}-\frac{1}{\left|B_{\rho}\right|}\left(\int_{B_{\rho}} \widehat{h_{x_{n}}}\right)^{2} . \tag{1.4.12}
\end{equation*}
$$

Similarly, we have

$$
\begin{equation*}
\int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{r}\right|^{2}=\int_{B_{r}}\left|\nabla h-\langle\nabla h\rangle_{r}\right|^{2}-\frac{1}{\left|B_{r}\right|}\left(\int_{B_{r}} \widehat{h_{x_{n}}}\right)^{2} . \tag{1.4.13}
\end{equation*}
$$

Now, using (1.4.12) and (1.4.13), we have for all $0<\rho<r$

$$
\begin{align*}
\int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{\rho}\right|^{2} & \leq \int_{B_{\rho}}\left|\nabla h-\langle\nabla h\rangle_{\rho}\right|^{2} \leq\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\nabla h-\langle\nabla h\rangle_{r}\right|^{2} \\
& \leq\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{r}\right|^{2}+\frac{1}{\left|B_{r}\right|}\left(\int_{B_{r}} \widehat{h_{x_{n}}}\right)^{2} \tag{1.4.14}
\end{align*}
$$

Next, note that if $r_{0} \leq(1 / 2)^{\frac{2 n+\alpha}{\alpha}}$, then $r \leq R / 2$. Then, for $\gamma:=1-\frac{3 \alpha}{8 n}$,

$$
\begin{aligned}
\sup _{B_{R / 2}}\left|D^{2} h\right| & \leq \frac{C(n)}{R} \sup _{B_{(3 / 4) R}}|\nabla h| \leq \frac{C(n)}{R^{1+\frac{n}{2}}}\left(\int_{B_{R}}|\nabla h|^{2}\right)^{1 / 2} \\
& \leq \frac{C(n)}{R^{1+\frac{n}{2}}}\left(\int_{B_{R}}|\nabla u|^{2}\right)^{1 / 2} \leq C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)} R^{\gamma-2},
\end{aligned}
$$

where the last inequality follows from (1.3.7). Thus, for $x=\left(x^{\prime}, x_{n}\right) \in B_{r}$, we have

$$
\begin{aligned}
\left|h_{x_{n}}\right| & \leq\left|x_{n}\right| \sup _{B_{R / 2}}\left|D^{2} h\right| \leq C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)} r R^{\gamma-2} \\
& \leq C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)} r^{1+\frac{2 n}{2 n+\alpha}(\gamma-2)},
\end{aligned}
$$

and hence

$$
\begin{align*}
\frac{1}{\left|B_{r}\right|}\left(\int_{B_{r}} \widehat{h_{x_{n}}}\right)^{2} & \leq C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+2+\frac{4 n}{2 n+\alpha}(\gamma-2)}  \tag{1.4.15}\\
& =C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} .
\end{align*}
$$

Combining (1.4.14) and (1.4.15), we obtain

$$
\begin{equation*}
\int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{\rho}\right|^{2} \leq\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{r}\right|^{2}+C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} . \tag{1.4.16}
\end{equation*}
$$

Finally, (1.4.4), (1.4.5), (1.4.8), and (1.4.16) give

$$
\begin{aligned}
\int_{B_{\rho}}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{\rho}\right|^{2} \leq & 3 \int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{\rho}\right|^{2}+6 \int_{B_{\rho}}|\widehat{\nabla u}-\widehat{\nabla h}|^{2} \\
\leq & 3\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{r}\right|^{2} \\
& +C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}}+6 \int_{B_{\rho}}|\widehat{\nabla u}-\widehat{\nabla h}|^{2} \\
\leq & 9\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{r}\right|^{2} \\
& +C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}}+24 \int_{B_{r}}|\widehat{\nabla u}-\widehat{\nabla h}|^{2}
\end{aligned}
$$

$$
\begin{aligned}
& \leq 9\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{r}\right|^{2}+C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} \\
& \quad+C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2 n+\alpha}\left(n-\frac{1}{2}\right)} \\
& \leq 9\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{r}\right|^{2}+C(n, \alpha, K)\|\nabla u\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} .
\end{aligned}
$$

From this and (1.4.11) we obtain (1.4.3) for $x_{0} \in K \cap B_{1}$.
Case 2. To extend (1.4.3) to any $x_{0} \in K$, we now assume $x_{0} \in K \cap B_{1}^{+}$. We use an argument similar to the one in Case 2 in the proof of Proposition 1.3.2.

Now, if $\rho \geq r / 4$, then

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, \rho}\right|^{2} \leq \int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, r}\right|^{2} \leq 4^{n+\alpha}\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{B_{r}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, r}\right|^{2},
$$

and thus we may assume $\rho<r / 4$. Let $d:=\operatorname{dist}\left(x_{0}, B_{1}^{\prime}\right)>0$ and choose $x_{1} \in \partial B_{d}\left(x_{0}\right) \cap B_{1}^{\prime}$. Note that from the assumption that $K$ is a ball centered at 0 , we have $x_{1} \in K \cap B_{1}^{\prime}$.

Case 2.1. If $\rho \geq d$, then from $B_{\rho}\left(x_{0}\right) \subset B_{2 \rho}\left(x_{1}\right) \subset B_{r / 2}\left(x_{1}\right) \subset B_{r}\left(x_{0}\right)$, we have

$$
\begin{aligned}
\int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, \rho}\right|^{2} \leq & \int_{B_{2 \rho}\left(x_{1}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{1}, 2 \rho}\right|^{2} \\
\leq & C(n, \alpha)\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{B_{r / 2}\left(x_{1}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{1}, r / 2}\right|^{2} \\
& \quad+C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+2 \beta} \\
\leq & C(n, \alpha)\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{B_{r}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, r}\right|^{2} \\
& \quad+C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+2 \beta}
\end{aligned}
$$

which gives (1.4.3) in this case.
Case 2.2. Now we suppose $d>\rho$. If also $d>r$, then $B_{r}\left(x_{0}\right) \subset B_{1}^{+}$and since $u$ is almost harmonic in $B_{1}^{+}$, we can apply Proposition 1.2.2, together with the growth estimate (1.3.7) in the proof of Theorem 1.3.1, to conclude

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, \rho}\right|^{2} \leq C(n, \alpha)\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{B_{r}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, r}\right|^{2}
$$

$$
+C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+2 \beta}
$$

Thus, we may assume $d \leq r$. Then, $B_{d}\left(x_{0}\right) \subset B_{1}^{+}$, and hence, again by the combination of Proposition 1.2.2 and the growth estimate (1.3.7), we have

$$
\begin{aligned}
& \int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, \rho}\right|^{2} \leq C(n, \alpha)\left(\frac{\rho}{d}\right)^{n+\alpha} \int_{B_{d}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, d}\right|^{2} \\
&+C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} d^{n+2 \beta}
\end{aligned}
$$

We need to consider further subcases.
Case 2.2.1. If $r / 4 \leq d$, then (since also $d \leq r$ )

$$
\int_{B_{d}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, d}\right|^{2} \leq 4^{n+\alpha}\left(\frac{d}{r}\right)^{n+\alpha} \int_{B_{r}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, r}\right|^{2}
$$

and combined with the previous inequality, we obtain (1.4.3) in this subcase.
Case 2.2.2. If $d<r / 4$, then we also have

$$
\begin{aligned}
\int_{B_{d}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, d}\right|^{2} \leq & \int_{B_{2 d}\left(x_{1}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{1}, 2 d}\right|^{2} \\
\leq & C(n, \alpha)\left(\frac{d}{r}\right)^{n+\alpha} \int_{B_{r / 2}\left(x_{1}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{1}, r / 2}\right|^{2} \\
& \quad+C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+2 \beta} \\
\leq & C(n, \alpha)\left(\frac{d}{r}\right)^{n+\alpha} \int_{B_{r}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, r}\right|^{2} \\
& +C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+2 \beta}
\end{aligned}
$$

Hence, the estimate (1.4.3) has been established in all possible cases.
To complete the proof of the theorem, we now apply Lemma 1.2.2 to the estimate (1.4.3) to obtain

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, \rho}\right|^{2} \leq C(n, \alpha)\left[\left(\frac{\rho}{r}\right)^{n+2 \beta} \int_{B_{r}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, r}\right|^{2}\right.
$$

$$
\left.+C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} \rho^{n+2 \beta}\right] .
$$

Taking $r \nearrow r_{0}=r_{0}(n, \alpha, K)$, we have

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla u}-\langle\widehat{\nabla u}\rangle_{x_{0}, \rho}\right|^{2} \leq C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} \rho^{n+2 \beta} .
$$

Then by the Campanato space embedding we conclude that

$$
\widehat{\nabla u} \in C^{0, \beta}(K)
$$

with

$$
\|\widehat{\nabla u}\|_{C^{0, \beta}(K)} \leq C(n, \alpha, K)\|u\|_{W^{1,2}\left(B_{1}\right)} .
$$

Having the $C^{1, \beta}$ regularity of almost minimizers, we can now talk about pointwise values of

$$
\partial_{x_{n}}^{+} u\left(x^{\prime}, 0\right)=\lim _{\substack{y \rightarrow\left(x^{\prime}, 0\right) \\ y \in B_{r}^{+}}} \partial_{x_{n}} u(y)
$$

for $x^{\prime} \in B_{1}^{\prime}$. The following complementarity condition is of crucial importance in the study of the free boundary.

Lemma 1.4.2 (Complementarity condition). Let u be an almost minimizer for the Signorini problem in $B_{1}$. Then $u$ satisfies the following complementarity condition

$$
u \partial_{x_{n}}^{+} u=0 \quad \text { on } \quad B_{1}^{\prime} .
$$

Moreover, if $x_{0} \in \Gamma(u)$ then

$$
u\left(x_{0}\right)=0 \quad \text { and } \quad\left|\widehat{\nabla u}\left(x_{0}\right)\right|=0
$$

Proof. Since $u \geq 0$ on $B_{1}^{\prime}$, the complementarity condition will follow once we show that $\partial_{x_{n}}^{+} u$ vanishes where $u>0$ on $B_{1}^{\prime}$. To this end, let $u\left(x^{\prime}, 0\right)>0$ for some $x^{\prime} \in B_{1}^{\prime}$. By the continuity of $u$ in $B_{1}$, (see Theorem 1.3.1), we have $u>0$ in some open neighborhood $U \subset B_{1}$ of ( $x^{\prime}, 0$ ).

If $B_{r}(y) \Subset U$ (not necessarily centered on $\left.B_{1}^{\prime}\right)$ and $v$ is a harmonic replacement of $u$ in $B_{r}(y)$, then by the minimum principle $v>0$ in $\overline{B_{r}(y)}$, and particularly $v>0$ on set $B_{r}(y) \cap B_{1}^{\prime}$. Then $v \in \mathfrak{K}_{0, u}\left(B_{r}(y), \mathcal{M}\right)$ and therefore we must have

$$
\int_{B_{r}(y)}|\nabla u|^{2} \leq\left(1+r^{\alpha}\right) \int_{B_{r}(y)}|\nabla v|^{2} .
$$

This means that $u$ is an almost harmonic function in $U$. Hence $u \in C^{1, \alpha / 2}(U)$ by Theorem 1.2.1. From the even symmetry of $u$ in $x_{n}$, it is then immediate that $\partial_{x_{n}}^{+} u\left(x^{\prime}, 0\right)=$ $\partial_{x_{n}} u\left(x^{\prime}, 0\right)=0$.

The second part of the lemma now follows by the $C^{1, \beta}$ regularity and the complementarity condition.

### 1.5 Weiss- and Almgren-type monotonicity formulas

In the rest of this chapter we study the free boundary of almost minimizers. In this section we introduce important technical tools, so-called Weiss- and Almgren-type monotonicity formulas, which play a significant role in our analysis.

We start with Weiss-type monotonicity formulas. They go back to the works of Weiss [49], [51] in the case of the classical obstacle problem and Alt-Caffarelli minimum problem, respectively, and to [14] for the solutions of the thin obstacle problems. In the context of almost minimizers, this type of monotonicity formulas has been used in a recent paper [43].

Theorem 1.5.1 (Weiss-type monitonicity formula). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$. For $x_{0} \in B_{1 / 2}^{\prime}$ and $0<\kappa<\kappa_{0}$ with a fixed $\kappa_{0} \geq 2$ set

$$
W_{\kappa}\left(t, u, x_{0}\right):=\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-2}}\left[\int_{B_{t}\left(x_{0}\right)}|\nabla u|^{2}-\kappa \frac{1-b t^{\alpha}}{t} \int_{\partial B_{t}\left(x_{0}\right)} u^{2}\right],
$$

with

$$
a=a_{\kappa}=\frac{n+2 \kappa-2}{\alpha}, \quad b=\frac{n+2 \kappa_{0}}{\alpha} .
$$

Then, for $0<t<t_{0}=t_{0}\left(n, \alpha, \kappa_{0}\right)$,

$$
\frac{d}{d t} W_{\kappa}\left(t, u, x_{0}\right) \geq \frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-2}} \int_{\partial B_{t}\left(x_{0}\right)}\left(u_{\nu}-\frac{\kappa\left(1-b t^{\alpha}\right)}{t} u\right)^{2}
$$

In particular, $W_{\kappa}\left(t, u, x_{0}\right)$ is nondecreasing in $t$ for $0<t<t_{0}$.

Remark 1.5.2. It is important to observe that while $a=a_{\kappa}$ depends on $\kappa$, the constant $b$ depends only on $\alpha, n$ and $\kappa_{0}$. We also note that in our version of Weiss's monotonicity formula, perturbations (from the case of the thin obstacle problem) appear in the form of multiplicative factors, rather than additive errors as in [43]. Because of the multiplicative nature of the perturbations, we can then use the one-parametric family of monotonicity formulas $\left\{W_{\kappa}\right\}_{0<\kappa<\kappa_{0}}$ to derive an Almgren-type monotonicity formula, see Theorem 1.5.4.

Remark 1.5.3. To avoid bulky notations, we will write $W_{\kappa}(t, u)$ for $W_{\kappa}\left(t, u, x_{0}\right)$ when $x_{0}=0$ or even simply $W_{\kappa}(t)$, when both $u$ and $x_{0}$ are clear from the context.

Proof. The proof uses an argument similar to the one in Theorem 1.2 in [51]. Essentially, it follows from a comparison (1.1.4) with special competitors, described below. Without loss of generality, assume $x_{0}=0$. Then for $t \in(0,1 / 2)$, define $w$ by

$$
w(x):=\left(\frac{|x|}{t}\right)^{\kappa} u\left(t \frac{x}{|x|}\right), \quad \text { for } x \in B_{t}
$$

Note that $w$ is $\kappa$-homogeneous in $B_{t}$, i.e., $w(\lambda x)=\lambda^{\kappa} w(\lambda x)$ for $\lambda>0, x, \lambda x \in B_{t}$, and coincides with $u$ on $\partial B_{t}$. Also note that $w \geq 0$ on $B_{t}^{\prime}$ and is therefore a valid competitor for $u$ in (1.1.4). We refer to this $w$ as the $\kappa$-homogeneous replacement of $u$ in $B_{t}$.

Now, in $B_{t}$, we have

$$
\nabla w(x)=\left(\frac{|x|}{t}\right)^{\kappa-1}\left[\frac{\kappa}{t} u\left(t \frac{x}{|x|}\right) \frac{x}{|x|}+\nabla u\left(t \frac{x}{|x|}\right)-\nabla u\left(t \frac{x}{|x|}\right) \cdot \frac{x}{|x|} \frac{x}{|x|}\right]
$$

which gives

$$
\int_{B_{t}}|\nabla w|^{2} d x=\int_{0}^{t} \int_{\partial B_{r}}|\nabla w(x)|^{2} d S_{x} d r
$$

$$
\begin{aligned}
& =\int_{0}^{t} \int_{\partial B_{r}}\left(\frac{r}{t}\right)^{2 \kappa-2}\left|\frac{\kappa}{t} u\left(t \frac{x}{r}\right) \nu-\left(\nabla u\left(t \frac{x}{r}\right) \cdot \nu\right) \nu+\nabla u\left(t \frac{x}{r}\right)\right|^{2} d S_{x} d r \\
& =\int_{0}^{t} \int_{\partial B_{t}}\left(\frac{r}{t}\right)^{n+2 \kappa-3}\left|\frac{\kappa}{t} u \nu-(\nabla u \cdot \nu) \nu+\nabla u\right|^{2} d S_{x} d r \\
& =\frac{t}{n+2 \kappa-2} \int_{\partial B_{t}}\left|\nabla u-(\nabla u \cdot \nu) \nu+\frac{\kappa}{t} u \nu\right|^{2} d S_{x} \\
& =\frac{t}{n+2 \kappa-2} \int_{\partial B_{t}}\left(|\nabla u|^{2}-(\nabla u \cdot \nu)^{2}+\left(\frac{\kappa}{t}\right)^{2} u^{2}\right) d S_{x} .
\end{aligned}
$$

The latter equality can be rewritten as

$$
\begin{equation*}
\int_{\partial B_{t}} u^{2} d S_{x}=\left(\frac{t}{\kappa}\right)^{2}\left[\frac{n+2 \kappa-2}{t} \int_{B_{t}}|\nabla w|^{2} d x+\int_{\partial B_{t}}\left(u_{\nu}^{2}-|\nabla u|^{2}\right) d S_{x}\right] . \tag{1.5.1}
\end{equation*}
$$

Since $w$ is a competitor for $u$, we have

$$
\begin{equation*}
\int_{B_{t}}|\nabla w|^{2} d x \geq \frac{1}{1+t^{\alpha}} \int_{B_{t}}|\nabla u|^{2} d x \geq\left(1-t^{\alpha}\right) \int_{B_{t}}|\nabla u|^{2} d x \tag{1.5.2}
\end{equation*}
$$

and combining (1.5.1) and (1.5.2) yields

$$
\begin{equation*}
\int_{\partial B_{t}} u^{2} d S_{x} \geq\left(\frac{t}{\kappa}\right)^{2}\left[(n+2 \kappa-2) \frac{1-t^{\alpha}}{t} \int_{B_{t}}|\nabla u|^{2} d x+\int_{\partial B_{t}}\left(u_{\nu}^{2}-|\nabla u|^{2}\right) d S_{x}\right] . \tag{1.5.3}
\end{equation*}
$$

Multiplying this by $\kappa^{2} \mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa}$ and rearranging terms, we obtain

$$
\begin{align*}
& \frac{d}{d t}\left(\mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa+2}\right) \int_{B_{t}}|\nabla u|^{2} d x \\
& \quad=-(n+2 \kappa-2) \mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa}\left(t-t^{\alpha+1}\right) \int_{B_{t}}|\nabla u|^{2} d x  \tag{1.5.4}\\
& \quad \geq \mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa+2} \int_{\partial B_{t}}\left(u_{\nu}^{2}-|\nabla u|^{2}\right) d S_{x}-\kappa^{2} \mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa} \int_{\partial B_{t}} u^{2} d S_{x} .
\end{align*}
$$

Define now an auxiliary function

$$
\psi(t)=\frac{\kappa \mathrm{e}^{a t^{\alpha}}\left(1-b t^{\alpha}\right)}{t^{n+2 \kappa-1}} .
$$

Then we write

$$
W_{\kappa}(t, u, 0)=\mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa+2} \int_{B_{t}}|\nabla u|^{2} d x-\psi(t) \int_{\partial B_{t}} u^{2} d S_{x}
$$

and, using (1.5.4), obtain

$$
\begin{aligned}
& \frac{d}{d t} W_{\kappa}(t, u, 0)= \frac{d}{d t}\left(\mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa+2}\right) \int_{B_{t}}|\nabla u|^{2} d x+\mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa+2} \int_{\partial B_{t}}|\nabla u|^{2} d S_{x} \\
& \quad-\psi^{\prime}(t) \int_{\partial B_{t}} u^{2} d S_{x}-2 \psi(t) \int_{\partial B_{t}} u u_{\nu} d S_{x}-(n-1) \frac{\psi(t)}{t} \int_{\partial B_{t}} u^{2} d S_{x} \\
& \geq \mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa+2} \int_{\partial B_{t}}\left(u_{\nu}^{2}-|\nabla u|^{2}\right) d S_{x}-\kappa^{2} \mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa} \int_{\partial B_{t}} u^{2} d S_{x} \\
& \quad+\mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa+2} \int_{\partial B_{t}}|\nabla u|^{2} d S_{x}-\psi^{\prime}(t) \int_{\partial B_{t}} u^{2} d S_{x} \\
& \quad-2 \psi(t) \int_{\partial B_{t}} u u_{\nu} d S_{x}-(n-1) \frac{\psi(t)}{t} \int_{\partial B_{t}} u^{2} d S_{x} \\
&= \mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa+2} \int_{\partial B_{t}} u_{\nu}^{2} d S_{x}-2 \psi(t) \int_{\partial B_{t}} u u_{\nu} d S_{x} \\
& \quad-\left(\kappa^{2} \mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa}+\psi^{\prime}(t)+(n-1) \frac{\psi(t)}{t}\right) \int_{\partial B_{t}} u^{2} d S_{x} .
\end{aligned}
$$

Now observe that $\psi(t)$ satisfies the inequality

$$
-\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-2}}\left(\kappa^{2} \mathrm{e}^{a t^{\alpha}} t^{-n-2 \kappa}+\psi^{\prime}(t)+(n-1) \frac{\psi(t)}{t}\right)-\psi^{2}(t) \geq 0
$$

for $0<t<t_{0}\left(n, \alpha, \kappa_{0}\right)$ and $0<\kappa<\kappa_{0}$. Indeed, a direct computation shows that the above inequality is equivalent to

$$
2 \alpha^{2}\left(1+\kappa_{0}-\kappa\right)-\left(n+2 \kappa_{0}\right)\left[\left(n+2 \kappa_{0}\right) \kappa-\alpha(n+2 \kappa-2)\right] t^{\alpha} \geq 0
$$

which holds for $0<\kappa<\kappa_{0}$ and small $t>0$ such that

$$
2 \alpha^{2}-4\left(n+2 \kappa_{0}\right)^{2} \kappa_{0} t^{\alpha} \geq 0
$$

Hence, recalling also the formula for $\psi(t)$, we can conclude that

$$
\begin{aligned}
\frac{d}{d t} W_{\kappa}(t, u, 0) & \geq \frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-2}}\left[\int_{\partial B_{t}} u_{\nu}^{2} d S_{x}-2 \frac{\kappa\left(1-b t^{\alpha}\right)}{t} \int_{\partial B_{t}} u u_{\nu} d S_{x}+\left(\frac{\kappa\left(1-b t^{\alpha}\right)}{t}\right)^{2} \int_{\partial B_{t}} u^{2} d S_{x}\right] \\
& =\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-2}} \int_{\partial B_{t}}\left(u_{\nu}-\frac{\kappa\left(1-b t^{\alpha}\right)}{t} u\right)^{2}
\end{aligned}
$$

for $0<t<t_{0}\left(n, \alpha, \kappa_{0}\right)$.
Next, for an almost minimizer $u$ in $B_{1}$ and $x_{0} \in B_{1 / 2}^{\prime}$, consider the quantity

$$
N\left(t, u, x_{0}\right):=\frac{t \int_{B_{t}\left(x_{0}\right)}|\nabla u|^{2}}{\int_{\partial B_{t}\left(x_{0}\right)} u^{2}}, \quad 0<t<1 / 2
$$

which is known as Almgren's frequency and goes back to Almgren's Big Regularity Paper [47]. This kind of quantities have also been used in unique continuation for a class of elliptic operators [52], [53] and have been instrumental in thin obstacle-type problems, starting with the works [12]-[14].

Before proceeding, we observe that Almgren's frequency is well defined when $x_{0}$ is a free boundary point, since $\int_{\partial B_{t}\left(x_{0}\right)} u^{2}>0$. Indeed, otherwise $u=0$ on $\partial B_{t}\left(x_{0}\right)$ and we can use $h \equiv 0$ in $B_{t}\left(x_{0}\right)$ as a competitor, to obtain that $\int_{B_{t}\left(x_{0}\right)}|\nabla u|^{2} \leq\left(1+t^{\alpha}\right) 0=0$, implying $u \equiv 0$ in $B_{t}\left(x_{0}\right)$, contradicting the assumption that $x_{0}$ is a free boundary point. Next, we also consider a modification of $N$ :

$$
\widetilde{N}\left(t, u, x_{0}\right):=\frac{1}{1-b t^{\alpha}} N\left(t, u, x_{0}\right)
$$

where $b$ is as in Theorem 1.5.1, as well as

$$
\widehat{N}_{\kappa_{0}}\left(t, u, x_{0}\right):=\min \left\{\widetilde{N}(t), \kappa_{0}\right\}, \quad 0<t<t_{0}
$$

which we call the truncated frequency.
For the frequencies $N, \widetilde{N}$, and $\widehat{N}_{\kappa_{0}}$, we will follow the same notational conventions as outlined in Remark 1.5.3 for Weiss's functionals $W_{\kappa}$.

With the Weiss type monotonicity formula at hand, we easily obtain the following monotonicity of $\widehat{N}_{\kappa_{0}}$.

Theorem 1.5.4 (Almgren-type monotonicity formula). Let $u$, $\kappa_{0}$, and $t_{0}$ be as in Theorem 1.5.1, and $x_{0}$ a free boundary point. Then $\widehat{N}_{\kappa_{0}}\left(t, u, x_{0}\right)$ is nondecreasing in $0<t<t_{0}$.

Proof. We assume $x_{0}=0$. It is quite important to observe that $t_{0}$ depends only on $n, \alpha$, and $\kappa_{0}$. Then, if $\widetilde{N}(t)<\kappa$ for some $t \in\left(0, t_{0}\right)$ and $\kappa \in\left(0, \kappa_{0}\right)$, then

$$
\begin{aligned}
W_{\kappa}(t) & =\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-1}}\left(\int_{\partial B_{t}} u^{2}\right)\left(N(t)-\kappa\left(1-b t^{\alpha}\right)\right) \\
& =\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-1}}\left(\int_{\partial B_{t}} u^{2}\right)\left(1-b t^{\alpha}\right)(\widetilde{N}(t)-\kappa)<0 .
\end{aligned}
$$

By Theorem 1.5.1 we also have $W_{\kappa}(s) \leq W_{\kappa}(t)<0$ for all $s \in(0, t)$, and thus $\widetilde{N}(s)<\kappa$. This completes the proof.

Remark 1.5.5. The proof above is rather indirect and establishes the monotonicity of $\widehat{N}_{\kappa_{0}}$ from that of Weiss-type formulas in one-parametric family $\left\{W_{\kappa}\right\}_{0<\kappa<\kappa_{0}}$. This kind of relation has been first observed in [14].

### 1.6 Almgren rescalings and blowups

In this section we prove a lower bound on Almgren's frequency for almost minimizers at free boundary points. The idea is to consider appropriate rescalings and blowups of almost minimizers to obtain solutions of the Signorini problem, for which a bound $N(0+) \geq 3 / 2$ is known.

Now, let $u$ be an almost minimizer for the Signorini problem in $B_{1}$, and $x_{0} \in B_{1 / 2}^{\prime}$ a free boundary point. For $0<r<1 / 2$ consider the Almgren rescaling ${ }^{3}$ of $u$ at $x_{0}$

$$
u_{x_{0}, r}^{A}(x):=\frac{u\left(r x+x_{0}\right)}{\left(\frac{1}{r^{n-1}} \int_{\partial B_{r}\left(x_{0}\right)} u^{2}\right)^{\frac{1}{2}}}, \quad x \in B_{1 /(2 r)} .
$$

[^1]When $x_{0}=0$, we also write $u_{r}^{A}$ instead of $u_{0, r}^{A}$. The Almgren rescalings have the following normalization and scaling properties

$$
\begin{aligned}
& \left\|u_{x_{0}, r}^{A}\right\|_{L^{2}\left(\partial B_{1}\right)}=1 \\
& N\left(\rho, u_{x_{0}, r}^{A}\right)=N\left(\rho r, u, x_{0}\right), \quad \rho<1 /(2 r) .
\end{aligned}
$$

We will call the limits of $u_{x_{0}, r}^{A}$ over any sequence $r=r_{\mathrm{j}} \rightarrow 0+$ Almgren blowups of $u$ at $x_{0}$ and denote by $u_{x_{0}, 0}^{A}$.

Proposition 1.6.1 (Existence of Almgren blowups). Let $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(u)$ be such that $\widehat{N}_{\kappa_{0}}\left(0+, u, x_{0}\right)=\kappa<\kappa_{0}$. Then every sequence of Almgren rescalings $u_{x_{0}, r_{\mathrm{j}}}^{A}$, with $r_{\mathrm{j}} \rightarrow 0+$ contains a subsequence, still denoted $r_{\mathrm{j}}$, such that for a function $u_{x_{0}, 0}^{A} \in W^{1,2}\left(B_{1}\right) \cap C_{\mathrm{loc}}^{1}\left(B_{1}^{ \pm} \cup\right.$ $B_{1}^{\prime}$ )

$$
\begin{array}{ll}
u_{x_{0}, r_{\mathrm{j}}}^{A} \rightarrow u_{x_{0}, 0}^{A} & \text { in } W^{1,2}\left(B_{1}\right), \\
u_{x_{0}, r_{\mathrm{j}}}^{A} \rightarrow u_{x_{0}, 0}^{A} & \text { in } L^{2}\left(\partial B_{1}\right), \\
u_{x_{0}, r_{\mathrm{j}}}^{A} \rightarrow u_{x_{0}, 0}^{A} & \text { in } C_{\mathrm{loc}}^{1}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right) .
\end{array}
$$

Moreover, $u_{x_{0}, 0}^{A}$ is a nonzero solution of the Signorini problem in $B_{1}$, even in $x_{n}$, and homogeneous of degree $\kappa$ in $B_{1}$, i.e.,

$$
u_{x_{0}, 0}^{A}(\lambda x)=\lambda^{\kappa} u_{x_{0}, 0}^{A}(x),
$$

for $\lambda>0$, provided $x, \lambda x \in B_{1}$.

Proof. Without loss of generality, we assume $x_{0}=0$. From the fact that $\widehat{N}(0+, u)=\kappa<\kappa_{0}$, it follows also that $N(0+, u)=\widehat{N}(0+, u)=\kappa$. In particular, $N\left(r_{\mathrm{j}}, u\right)<\kappa_{0}$ for large j . Then, for such j

$$
\int_{B_{1}}\left|\nabla u_{r_{\mathrm{j}}}^{A}\right|^{2}=N\left(1, u_{r_{\mathrm{j}}}^{A}\right)=N\left(r_{\mathrm{j}}, u\right) \leq \kappa_{0}
$$

and combined with the normalization $\int_{\partial B_{1}}\left(u_{r_{\mathrm{j}}}^{A}\right)^{2}=1$, we see that the sequence $u_{r_{\mathrm{j}}}^{A}$ is bounded in $W^{1,2}\left(B_{1}\right)$. Hence, there is a function $u_{0}^{A} \in W^{1,2}\left(B_{1}\right)$ such that, over a subsequence,

$$
\begin{aligned}
& u_{r_{\mathrm{j}}}^{A} \rightarrow u_{0}^{A} \quad \text { weakly in } W^{1,2}\left(B_{1}\right), \\
& u_{r_{\mathrm{j}}}^{A} \rightarrow u_{0}^{A} \quad \text { strongly in } L^{2}\left(\partial B_{1}\right) .
\end{aligned}
$$

In particular, $\int_{\partial B_{1}}\left(u_{0}^{A}\right)^{2}=1$, implying that $u_{0}^{A} \not \equiv 0$ in $B_{1}$.
Next, we observe that since $u$ is an almost minimizer in $B_{1}$ with gauge function $\omega(t)=t^{\alpha}$, $u_{r}^{A}$ is also an almost minimizer in $B_{1 /(2 r)}$ with gauge function $\omega_{r}(t)=(r t)^{\alpha}$. This is rather easy to see, since $u_{r}^{A}(x)$ up to a positive constant factor is $u(r x)$ and the multiplication (or the division) by a positive number preserves the almost minimizing property. Since $\omega_{r}(t) \leq \omega(t)$, Theorem 1.4.1 is applicable to rescalings $u_{r_{j}}^{A}$, from where we can deduce that over yet another subsequence,

$$
\begin{equation*}
u_{r_{\mathrm{j}}}^{A} \rightarrow u_{0}^{A} \quad \text { in } C_{\mathrm{loc}}^{1}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right) \tag{1.6.1}
\end{equation*}
$$

Now, we claim that since the gauge functions $\omega_{r}(t)=(r t)^{\alpha} \rightarrow 0$ as $r \rightarrow 0$, the blowup $u_{0}^{A}$ is a solution of the Signorini problem in $B_{1}$. Indeed, for a fixed $r_{\mathrm{j}}$, let $h_{r_{\mathrm{j}}}$ be the Signorini replacement of $u_{r_{\mathrm{j}}}^{A}$ in $B_{1}$. Then, by repeating the argument as in the proof of Proposition 1.3.2

$$
\int_{B_{1}}\left|\nabla\left(u_{r_{\mathrm{j}}}^{A}-h_{r_{\mathrm{j}}}\right)\right|^{2} \leq r_{\mathrm{j}}^{\alpha} \int_{B_{1}}\left|\nabla u_{r_{\mathrm{j}}}^{A}\right|^{2} .
$$

This implies that $h_{r_{\mathrm{j}}} \rightarrow u_{0}^{A}$ weakly in $W^{1,2}\left(B_{1}\right)$. On the other hand, by the boundedness of the sequence $h_{r_{\mathrm{j}}}$ in $W^{1,2}\left(B_{1}\right)$, we have also boundedness in $C_{\mathrm{loc}}^{1,1 / 2}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)$ and hence, over a subsequence, $h_{r_{\mathrm{j}}} \rightarrow u_{0}^{A}$ in $C_{\mathrm{loc}}^{1}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)$. By this convergence we then conclude that $u_{0}^{A}$ satisfies

$$
\begin{aligned}
\Delta u_{0}^{A}=0 & \text { in } B_{1} \backslash B_{1}^{\prime} \\
u_{0}^{A} \geq 0, & -\partial_{x_{n}}^{+} u_{0}^{A} \geq 0, \quad u_{0}^{A} \partial_{x_{n}}^{+} u_{0}^{A}=0
\end{aligned} \quad \text { on } B_{1}^{\prime}, ~ \$ ~ \$
$$

and hence $u_{0}^{A}$ itself solves the Signorini problem in $B_{1}$.

Using the $C_{\text {loc }}^{1}$ convergence again, we have that for any $0<\rho<1$

$$
N\left(\rho, u_{0}^{A}\right)=\lim _{r_{\mathrm{j}} \rightarrow 0} N\left(\rho, u_{r_{\mathrm{j}}}^{A}\right)=\lim _{r_{\mathrm{j}} \rightarrow 0} N\left(\rho r_{\mathrm{j}}, u\right)=N(0+, u)=\kappa .
$$

Thus, the Almgren frequency of $u_{0}^{A}$ is constant $\kappa$, which is possible only if $u_{0}^{A}$ is a $\kappa$ homogeneous solution of the Signorini problem in $B_{1}$, see Theorem 9.4 in [48].

In what follows, it will be sufficient for us to fix $\kappa_{0} \geq 2$ (say $\kappa_{0}=2$ ), in the definition of $\widehat{N}_{\kappa_{0}}$ and we will simply write

$$
\widehat{N}=\widehat{N}_{\kappa_{0}} .
$$

Lemma 1.6.1 (Minimal frequency). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$. If $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(u)$, then

$$
\widehat{N}\left(0+, u, x_{0}\right)=\lim _{r \rightarrow 0+} \widehat{N}\left(r, u, x_{0}\right) \geq \frac{3}{2} .
$$

Consequently, we also have

$$
\widehat{N}\left(t, u, x_{0}\right) \geq 3 / 2 \quad \text { for } 0<t<t_{0}
$$

Proof. As before, let $x_{0}=0$. Assume to the contrary that $\widehat{N}(0+, u)=\kappa<3 / 2$. Since $\kappa<\kappa_{0}$ we can apply Proposition 1.6.1 to obtain that over a sequence $r_{\mathrm{j}} \rightarrow 0+, u_{r_{\mathrm{j}}}^{A} \rightarrow u_{0}^{A}$ in $C_{\text {loc }}^{1}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)$, where $u_{0}^{A}$ is a nonzero $\kappa$-homogeneous solution of the Signorini problem in $B_{1}$, even in $x_{n}$. Moreover, since $0 \in \Gamma(u)$, by Lemma 1.4.2 we have that $u(0)=|\widehat{\nabla u}(0)|=0$, implying that $u_{r_{\mathrm{j}}}^{A}(0)=\left|\widehat{\nabla u_{r_{j}}^{A}}(0)\right|=0$ and, by passing to the limit, $u_{0}^{A}(0)=\left|\widehat{\nabla u_{0}^{A}}(0)\right|=0$. Now, to arrive at a contradiction, we argue as in the proof of Proposition 9.9 in [48] to reduce the problem to dimension $n=2$, where we can classify all possible homogeneous solutions of the Signorini problem, even in $x_{n}$. The only nonzero homogeneous solutions with $\kappa<3 / 2$ in dimension $n=2$ are possible for $\kappa=1$ and have the form $u_{0}^{A}(x)=-c x_{n}$ for some $c>0$, but they fails to satisfy the condition $\left|\widehat{\nabla u_{0}^{A}}(0)\right|=0$. Thus, we arrived at contradiction, implying that $\widehat{N}(0+, u) \geq 3 / 2$. Finally, applying Theorem 1.5.4, we obtain $\widehat{N}(t, u) \geq \widehat{N}(0+, u) \geq 3 / 2$, for $0<t<t_{0}$.

Corollary 1.6.2. Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$ and $x_{0} a$ free boundary points. Then

$$
W_{3 / 2}\left(t, u, x_{0}\right) \geq 0 \quad \text { for } 0<t<t_{0} .
$$

Proof. We simply observe that $\widetilde{N}(t) \geq \widehat{N}(t) \geq 3 / 2$ for $0<t<t_{0}$ and hence

$$
W_{3 / 2}\left(t, u, x_{0}\right)=\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-1}}\left(\int_{\partial B_{t}} u^{2}\right)\left(1-b t^{\alpha}\right)\left(\widetilde{N}(t)-\frac{3}{2}\right) \geq 0
$$

### 1.7 Growth estimates

An important step in the study of the free boundary in the Signorini problem (and in many other free boundary problems) is the proof of the optimal regularity of solutions, which in this case is $C^{1,1 / 2}$ on each side of the thin space. This allows to make proper blowup arguments to establish the regularity of the so-called regular part of the free boundary. However, in the case of almost minimizers, we only know $C^{1, \beta}$ regularity for some small $\beta>0$ and do not expect to have anything better. Yet, in this section, we establish the optimal growth of the almost minimizers at free boundary points with the help of the Weisstype monotonicity formula and the epiperimetric inequality.

Finally, we want to point out that the results in this section are rather immediate in the case of minimizers, as they follow easily from the differentiation formulas for the quantities involved in the Almgren's frequency formula. This is completely unavailable for almost minimizers.

We start by defining a new type of rescalings. Fix $\kappa \geq 3 / 2$. For a free boundary point $x_{0}$ in $B_{1 / 2}^{\prime}$ and $r>0$, we define the $\kappa$-homogeneous rescaling by

$$
u_{x_{0}, r}(x):=u_{x_{0}, r}^{(\kappa)}(x)=\frac{u\left(r x+x_{0}\right)}{r^{\kappa}}, \quad x \in B_{1 /(2 r)} .
$$

To take advantage of the Weiss-type monotonicity formula, we need a slight modification of this rescaling. With the help of an auxiliary function

$$
\phi(r)=\phi_{\kappa}(r):=\mathrm{e}^{-(\kappa b / \alpha) r^{\alpha}} r^{\kappa}, \quad r>0,
$$

which is a solution of the differential equation

$$
\phi^{\prime}(r)=\kappa \phi(r) \frac{1-b r^{\alpha}}{r}, \quad r>0
$$

we define the $\kappa$-almost homogeneous rescalings by

$$
u_{x_{0}, r}^{\phi}(x):=\frac{u\left(r x+x_{0}\right)}{\phi(r)}, \quad x \in B_{1 /(2 r)} .
$$

Lemma 1.7.1 (Weak growth estimate). Let $u$ be an almost minimizer of the Signorini problem in $B_{1}$ and $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(u)$ be such that $\widehat{N}\left(0+, u, x_{0}\right) \geq \kappa$ for $\kappa \leq \kappa_{0}$. Then

$$
\begin{aligned}
\int_{\partial B_{t}\left(x_{0}\right)} u^{2} & \leq C\left(n, \alpha, \kappa_{0}\right)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2}\left(\log \frac{1}{t}\right) t^{n+2 \kappa-1}, \\
\int_{B_{t}\left(x_{0}\right)}|\nabla u|^{2} & \leq C\left(n, \alpha, \kappa_{0}\right)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2}\left(\log \frac{1}{t}\right) t^{n+2 \kappa-2},
\end{aligned}
$$

for $0<t<t_{0}=t_{0}\left(n, \alpha, \kappa_{0}\right)$.
Proof. Without loss of generality, assume $x_{0}=0$. We first note that the condition $\widehat{N}(0+, u) \geq$ $\kappa$ implies that $\widehat{N}(t, u) \geq \kappa$ for $0<t<t_{0}=t_{0}\left(n, \alpha, \kappa_{0}\right)$. Then also $\widetilde{N}(t, u) \geq \kappa$ for such $t$ and consequently,

$$
W_{\kappa}(t, u)=\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-1}}\left(\int_{\partial B_{t}} u^{2}\right)\left(1-b t^{\alpha}\right)(\widetilde{N}(t, u)-\kappa) \geq 0 .
$$

Next, for $\phi=\phi_{\kappa}$, we have that

$$
\begin{aligned}
\frac{d}{d r} u_{r}^{\phi}(x) & =\frac{\nabla u(r x) \cdot x}{\phi(r)}-\frac{u(r x)\left[\phi^{\prime}(r) / \phi(r)\right]}{\phi(r)} \\
& =\frac{1}{\phi(r)}\left(\nabla u(r x) \cdot x-\frac{\kappa\left(1-b r^{\alpha}\right)}{r} u(r x)\right) .
\end{aligned}
$$

Now let

$$
m(r)=\left(\int_{\partial B_{1}}\left(u_{r}^{\phi}(\xi)\right)^{2} d S_{\xi}\right)^{1 / 2}, \quad r>0 .
$$

Then,

$$
m^{\prime}(r)=\left(\int_{\partial B_{1}} u_{r}^{\phi}(\xi) \frac{d}{d r} u_{r}^{\phi}(\xi) d S_{\xi}\right)\left(\int_{\partial B_{1}}\left(u_{r}^{\phi}(\xi)\right)^{2} d S_{\xi}\right)^{-1 / 2}
$$

and consequently, by Cauchy-Schwarz,

$$
\left|m^{\prime}(r)\right| \leq\left(\int_{\partial B_{1}}\left[\frac{d}{d r} u_{r}^{\phi}(\xi)\right]^{2} d S_{\xi}\right)^{1 / 2}
$$

Hence,

$$
\begin{aligned}
\left|m^{\prime}(r)\right| & \leq \frac{1}{\phi(r)}\left(\int_{\partial B_{1}}\left(\nabla u(r \xi) \cdot \xi-\frac{\kappa\left(1-b r^{\alpha}\right)}{r} u(r \xi)\right)^{2} d S_{\xi}\right)^{1 / 2} \\
& =\frac{1}{\phi(r)}\left(\frac{1}{r^{n-1}} \int_{\partial B_{r}}\left(\partial_{\nu} u(x)-\frac{\kappa\left(1-b r^{\alpha}\right)}{r} u(x)\right)^{2} d S_{x}\right)^{1 / 2} \\
& \leq \frac{1}{\phi(r)}\left(\frac{1}{r^{n-1}} \frac{r^{n+1}}{\mathrm{e}^{a r^{\alpha}}} \frac{d}{d r} W_{\kappa}(r)\right)^{1 / 2}=\frac{\mathrm{e}^{c r^{\alpha}}}{r^{1 / 2}}\left(\frac{d}{d r} W_{\kappa}(r)\right)^{1 / 2}, \quad c=\kappa \frac{b}{\alpha}-\frac{a}{2},
\end{aligned}
$$

for $0<r<t_{0}=t_{0}\left(n, \alpha, \kappa_{0}\right)$. Thus, we have shown

$$
\left|m^{\prime}(r)\right| \leq \frac{\mathrm{e}^{c r^{\alpha}}}{r^{1 / 2}}\left(\frac{d}{d r} W_{\kappa}(r)\right)^{1 / 2}, \quad 0<r<t_{0}
$$

Integrating in $r$ over the interval $(s, t) \subset\left(0, t_{0}\right)$, we obtain

$$
\begin{aligned}
|m(t)-m(s)| & \leq \int_{s}^{t} \frac{\mathrm{e}^{c r^{\alpha}}}{r^{1 / 2}}\left(\frac{d}{d r} W_{\kappa}(r)\right)^{1 / 2} d r \leq\left(\int_{s}^{t} \frac{\mathrm{e}^{2 c r^{\alpha}}}{r} d r\right)^{1 / 2}\left(\int_{s}^{t} \frac{d}{d r} W_{\kappa}(r)\right)^{1 / 2} \\
& \leq C_{0}\left(\log \frac{t}{s}\right)^{1 / 2}\left[W_{\kappa}(t)-W_{\kappa}(s)\right]^{1 / 2}
\end{aligned}
$$

In particular (recalling that $W_{\kappa}(s) \geq 0$ ), we obtain

$$
m(t) \leq m\left(t_{0}\right)+C_{0}\left(\log \frac{t_{0}}{t}\right)^{1 / 2}\left[W_{\kappa}\left(t_{0}\right)\right]^{1 / 2}
$$

Varying $t_{0}$ by an absolute factor, we can guarantee that

$$
m\left(t_{0}\right) \leq C\left(n, \alpha, \kappa_{0}\right)\|u\|_{L^{2}\left(B_{1}\right)}, \quad W_{\kappa}\left(t_{0}\right) \leq C\left(n, \alpha, \kappa_{0}\right)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} .
$$

Hence, we can conclude

$$
\int_{\partial B_{t}} u^{2} \leq C\left(n, \alpha, \kappa_{0}\right)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2}\left(\log \frac{1}{t}\right) t^{n+2 \kappa-1}
$$

for $0<t<t_{0}=t_{0}\left(n, \alpha, \kappa_{0}\right)$. This implies the first bound. The second bound follows immediately from the first one by using that $W_{\kappa}(t, u) \leq W_{\kappa}\left(t_{0}, u\right)$ :

$$
\begin{aligned}
\frac{1}{t^{n+2 \kappa-2}} \int_{B_{t}}|\nabla u|^{2} & \leq \frac{\kappa\left(1-b t^{\alpha}\right)}{t^{n+2 \kappa-1}} \int_{\partial B_{t}} u^{2}+\mathrm{e}^{-a t^{\alpha}} W_{\kappa}\left(t_{0}, u\right) \\
& \leq C\left(n, \alpha, \kappa_{0}\right)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2}\left(\log \frac{1}{t}\right)+\frac{\mathrm{e}^{a t_{0}^{\alpha}}}{t_{0}^{n+2 \kappa-2}} \int_{B_{t_{0}}}|\nabla u|^{2} \\
& \leq C\left(n, \alpha, \kappa_{0}\right)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2}\left(\log \frac{1}{t}\right) .
\end{aligned}
$$

The logarithmic term in Lemma 1.7.1 does not allow to conclude that the sequence of $\kappa$-homogeneous or almost homogeneous rescaling is uniformly bounded say in $W^{1,2}\left(B_{1}\right)$. In the rest of this section we show that in the case of the minimal frequency $\kappa=3 / 2$ we can do that with the help of the so-called epiperimetric inequality for the Signorini problem for the Weiss energy

$$
W_{3 / 2}^{0}(w):=\int_{B_{1}}|\nabla w|^{2}-\frac{3}{2} \int_{\partial B_{1}} w^{2} .
$$

To state this result, we let

$$
\begin{equation*}
\mathcal{A}:=\left\{w \in W^{1,2}\left(B_{1}\right): w \geq 0 \text { on } B_{1}^{\prime}, w\left(x^{\prime}, x_{n}\right)=w\left(x^{\prime},-x_{n}\right)\right\} \tag{1.7.1}
\end{equation*}
$$

Theorem 1.7.2 (Epiperimetric inequality). There exists $\eta \in(0,1)$ such that if $w \in \mathcal{A}$ is homogeneous of degree $3 / 2$ in $B_{1}$, then there exists $v \in \mathcal{A}$ with $v=w$ on $\partial B_{1}$ such that

$$
W_{3 / 2}^{0}(v) \leq(1-\eta) W_{3 / 2}^{0}(w)
$$

This kind of inequalities go back to the work of Weiss [49], in the case of the classical obstacle problem. For the Signorini problem, a version of this theorem was proved in [20] and [54]. In fact, the theorem above is the version in [55]. The inequality in [20] and [54] requires $w$ to be close to the blowup profile, but this can be easily removed by a scaling argument (see [55]). We also refer to [27], for a more direct proof of this inequality with an explicit constant $\eta=1 /(2 n+3)$.

Now, with the help of the epiperimetric inequality, we can prove a decay estimate for the Weiss-type energy functional $W_{3 / 2}$. For the rest of the section, we will assume

$$
\kappa_{0}=2,
$$

which will make some of the constants independent of $\kappa_{0}$, but the results hold also for any other value of $\kappa_{0} \geq 2$, with possible added dependence of constants on $\kappa_{0}$.

Lemma 1.7.3. Let $x_{0} \in B_{1 / 2}^{\prime}$ be a free boundary point. Then, there exist $\delta=\delta(n, \alpha)>0$ such that

$$
0 \leq W_{3 / 2}\left(t, u, x_{0}\right) \leq C t^{\delta}, \quad 0<t<t_{0}=t_{0}(n, \alpha)
$$

with $C=C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2}$.
Proof. As before, without loss of generality we assume that $x_{0}=0$.
The proof will follow from a differential inequality that we derive by using our earlier computations and the epiperimetric inequality. Recalling the proof of the Weiss-type monotonicity formula (Theorem 1.5.1), for small $t>0$, we have

$$
\begin{aligned}
& \frac{d}{d t} W_{3 / 2}(t, u)= \frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+1}} \int_{\partial B_{t}}|\nabla u|^{2}-\frac{(n+1)\left(1-t^{\alpha}\right) \mathrm{e}^{a t^{\alpha}}}{t^{n+2}} \int_{B_{t}}|\nabla u|^{2} \\
& \quad-\psi^{\prime}(t) \int_{\partial B_{t}} u^{2}-(n-1) \frac{\psi(t)}{t} \int_{\partial B_{t}} u^{2}-2 \psi(t) \int_{\partial B_{t}} u \partial_{\nu} u
\end{aligned}
$$

$$
\begin{aligned}
=- & \frac{(n+1)\left(1-t^{\alpha}\right)}{t} W_{3 / 2}(t, u)+\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+1}} \int_{\partial B_{t}}|\nabla u|^{2} \\
& -\left(\left[(n+1)\left(1-t^{\alpha}\right)+(n-1)\right] \frac{\psi(t)}{t}+\psi^{\prime}(t)\right) \int_{\partial B_{t}} u^{2}-2 \psi(t) \int_{\partial B_{t}} u \partial_{\nu} u \\
\geq- & \frac{(n+1)\left(1-t^{\alpha}\right)}{t} W_{3 / 2}(t, u) \\
& +\frac{\mathrm{e}^{a t^{\alpha}}\left(1-b t^{\alpha}\right)}{t^{n+1}} \int_{\partial B_{t}}\left(|\nabla u|^{2}-\frac{3}{t} u \partial_{\nu} u\right. \\
& \left.\quad-\frac{3}{2 t}\left[\frac{(n+1)\left(1-t^{\alpha}\right)+(n-1)}{t}+\frac{\psi^{\prime}(t)}{\psi(t)}\right] u^{2}\right)
\end{aligned}
$$

To proceed, note that

$$
\frac{(n+1)\left(1-t^{\alpha}\right)+(n-1)}{t}+\frac{\psi^{\prime}(t)}{\psi(t)}=\frac{(n-2)+O\left(t^{\alpha}\right)}{t}
$$

Now, for the homogeneous rescalings

$$
u_{t}(x)=\frac{u(t x)}{t^{3 / 2}}
$$

we can write

$$
\begin{aligned}
\int_{\partial B_{t}} & |\nabla u|^{2}-\frac{3}{t} u \partial_{\nu} u-\frac{3}{2} \frac{(n-2)+O\left(t^{\alpha}\right)}{t^{2}} u^{2} \\
& =t^{n} \int_{\partial B_{1}}\left|\nabla u_{t}\right|^{2}-3 u_{t} \partial_{\nu} u_{t}-\frac{3}{2}\left[(n-2)+O\left(t^{\alpha}\right)\right] u_{t}^{2} \\
& =t^{n} \int_{\partial B_{1}}\left(\partial_{\nu} u_{t}-\frac{3}{2} u_{t}\right)^{2}+\left(\partial_{\tau} u_{t}\right)^{2}-\frac{3}{2}\left[\left(n-\frac{1}{2}\right)+O\left(t^{\alpha}\right)\right] u_{t}^{2}
\end{aligned}
$$

where $\partial_{\tau} u_{t}$ is the tangential component of $\nabla u_{t}$ on the unit sphere. We can summarize for now that

$$
\begin{aligned}
\frac{d}{d t} W_{3 / 2}(t, u) \geq- & \frac{(n+1)\left(1-t^{\alpha}\right)}{t} W_{3 / 2}(t, u) \\
& +\frac{\mathrm{e}^{a t^{\alpha}}\left(1-b t^{\alpha}\right)}{t} \int_{\partial B_{1}}\left[\left(\partial_{\nu} u_{t}-\frac{3}{2} u_{t}\right)^{2}+\left(\partial_{\tau} u_{t}\right)^{2}-\frac{3}{2}\left(n-\frac{1}{2}\right) u_{t}^{2}\right] \\
& +O\left(t^{\alpha-1}\right) \int_{\partial B_{1}} u_{t}^{2}
\end{aligned}
$$

On the other hand, if $w_{t}$ is a $3 / 2$-homogeneous replacement of $u_{t}$ in $B_{1}$, i.e.,

$$
w_{t}(x)=|x|^{3 / 2} u_{t}(x /|x|)
$$

then

$$
\int_{\partial B_{1}}\left(\partial_{\tau} u_{t}\right)^{2}-\frac{3}{2}\left(n-\frac{1}{2}\right) u_{t}^{2}=\int_{\partial B_{1}}\left(\partial_{\tau} w_{t}\right)^{2}-\frac{3}{2}\left(n-\frac{1}{2}\right) w_{t}^{2}=(n+1) W_{3 / 2}^{0}\left(w_{t}\right),
$$

where

$$
W_{3 / 2}^{0}\left(w_{t}\right)=\int_{B_{1}}\left|\nabla w_{t}\right|^{2}-\frac{3}{2} \int_{\partial B_{1}} w_{t}^{2}
$$

The last equality follows by repeating the arguments in the beginning of the proof of Theorem 1.5.1 with $\kappa=3 / 2$. Let $v_{t}$ be the solution of the Signorini problem in $B_{1}$ with $v_{t}=u_{t}=w_{t}$ on $\partial B_{1}$. Then by the epiperimetric inequality

$$
W_{3 / 2}^{0}\left(v_{t}\right) \leq(1-\eta) W_{3 / 2}^{0}\left(w_{t}\right)
$$

On the other hand, since $u$ is an almost minimizer, we have

$$
\int_{B_{1}}\left|\nabla u_{t}\right|^{2} \leq\left(1+t^{\alpha}\right) \int_{B_{1}}\left|\nabla v_{t}\right|^{2}
$$

and since also $u_{t}=v_{t}$ on $\partial B_{1}$, we have

$$
\begin{aligned}
W_{3 / 2}(t, u) & =\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+1}}\left[\int_{B_{t}}|\nabla u|^{2}-\frac{(3 / 2)\left(1-b t^{\alpha}\right)}{t} \int_{\partial B_{t}} u^{2}\right] \\
& \leq\left(1+O\left(t^{\alpha}\right)\right) W_{3 / 2}^{0}\left(v_{t}\right)+O\left(t^{\alpha}\right) \int_{\partial B_{1}} u_{t}^{2} \\
& \leq\left(1-\frac{\eta}{2}\right) W_{3 / 2}^{0}\left(w_{t}\right)+O\left(t^{\alpha}\right) \int_{\partial B_{1}} u_{t}^{2}, \quad \text { for } 0<t<t_{0}=t_{0}(n, \alpha) .
\end{aligned}
$$

We can therefore write

$$
\begin{aligned}
\frac{d}{d t} W_{3 / 2}(t, u) \geq- & \frac{(n+1)\left(1-t^{\alpha}\right)}{t} W_{3 / 2}(t, u) \\
& +\frac{(n+1) \mathrm{e}^{a t^{\alpha}}\left(1-b t^{\alpha}\right)}{t} W_{3 / 2}^{0}\left(w_{t}\right)+O\left(t^{\alpha-1}\right) \int_{\partial B_{1}} u_{t}^{2}
\end{aligned}
$$

$$
\begin{aligned}
& \geq \frac{n+1}{t}\left(-1+\frac{1}{1-\eta / 2}+O\left(t^{\alpha}\right)\right) W_{3 / 2}(t, u)+\frac{O\left(t^{\alpha}\right)}{t^{n+3}} \int_{\partial B_{t}} u^{2} \\
& \geq \frac{\eta}{4 t} W_{3 / 2}(t, u)-C t^{\alpha / 2-1}
\end{aligned}
$$

for small $t$, where we have also used the growth estimate in Lemma 1.7.1. Taking now $\delta$ such that

$$
0<\delta<\min \left\{\frac{\eta}{4}, \frac{\alpha}{2}\right\}
$$

we have

$$
\begin{aligned}
\frac{d}{d t}\left[W_{3 / 2}(t, u) t^{-\delta}+\frac{C}{\alpha / 2-\delta} t^{\alpha / 2-\delta}\right] & =t^{-\delta}\left(\frac{d}{d t} W_{3 / 2}(t, u)-\frac{\delta}{t} W_{3 / 2}(t, u)\right)+C t^{\alpha / 2-\delta-1} \\
& \geq t^{-\delta-1}\left[\frac{\eta}{4}-\delta\right] W_{3 / 2}(t, u)-C t^{\alpha / 2-\delta-1}+C t^{\alpha / 2-\delta-1} \\
& \geq 0
\end{aligned}
$$

for small $t$, where we have used again that $W_{3 / 2}(t, u) \geq 0$. Thus, we can conclude that

$$
0 \leq W_{3 / 2}(t, u) \leq C t^{\delta}, \quad 0<t<t_{0}=t_{0}(n, \alpha)
$$

with $C=C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2}$.
Using the estimate on $W_{3 / 2}(t, u)$ in Lemma 1.7.3, we can improve on Lemma 1.7.1 in the case $\kappa=3 / 2$.

Lemma 1.7.4 (Optimal growth estimate). Let $x_{0} \in B_{1 / 2}^{\prime}$ be a free boundary point. Then, for $0<t<t_{0}=t_{0}(n, \alpha)$,

$$
\begin{aligned}
\int_{\partial B_{t}\left(x_{0}\right)} u^{2} & \leq C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} t^{n+2} \\
\int_{B_{t}\left(x_{0}\right)}|\nabla u|^{2} & \leq C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} t^{n+1} .
\end{aligned}
$$

Proof. We proceed as in the proof of Lemma 1.7.1 up to the estimate

$$
|m(t)-m(s)| \leq C_{0}\left(\log \frac{t}{s}\right)^{1 / 2}\left[W_{3 / 2}(t, u)-W_{3 / 2}(s, u)\right]^{1 / 2}
$$

From there, using Lemma 1.7.3, we now have an improved bound

$$
|m(t)-m(s)| \leq C\left(\log \frac{t}{s}\right)^{1 / 2} t^{\delta / 2}, \quad s<t<t_{0}
$$

with $C=C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}$. Then, by a dyadic argument, we can conclude that

$$
|m(t)-m(s)| \leq C t^{\delta / 2}
$$

Indeed, let $k=0,1,2, \ldots$ be such that $t / 2^{k+1} \leq s<t / 2^{k}$. Then,

$$
\begin{aligned}
|m(t)-m(s)| & \leq \sum_{\mathrm{j}=1}^{k}\left|m\left(t / 2^{\mathrm{j}-1}\right)-m\left(t / 2^{\mathrm{j}}\right)\right|+\left|m\left(t / 2^{k}\right)-m(s)\right| \\
& \leq C(\log 2)^{1 / 2} \sum_{\mathrm{j}=1}^{k+1}\left(t / 2^{\mathrm{j}-1}\right)^{\delta / 2} \leq C(\log 2)^{1 / 2} \frac{t^{\delta / 2}}{1-2^{-\delta / 2}}=C t^{\delta / 2}
\end{aligned}
$$

In particular, we have

$$
m(t) \leq m\left(t_{0}\right)+C t_{0}^{\delta / 2} \leq C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}, \quad t<t_{0}
$$

This implies the first bound. The second bound follows immediately from the first one by using that $W_{3 / 2}(t, u) \leq W_{3 / 2}\left(t_{0}, u\right)$ :

$$
\begin{aligned}
\frac{1}{t^{n+1}} \int_{B_{t}}|\nabla u(x)|^{2} d x & \leq \frac{(3 / 2)\left(1-b t^{\alpha}\right)}{t^{n+2}} \int_{\partial B_{t}} u(x)^{2} d S_{x}+\mathrm{e}^{-a t^{\alpha}} W_{3 / 2}\left(t_{0}, u\right) \\
& \leq C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2}+\frac{\mathrm{e}^{a t_{0}^{\alpha}}}{t_{0}^{n+1}} \int_{B_{t_{0}}}|\nabla u(x)|^{2} d x \\
& \leq C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} .
\end{aligned}
$$

### 1.8 3/2-Homogeneous blowups

For a free boundary point $x_{0} \in B_{1 / 2}^{\prime}$, we consider again the $3 / 2$-almost homogeneous rescalings

$$
u_{x_{0}, t}^{\phi}(x)=\frac{u\left(t x+x_{0}\right)}{\phi(t)}, \quad x \in B_{1 /(2 t)}
$$

with $\phi=\phi_{3 / 2}$. We now observe that the optimal growth estimates in Lemma 1.7.4 implies the boundedness of this family of rescalings in $W^{1,2}\left(B_{R}\right)$ for any $R>1$. Indeed, the rescalings above will be defined in $B_{R}$ if $t<1 /(2 R)$, and by Lemma 1.7.4, we will have

$$
\begin{aligned}
\int_{B_{R}}\left|\nabla u_{x_{0}, t}^{\phi}\right|^{2} & =\frac{\mathrm{e}^{\frac{3 b}{\alpha} t^{\alpha}}}{t^{n+1}} \int_{B_{R t}\left(x_{0}\right)}|\nabla u|^{2} \leq C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} R^{n+1} \\
\int_{\partial B_{R}}\left(u_{x_{0}, t}^{\phi}\right)^{2} & =\frac{\mathrm{e}^{\frac{3 b}{\alpha} t^{\alpha}}}{t^{n+2}} \int_{\partial B_{R t}\left(x_{0}\right)} u^{2} \leq C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}^{2} R^{n+2}
\end{aligned}
$$

for $0<t<t_{0} / R$. Arguing as in the proof of Proposition 1.6.1, we have for a sequence $t=t_{\mathrm{j}} \rightarrow 0+$

$$
u_{x_{0}, t_{\mathrm{j}}}^{\phi} \rightarrow u_{x_{0}, 0}^{\phi} \quad \text { in } \quad C_{\mathrm{loc}}^{1}\left(B_{R}^{ \pm} \cup B_{R}^{\prime}\right) .
$$

By letting $R \rightarrow \infty$ and using Cantor's diagonal argument, we therefore have that over a subsequence $t=t_{\mathrm{j}} \rightarrow 0+$

$$
u_{x_{0}, t_{\mathrm{j}}}^{\phi} \rightarrow u_{x_{0}, 0}^{\phi} \quad \text { in } \quad C_{\mathrm{loc}}^{1}\left(\mathbb{R}_{ \pm}^{n} \cup \mathbb{R}^{n-1}\right) .
$$

We call such $u_{x_{0}, 0}^{\phi}$ a 3/2-homogeneous blowup of $u$ at $x_{0}$. The name is explained by the fact that

$$
\lim _{t \rightarrow 0} \frac{\phi(t)}{t^{3 / 2}}=1
$$

which implies that if we consider the 3/2-homogeneous rescalings

$$
u_{x_{0}, t}^{(3 / 2)}(x)=\frac{u\left(t x+x_{0}\right)}{t^{3 / 2}},
$$

then we will have

$$
u_{x_{0}, 0}^{\phi}=\lim _{t_{\mathrm{j}} \rightarrow 0} u_{x_{0}, t_{\mathrm{j}}}^{\phi}=\lim _{t_{\mathrm{j}} \rightarrow 0} u_{x_{0}, t_{\mathrm{j}}}^{(3 / 2)}=: u_{x_{0}, 0}^{(3 / 2)}
$$

and thus $u_{x_{0}, 0}^{\phi}=u_{x_{0}, 0}^{(3 / 2)}$.
Remark 1.8.1. Because of the logarithmic term in the weak growth estimates in Lemma 1.7.1, at the moment we are unable to consider $\kappa$-homogeneous blowups as above for frequencies other than $\kappa=3 / 2$. However, once the logarithmic term is removed, the same construction as for $\kappa=3 / 2$ applies. In particular, we note that in Lemma 1.10.6 we prove the optimal growth estimates for frequencies $\kappa=2 m<\kappa_{0}, m \in \mathbb{N}$, enabling us to consider the $\kappa$-homogeneous blowups for these values of $\kappa$.

We show next that the $3 / 2$-homogeneous blowups are unique at free boundary points. This is achieved by the control on the "rotation" of the rescalings $u_{x_{0}, r}^{\phi}(x)$.

Lemma 1.8.2 (Rotation estimate). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}, x_{0} \in B_{1 / 2}^{\prime}$ a free boundary point, and $\delta$ as in Lemma 1.7.3. Then for $\kappa=3 / 2$ and $\phi=\phi_{3 / 2}$

$$
\int_{\partial B_{1}}\left|u_{x_{0}, t}^{\phi}-u_{x_{0}, s}^{\phi}\right| \leq C t^{\delta / 2}, \quad s<t<t_{0}=t_{0}(n, \alpha)
$$

for $C=C(n, \alpha)\|u\|_{W^{1,2}\left(B_{1}\right)}$.
Proof. The proof uses computations similar to the proof of Lemma 1.7.1 combined with the growth estimated for $W_{3 / 2}(t, u)$ in Lemma 1.7.3. We assume $x_{0}=0$, and have

$$
\begin{aligned}
\int_{\partial B_{1}}\left|u_{t}^{\phi}-u_{s}^{\phi}\right| & \leq \int_{\partial B_{1}} \int_{s}^{t}\left|\frac{d}{d r} u_{r}^{\phi}\right| d r=\int_{s}^{t} \int_{\partial B_{1}}\left|\frac{d}{d r} u_{r}^{\phi}\right| d r \\
& \leq C_{n} \int_{s}^{t}\left(\int_{\partial B_{1}}\left|\frac{d}{d r} u_{r}^{\phi}\right|^{2}\right)^{1 / 2} \\
& \leq C_{n}\left(\int_{s}^{t} \frac{1}{r} d r\right)^{1 / 2}\left(\int_{s}^{t} r \int_{\partial B_{1}}\left|\frac{d}{d r} u_{r}^{\phi}\right|^{2}\right)^{1 / 2} \\
& \leq C_{n} \mathrm{e}^{c t^{\alpha}}\left(\log \frac{t}{s}\right)^{1 / 2}\left(\int_{s}^{t} \frac{d}{d r} W_{3 / 2}(r, u) d r\right)^{1 / 2}, \quad c=\frac{3 b}{2 \alpha}-\frac{a}{2},
\end{aligned}
$$

where we have re-used the computation made in the proof of Lemma 1.7.1. Thus, we obtain

$$
\int_{\partial B_{1}}\left|u_{t}^{\phi}-u_{s}^{\phi}\right| \leq C(n, \alpha)\left(\log \frac{t}{s}\right)^{1 / 2}\left(W_{3 / 2}(t, u)-W_{3 / 2}(s, u)\right)^{1 / 2} \leq C\left(\log \frac{t}{s}\right)^{1 / 2} t^{\delta / 2}
$$

Then, using a dyadic argument as Lemma 1.7.4, we can conclude that

$$
\int_{\partial B_{1}}\left|u_{t}^{\phi}-u_{s}^{\phi}\right| \leq C t^{\delta / 2}, \quad s<t<t_{0}
$$

as required. Indeed, let $k=0,1,2, \ldots$ be such that $t / 2^{k+1} \leq s<t / 2^{k}$. Then

$$
\begin{aligned}
\int_{\partial B_{1}}\left|u_{t}^{\phi}-u_{s}^{\phi}\right| & \leq \sum_{\mathrm{j}=1}^{k} \int_{\partial B_{1}}\left|u_{t / 2^{\mathrm{j}-1}}^{\phi}-u_{t / 2^{\mathrm{j}}}^{\phi}\right|+\int_{\partial B_{1}}\left|u_{t / 2^{k}}^{\phi}-u_{s}^{\phi}\right| \\
& \leq C(\log 2)^{1 / 2} \sum_{\mathrm{j}=1}^{k+1}\left(t / 2^{\mathrm{j}-1}\right)^{\delta / 2} \leq C(\log 2)^{1 / 2} \frac{t^{\delta / 2}}{1-2^{-\delta / 2}}
\end{aligned}
$$

This completes the proof.
The uniqueness of $3 / 2$-homogeneous blowup now follows.
Lemma 1.8.3. Let $u_{x_{0}, 0}^{\phi}$ be a blowup at a free boundary point $x_{0} \in B_{1 / 2}^{\prime}$. Then for $\kappa=3 / 2$

$$
\int_{\partial B_{1}}\left|u_{x_{0}, t}^{\phi}-u_{x_{0}, 0}^{\phi}\right| \leq C t^{\delta / 2}, \quad 0<t<t_{0}
$$

where $C=C\left(n, \alpha,\|u\|_{W^{1,2}\left(B_{1}\right)}\right)$ and $\delta=\delta(n, \alpha)>0$ are as in Lemma 1.8.2. In particular, the blowup $u_{x_{0}, 0}^{\phi}$ is unique.

Proof. If $u_{x_{0}, 0}$ is the limit of $u_{x_{0}, t_{\mathrm{j}}}^{\phi}$ for $t_{\mathrm{j}} \rightarrow 0$, then first part of the lemma follows immediately from Lemma 1.8.2, by taking $s=t_{\mathrm{j}} \rightarrow 0$ and passing to the limit.

To see the uniqueness of blowup, we observe that $u_{x_{0}, 0}^{\phi}$ is a solution of the Signorini problem in $B_{1}$, by arguing as in the proof of Lemma 1.6.1 for Almgren blowups. Now, if $\tilde{u}_{x_{0}, 0}^{\phi}$ is another blowup, then from the first part of the lemma we will have

$$
\int_{\partial B_{1}}\left|\tilde{u}_{x_{0}, 0}^{\phi}-u_{x_{0}, 0}^{\phi}\right|^{2}=0
$$

implying that both $\tilde{u}_{x_{0}, 0}^{\phi}$ and $u_{x_{0}, 0}^{\phi}$ are solutions of the Signorini problem in $B_{1}$ with the same boundary values on $\partial B_{1}$. By the uniqueness of such solutions, we have $\tilde{u}_{x_{0}, 0}^{\phi}=u_{x_{0}, 0}^{\phi}$ in $B_{1}$. The equality propagates to all of $\mathbb{R}^{n}$ by the unique continuation of harmonic functions in $\mathbb{R}_{ \pm}^{n}$.

We next show that not only the blowups are unique, but also depend continuously on a free boundary point.

Lemma 1.8.4 (Continuous dependence of blowups). There exists $\rho=\rho(n, \alpha)>0$ such that if $x_{0}, y_{0} \in B_{\rho}$ are free boundary points, then

$$
\int_{\partial B_{1}}\left|u_{x_{0}, 0}^{\phi}-u_{y_{0}, 0}^{\phi}\right| \leq C\left|x_{0}-y_{0}\right|^{\gamma}
$$

with $C=C\left(n, \alpha,\|u\|_{W^{1,2}\left(B_{1}\right)}\right)$ and $\gamma=\gamma(n, \alpha)>0$.
Proof. Let $d=\left|x_{0}-y_{0}\right|$ and $d^{\mu} \leq r \leq 2 d^{\mu}$ with $\mu \in(0,1]$ to be determined. By Lemma 1.8.3 we have

$$
\begin{aligned}
\int_{\partial B_{1}}\left|u_{x_{0}, 0}^{\phi}-u_{y_{0}, 0}^{\phi}\right| & \leq 2 C r^{\delta / 2}+\int_{\partial B_{1}}\left|u_{x_{0}, r}^{\phi}-u_{y_{0}, r}^{\phi}\right| \\
& \leq C d^{\mu \delta / 2}+\frac{C}{d^{\mu(n+1 / 2)}} \int_{\partial B_{r}}\left|u\left(x_{0}+z\right)-u\left(y_{0}+z\right)\right| d S_{z}
\end{aligned}
$$

and taking the average over $d^{\mu} \leq r \leq 2 d^{\mu}$, we have

$$
\int_{\partial B_{1}}\left|u_{x_{0}, 0}^{\phi}-u_{y_{0}, 0}^{\phi}\right| \leq C d^{\mu \delta / 2}+\frac{C}{d^{\mu(n+3 / 2)}} \int_{B_{2 d^{\mu} \backslash B_{d \mu}}}\left|u\left(x_{0}+z\right)-u\left(y_{0}+z\right)\right| d z
$$

On the other hand, by using Lemma 1.7.4,

$$
\begin{aligned}
\int_{B_{2 d^{\mu}} \backslash B_{d \mu}}\left|u\left(x_{0}+z\right)-u\left(y_{0}+z\right)\right| d z & \leq \int_{B_{2 d^{\mu}} \backslash B_{d \mu}}\left|\int_{0}^{1} \frac{d}{d s} u\left(z+x_{0}(1-s)+y_{0} s\right) d s\right| d z \\
& \leq\left|x_{0}-y_{0}\right| \int_{0}^{1} \int_{B_{2 d^{\mu}}}\left|\nabla u\left(z+x_{0}(1-s)+y_{0} s\right)\right| d z d s \\
& \leq d \int_{0}^{1}\left(\int_{\left.B_{2 d^{\mu}\left(x_{0}(1-s)+y_{0} s\right)}|\nabla u|\right) d s}\right. \\
& \leq d \int_{B_{2 d^{\mu}+d}\left(x_{0}\right)}|\nabla u| \leq d \int_{B_{3 d^{\mu}}\left(x_{0}\right)}|\nabla u| \\
& \leq C d^{1+\mu n / 2}\left(\int_{B_{3 d^{\mu}\left(x_{0}\right)}}|\nabla u|^{2}\right)^{1 / 2} \leq C d^{1+\mu n / 2} d^{\mu(n+1) / 2} \\
& \leq C d^{1+\mu(n+1 / 2)}
\end{aligned}
$$

provided $3 d^{\mu}<t_{0}$, which will hold if $d<\rho(n, \alpha)$.

Combining the estimates, we infer that

$$
\int_{\partial B_{1}}\left|u_{x_{0}, 0}^{\phi}-u_{y_{0}, 0}^{\phi}\right| \leq C d^{\mu \delta / 2}+C d^{1-\mu} .
$$

Now choosing $\mu$ so that $\mu \delta / 2=1-\mu$, that is $\mu=1 /(1+\delta / 2)$, we obtain

$$
\int_{\partial B_{1}}\left|u_{x_{0}, 0}^{\phi}-u_{y_{0}, 0}^{\phi}\right| \leq C\left|x_{0}-y_{0}\right|^{\gamma}, \quad x_{0}, y_{0} \in B_{\rho}^{\prime}
$$

with

$$
\gamma=\frac{\delta}{\delta+2}
$$

### 1.9 Regularity of the regular set

In this section we establish one of the main result of this chapter, the $C^{1, \gamma}$ regularity of the regular set. In fact, the most technical part of the proof has already been done in the previous section, where we proved the uniqueness of the $3 / 2$-homogeneous blowups, as well as their Hölder continuous dependence on the free boundary points.

We start by defining the regular set.

Definition 1.9.1 (Regular points). For an almost minimizer u for the Signorini problem in $B_{1}$, we say that a free boundary point $x_{0}$ is regular if

$$
\widehat{N}\left(0+, u, x_{0}\right)=3 / 2
$$

Note that since $3 / 2<2 \leq \kappa_{0}$, we will have that $\widehat{N}(r)<\kappa_{0}$ for small $r>0$, implying that $\widetilde{N}(r)=\widehat{N}(r)$ for such $r$ and consequently that

$$
N(0+)=\widetilde{N}(0+)=\widehat{N}(0+)=3 / 2 .
$$

In particular, the condition above does not depend on the choice of $\kappa_{0} \geq 2$.
We denote the set of all regular points of $u$ by $\mathcal{R}(u)$ and call it the regular set.

An important ingredient in the analysis of the regular set is the following nondegeneracy lemma.

Lemma 1.9.1 (Nondegeneracy at regular points). Let $x_{0} \in B_{1 / 2}^{\prime} \cap \mathcal{R}(u)$ for an almost minimizer $u$ for the Signorini problem in $B_{1}$. Then, for $\kappa=3 / 2$,

$$
\liminf _{t \rightarrow 0} \int_{\partial B_{1}}\left(u_{x_{0}, t}^{\phi}\right)^{2}=\liminf _{t \rightarrow 0} \frac{1}{t^{n+2}} \int_{\partial B_{t}\left(x_{0}\right)} u^{2}>0 .
$$

Proof. As before, assume $x_{0}=0$. In terms of the quantities defined in the proofs of Lemmas 1.7.1 and 1.7.4, we want to prove that

$$
\liminf _{t \rightarrow 0} m(t)>0
$$

Assume, towards a contradiction, that $m\left(t_{\mathrm{j}}\right) \rightarrow 0$ for some sequence $t_{\mathrm{j}} \rightarrow 0$. Recall that by the proof of Lemma 1.7.4, we have

$$
|m(t)-m(s)| \leq C t^{\delta / 2}, \quad 0<s<t<t_{0}
$$

Now, setting $s=t_{\mathrm{j}} \rightarrow 0$, we conclude that

$$
|m(t)| \leq C t^{\delta / 2}, \quad 0<t<t_{0}
$$

Equivalently, we can rewrite this as

$$
\int_{\partial B_{t}} u^{2} \leq C t^{n+2+\delta}
$$

Next, take $\tilde{\kappa}=3 / 2+\delta / 4$ and consider Weiss's monotonicity formula

$$
W_{\tilde{\kappa}}(t, u)=\frac{\mathrm{e}^{a_{\tilde{\kappa}} t^{\alpha}}}{t^{n+2 \tilde{\kappa}-2}}\left[\int_{B_{t}}|\nabla u|^{2}-\tilde{\kappa} \frac{1-b t^{\alpha}}{t} \int_{\partial B_{t}} u^{2}\right] .
$$

Now observe that

$$
\frac{1}{t^{n+2 \tilde{\kappa}-1}} \int_{\partial B_{t}} u^{2} \leq C t^{\delta / 2} \rightarrow 0
$$

which readily implies that

$$
W_{\tilde{\kappa}}(0+, u) \geq 0 .
$$

In particular, by monotonicity, $W_{\tilde{\kappa}}(t, u) \geq 0$, for small $t>0$, which also implies that $\widetilde{N}(t, u) \geq \tilde{\kappa}$. But then $N(0+, u)=\widetilde{N}(0+, u) \geq \tilde{\kappa}=3 / 2+\delta / 4$ contrary to the assumption in the lemma. This completes the proof.

The next result provides two important facts: a gap in possible values of Almgren's frequency $N(0+)$ as well as the classification of $3 / 2$-homogeneous blowups.

Proposition 1.9.1. If $\widehat{N}\left(0+, u, x_{0}\right)=\kappa<2$, then $\kappa=3 / 2$ and

$$
u_{x_{0}, 0}^{\phi}(x)=a_{x_{0}} \operatorname{Re}\left(x^{\prime} \cdot \nu_{x_{0}}+i\left|x_{n}\right|\right)^{3 / 2}
$$

for some $a_{x_{0}}>0, \nu_{x_{0}} \in \partial B_{1}^{\prime}$.
Proof. Without loss of generality, we may assume $x_{0}=0$. Let $r_{\mathrm{j}} \rightarrow 0+$ be a sequence such that $u_{r_{\mathrm{j}}}^{\phi} \rightarrow u_{0}^{\phi}$ in $C_{\text {loc }}^{1}\left(\mathbb{R}_{ \pm}^{n} \cup \mathbb{R}^{n-1}\right)$. Comparing 3/2-almost homogeneous and Almgren rescalings, we have

$$
u_{r}^{\phi}(x)=u_{r}^{A}(x) \mu(r), \quad \mu(r):=\frac{\left(\frac{1}{r^{n-1}} \int_{\partial B_{r}} u^{2}\right)^{1 / 2}}{\phi(r)}
$$

By the optimal growth estimate (Lemma 1.7.4) and the nondegeneracy at regular points (Lemma 1.9.1) we have

$$
0<\liminf _{r \rightarrow 0+} \mu(r) \leq \limsup _{r \rightarrow 0+} \mu(r)<\infty
$$

Thus, we may assume that, over a subsequence of $r_{\mathrm{j}}, \mu\left(r_{\mathrm{j}}\right) \rightarrow \mu_{0} \in(0, \infty)$, and therefore

$$
u_{r_{\mathrm{j}}}^{\phi} \rightarrow \mu_{0} u_{0}^{A} \quad \text { in } C_{\mathrm{loc}}^{1}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)
$$

where $u_{0}^{A}$ is an Almgren blowup of $u$ at $x_{0}=0$. Now, since $\kappa<\kappa_{0}$, we can apply Proposition 1.6.1 to obtain that $u_{0}^{A}$ is a nonzero $\kappa$-homogeneous solution of the Signorini problem
in $B_{1}$, even in $x_{n}$-variable. Next, applying Lemma 1.6.1, we have $3 / 2 \leq \kappa<2$ and thus by Proposition 9.9 in [48], we must have $\kappa=3 / 2$ and

$$
u_{0}^{A}(x)=C_{n} \operatorname{Re}\left(x^{\prime} \cdot \nu_{0}+i\left|x_{n}\right|\right)^{3 / 2}
$$

for some $C_{n}>0, \nu_{0} \in \partial B_{1}$. (The constant $C_{n}$ comes from the normalization $\int_{\partial B_{1}}\left(u_{0}^{A}\right)^{2}=1$.) Thus,

$$
u_{0}^{\phi}(x)=a_{0} \operatorname{Re}\left(x^{\prime} \cdot \nu_{0}+i\left|x_{n}\right|\right)^{3 / 2} \quad \text { in } B_{1}
$$

with $a_{0}=C_{n} \mu_{0}$. By the unique continuation of harmonic functions in $\mathbb{R}_{ \pm}^{n}$, we obtain that the above formula for $u_{0}^{\phi}$ propagates to all of $\mathbb{R}^{n}$.

Proposition 1.9.1 has an immediate corollary.
Corollary 1.9.2 (Almgren frequency gap). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$ and $x_{0}$ a free boundary point. Then either

$$
\widehat{N}(0+, u)=3 / 2 \quad \text { or } \quad \widehat{N}(0+, u) \geq 2
$$

Yet another important fact is as follows.

Corollary 1.9.3. The regular set $\mathcal{R}(u)$ is a relatively open subset of the free boundary.
Proof. For a fixed $0<t<t_{0}$, the mapping $x \mapsto \widehat{N}(t, u, x)$ is continuous on $\Gamma(u)$. Then, by the monotonicity of $\widehat{N}$, the mapping $x \mapsto \widehat{N}\left(0+, u, x_{0}\right)$ is upper semicontinuous on $\Gamma(u)$. Moreover, by Proposition 1.9.1,

$$
\mathcal{R}(u)=\{x \in \Gamma(u): \widehat{N}(0+, u, x)<2\}
$$

which implies that $\mathcal{R}(u)$ is relatively open in $\Gamma(u)$.
The combination of Proposition 1.9.1 and Lemma 1.8.4 implies the following lemma.
Lemma 1.9.4. Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$, and $x_{0} \in$ $\mathcal{R}(u)$. Then there exists $\rho>0$, depending on $x_{0}$ such that $B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u) \subset \mathcal{R}(u)$ and if
$u_{\bar{x}, 0}^{\phi}(x)=a_{\bar{x}} \operatorname{Re}\left(x^{\prime} \cdot \nu_{\bar{x}}+i\left|x_{n}\right|\right)^{3 / 2}$ is the unique 3/2-homogeneous blowup at $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u)$, then

$$
\begin{aligned}
& \left|a_{\bar{x}}-a_{\bar{y}}\right| \leq C_{0}|\bar{x}-\bar{y}|^{\gamma}, \\
& \left|\nu_{\bar{x}}-\nu_{\bar{y}}\right| \leq C_{0}|\bar{x}-\bar{y}|^{\gamma},
\end{aligned}
$$

for any $\bar{x}, \bar{y} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u)$ with a constant $C_{0}$ depending on $x_{0}$.
Proof. The proof follows by repeating the argument in Lemma 7.5 in [20].
Now we are ready to prove the main result on the regularity of the regular set.
Theorem 1.9.5 ( $C^{1, \gamma}$ regularity of the regular set). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$. Then, if $x_{0} \in B_{1 / 2}^{\prime} \cap \mathcal{R}(u)$, there exists $\rho>0$, depending on $x_{0}$ such that, after a possible rotation of coordinate axes in $\mathbb{R}^{n-1}$, one has $B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u) \subset \mathcal{R}(u)$, and

$$
B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u)=B_{\rho}^{\prime}\left(x_{0}\right) \cap\left\{x_{n-1}=g\left(x_{1}, \ldots, x_{n-2}\right)\right\},
$$

for $g \in C^{1, \gamma}\left(\mathbb{R}^{n-2}\right)$ with an exponent $\gamma=\gamma(n, \alpha) \in(0,1)$.

Proof. The proof of the theorem is similar to that of Theorem 1.2 in [20]. However, we provide full details since there are technical differences.

Step 1. By relative openness of $\mathcal{R}(u)$ in $\Gamma(u)$, for small $\rho>0$ we have $B_{2 \rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u) \subset \mathcal{R}(u)$. We then claim that for any $\varepsilon>0$, there is $r_{\varepsilon}>0$ such that for $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u), r<r_{\varepsilon}$, we have that for $\phi=\phi_{3 / 2}$

$$
\left\|u_{\bar{x}, r}^{\phi}-u_{\bar{x}, 0}^{\phi}\right\|_{C^{1}\left(\overline{B_{1}^{ \pm}}\right)}<\varepsilon .
$$

Assuming the contrary, there is a sequence of points $\bar{x}_{\mathrm{j}} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u)$ and radii $r_{\mathrm{j}} \rightarrow 0$ such that

$$
\left\|u_{\bar{x}_{j}, r_{j}}^{\phi}-u_{\overline{x_{j}}, 0}^{\phi}\right\|_{C^{1}\left(\overline{B_{1}^{ \pm}}\right)} \geq \varepsilon_{0}
$$

for some $\varepsilon_{0}>0$. Taking a subsequence, if necessary, we may assume $\bar{x}_{\mathrm{j}} \rightarrow \bar{x}_{0} \in \overline{B_{\rho}^{\prime}\left(x_{0}\right)} \cap \Gamma(u)$. Using estimates (1.3.1), (1.4.2) and Lemma 1.7.4, we can see that $u_{\bar{x}_{\mathrm{j}}, r_{\mathrm{j}}}^{\phi}$ are uniformly bounded in $C^{1, \beta}\left(B_{2}^{ \pm} \cup B_{2}^{\prime}\right)$. Thus, we may assume that for some $w$

$$
u_{\overline{x_{j}}, r_{\mathrm{j}}}^{\phi} \rightarrow w \quad \text { in } C^{1}\left(\overline{B_{1}^{ \pm}}\right) .
$$

By arguing as in the proof of Proposition 1.6.1, we see that the limit $w$ is a solution of the Signorini problem in $B_{1}$. Further, by Lemma 1.8.3, we have

$$
\left\|u_{\bar{x}_{\mathrm{j}}, r_{\mathrm{j}}}^{\phi}-u_{\bar{x}_{\mathrm{j}}, 0}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)} \rightarrow 0 .
$$

On the other hand, by Lemma 1.9.4, we have

$$
u_{\bar{x}_{j}, 0}^{\phi} \rightarrow u_{\bar{x}_{0}, 0}^{\phi} \quad \text { in } C^{1}\left(\overline{B_{1}^{ \pm}}\right)
$$

and thus

$$
w=u_{\bar{x}_{0}, 0}^{\phi} \quad \text { on } \partial B_{1} .
$$

Since both $w$ and $u_{\bar{x}_{0}, 0}^{\phi}$ are solutions of the Signorini problem, they must coincide also in $B_{1}$. Therefore

$$
u_{\bar{x}_{j}, r_{j}}^{\phi} \rightarrow u_{\bar{x}_{0}, 0}^{\phi} \quad \text { in } C^{1}\left(\overline{B_{1}^{ \pm}}\right),
$$

implying also that

$$
\left\|u_{\bar{x}_{\mathrm{j}}, r_{\mathrm{j}}}^{\phi}-u_{\bar{x}_{\mathrm{j}}, 0}^{\phi}\right\|_{C^{1}\left(\overline{B_{1}^{ \pm}}\right)} \rightarrow 0
$$

which contradicts our assumption.
Step 2. As [20], for a given $\varepsilon>0$ and a unit vector $\nu \in \mathbb{R}^{n-1}$ define the cone

$$
\mathcal{C}_{\varepsilon}(\nu)=\left\{x^{\prime} \in \mathbb{R}^{n-1}: x^{\prime} \cdot \nu>\varepsilon\left|x^{\prime}\right|\right\} .
$$

By Lemma 1.9.4, we may assume $a_{\bar{x}} \geq \frac{a_{x_{0}}}{2}$ for $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u)$ by taking $\rho$ small. For such $\rho$ we then claim that for any $\varepsilon>0$ there is $r_{\varepsilon}>0$ such that for any $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u)$ we have

$$
\bar{x}+\left(\mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) \subset\{u(\cdot, 0)>0\} .
$$

Indeed, denoting $\mathcal{K}_{\varepsilon}(\nu)=\mathcal{C}_{\varepsilon} \cap \partial B_{1 / 2}^{\prime}$, we have for some universal $C_{\varepsilon}>0$

$$
\mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right) \Subset\left\{u_{\bar{x}, 0}^{\phi}(\cdot, 0)>0\right\} \cap B_{1}^{\prime} \quad \text { and } \quad u_{\bar{x}, 0}^{\phi}(\cdot, 0) \geq a_{\bar{x}} C_{\varepsilon} \geq \frac{a_{x_{0}}}{2} C_{\varepsilon} \quad \text { on } \mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right) .
$$

Since $\frac{a_{x_{0}}}{2} C_{\varepsilon}$ is independent of $\bar{x}$, by Step 1 we can find $r_{\varepsilon}>0$ such that for $r<2 r_{\varepsilon}$,

$$
u_{\bar{x}, r}^{\phi}(\cdot, 0)>0 \quad \text { on } \mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right) .
$$

This implies that for $r<2 r_{\varepsilon}$,

$$
u(\cdot, 0)>0 \quad \text { on } \quad \bar{x}+r \mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right)=\bar{x}+\left(\mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap \partial B_{r / 2}^{\prime}\right) .
$$

Taking the union over all $r<2 r_{\varepsilon}$, we obtain

$$
u(\cdot, 0)>0 \quad \text { on } \bar{x}+\left(\mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) .
$$

Step 3. We claim that for given $\varepsilon>0$, there exists $r_{\varepsilon}>0$ such that for any $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u)$ we have $\bar{x}-\left(\mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) \subset\{u(\cdot, 0)=0\}$.

Indeed, we first note that

$$
-\partial_{x_{n}}^{+} u_{\bar{x}, 0}^{\phi} \geq a_{\bar{x}} C_{\varepsilon}>\left(\frac{a_{x_{0}}}{2}\right) C_{\varepsilon} \quad \text { on }-\mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right)
$$

for a universal constant $C_{\varepsilon}>0$. From Step 1, there exists $r_{\varepsilon}>0$ such that for $r<2 r_{\varepsilon}$,

$$
-\partial_{x_{n}}^{+} u_{\bar{x}, r}^{\phi}(\cdot, 0)>0 \quad \text { on }-\mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right)
$$

By arguing as in Step 2, we obtain

$$
-\partial_{x_{n}}^{+} u(\cdot, 0)>0 \quad \text { on } \bar{x}-\left(\mathcal{C}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) .
$$

By the complementarity condition in Lemma 1.4.2, we therefore conclude that

$$
\bar{x}-\left(\mathcal{C}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) \subset\left\{-\partial_{x_{n}}^{+} u(\cdot, 0)>0\right\} \subset\{u(\cdot, 0)=0\} .
$$

Step 4. By rotation in $\mathbb{R}^{n-1}$ we may assume $\nu_{x_{0}}=\mathrm{e}_{n-1}$. For any $\varepsilon>0$, by Lemma 1.9.4 again, we can take $\rho_{\varepsilon}=\rho\left(x_{0}, \varepsilon\right)$, possibly smaller than $\rho$ in the previous steps, such that

$$
\mathcal{C}_{2 \varepsilon}\left(\mathrm{e}_{n-1}\right) \cap B_{r_{\varepsilon}}^{\prime} \subset \mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime} \text { for } \bar{x} \in B_{\rho_{\varepsilon}}^{\prime}\left(x_{0}\right) \cap \Gamma(u) .
$$

By Step 2 and Step 3, for $\bar{x} \in B_{\rho_{\varepsilon}}^{\prime}\left(x_{0}\right) \cap \Gamma(u)$,

$$
\begin{aligned}
& \bar{x}+\left(\mathcal{C}_{2 \varepsilon}\left(\mathrm{e}_{n-1}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) \subset\{u(\cdot, 0)>0\} \\
& \bar{x}-\left(\mathcal{C}_{2 \varepsilon}\left(\mathrm{e}_{n-1}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) \subset\{u(\cdot, 0)=0\}
\end{aligned}
$$

Now, fixing $\varepsilon=\varepsilon_{0}$, by the standard arguments, we conclude that there exists a Lipschitz function $g: \mathbb{R}^{n-2} \rightarrow \mathbb{R}$ with $|\nabla g| \leq C_{n} / \varepsilon_{0}$ such that

$$
\begin{aligned}
& B_{\rho_{e_{0}}}^{\prime}\left(x_{0}\right) \cap\{u(\cdot, 0)=0\}=B_{\rho_{e_{0}}}^{\prime}\left(x_{0}\right) \cap\left\{x_{n-1} \leq g\left(x^{\prime \prime}\right)\right\}, \\
& B_{\rho_{e_{0}}}^{\prime}\left(x_{0}\right) \cap\{u(\cdot, 0)>0\}=B_{\rho_{e_{0}}}^{\prime}\left(x_{0}\right) \cap\left\{x_{n-1}>g\left(x^{\prime \prime}\right)\right\} .
\end{aligned}
$$

Step 5. Taking $\varepsilon \rightarrow 0$ in Step $4, \Gamma(u)$ is differentiable at $x_{0}$ with normal $\nu_{x_{0}}$. Recentering at any $\bar{x} \in B_{\rho_{\varepsilon_{0}}}^{\prime}\left(x_{0}\right) \cap \Gamma(u)$, we see that $\Gamma(u)$ has a normal $\nu_{\bar{x}}$ at $\bar{x}$. By Lemma 1.9.4, we conclude that $g$ in Step 4 is $C^{1, \gamma}$. This completes the proof of the theorem.

### 1.10 Singular points

In this section we study the set of so-called singular free boundary points. An important technical tool to accomplish this is the logarithmic epiperimetric inequality of [27]. We use it for two purposes: to establish the optimal growth at singular points as well as the rate of convergence of rescalings to blowups, ultimately implying a structural theorem for the singular set.

Definition 1.10.1 (Singular points). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$. We say that a free boundary point $x_{0}$ is singular if the coincidence set $\Lambda(u)=\{u(\cdot, 0)=$ $0\} \subset B_{1}^{\prime}$ has zero $H^{n-1}$-density at $x_{0}$, i.e.,

$$
\lim _{r \rightarrow 0+} \frac{H^{n-1}\left(\Lambda(u) \cap B_{r}^{\prime}\left(x_{0}\right)\right)}{H^{n-1}\left(B_{r}^{\prime}\left(x_{0}\right)\right)}=0 .
$$

By using Almgren's rescalings $u_{x_{0}, r}^{A}$, we can rewrite this condition as

$$
\lim _{r \rightarrow 0+} H^{n-1}\left(\Lambda\left(u_{x_{0}, r}^{A}\right) \cap B_{1}^{\prime}\right)=0
$$

We denote the set of all singular points by $\Sigma(u)$ and call it the singular set.
Throughout the section we will assume that

$$
\kappa_{0}>2
$$

We can take $\kappa_{0}$ as large as we like, however, we have to remember that the constants in $\widehat{N}=\widehat{N}_{\kappa_{0}}$ and $W_{\kappa}$ do depend on $\kappa_{0}$.

We then have the following characterization of singular points, similar to Proposition 9.22 in [48] for the solutions of the Signorini problem.

Proposition 1.10.1 (Characterization of singular points). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$, and $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(u)$ be such that $\widehat{N}\left(0+, u, x_{0}\right)=\kappa<\kappa_{0}$. Then the following statements are equivalent.
(i) $x_{0} \in \Sigma(u)$.
(ii) any Almgren blowup of $u$ at $x_{0}$ is a nonzero polynomial from the class

$$
\mathcal{Q}_{\kappa}=\{q: q \text { is homogeneous polynomial of degree } \kappa \text { such that }
$$

$$
\left.\Delta q=0, q\left(x^{\prime}, 0\right) \geq 0, q\left(x^{\prime}, x_{n}\right)=q\left(x^{\prime},-x_{n}\right)\right\}
$$

(iii) $\kappa=2 m$ for some $m \in \mathbb{N}$.

Note that for $\kappa<\kappa_{0}$, the condition $\widehat{N}(0+)=\kappa$ is equivalent to $N(0+)=\kappa$.

Proof. Without loss of generality we may assume $x_{0}=0$. By Proposition 1.6.1, any Almgren blowup $u_{0}^{A}$ of $u$ at 0 is a nonzero global solution of the Signorini problem, homogeneous of degree $\kappa$. Moreover $u_{0}^{A}$ is a $C_{\text {loc }}^{1}$ limit of Almgren rescalings $u_{r_{\mathrm{j}}}^{A}$ in $B_{1}^{ \pm} \cup B_{1}^{\prime}$. Because of that, most parts of the proof of this proposition are just the repetitions of Proposition 9.22 in [48]. Thus, by following Proposition 9.22 in [48], we can easily see the implications (ii) $\Rightarrow$ (iii), (iii) $\Rightarrow$ (ii), (ii) $\Rightarrow$ (i). Moreover, in the proof of the remaining implication (i) $\Rightarrow$ (ii), the only nontrivial part is that any blowup $u_{0}^{A}$ is harmonic in $B_{1}$. But this comes from the complementarity condition in Lemma 1.4.2. Indeed, assuming (i), we claim that

$$
\partial_{x_{n}}^{+} u_{0}^{A}=0 \quad \text { in } \quad B_{1}^{\prime}
$$

Otherwise,

$$
H^{n-1}\left(\left\{-\partial_{x_{n}}^{+} u_{0}^{A}(\cdot, 0)>0\right\} \cap B_{1}^{\prime}\right) \geq \delta
$$

for some $\delta>0$. Then using the continuity from the below we also have that for some $\rho>0$,

$$
H^{n-1}\left(\left\{-\partial_{x_{n}}^{+} u_{0}^{A}(\cdot, 0)>\rho\right\} \cap B_{1-\rho}^{\prime}\right) \geq \delta / 2
$$

Using $C_{\text {loc }}^{1}$ convergence $u_{r_{\mathrm{j}}}^{A} \rightarrow u_{0}^{A}$ in $B_{1}^{ \pm} \cup B_{1}^{\prime}$ and applying the complementarity condition in Lemma 1.4.2 to rescalings $u_{r_{\mathrm{j}}}^{A}$, we obtain that for small $r_{\mathrm{j}}$,

$$
H^{n-1}\left(\Lambda\left(u_{r_{\mathrm{j}}}^{A}\right) \cap B_{1}^{\prime}\right) \geq H^{n-1}\left(\left\{-\partial_{x_{n}}^{+} u_{r_{\mathrm{j}}}^{A}(\cdot, 0)>0\right\} \cap B_{1}^{\prime}\right) \geq \delta / 4
$$

which contradicts (i). Now recalling that $u_{0}^{A}$ is a solution of the Signorini problem, even in $x_{n}$-variable, it satisfies

$$
\Delta u_{0}^{A}=\left.2\left(\partial_{x_{n}}^{+} u_{0}^{A}\right) H^{n-1}\right|_{\Lambda\left(u_{0}^{A}\right)}=0 \quad \text { in } \quad B_{1} .
$$

By homogeneity, we obtain that $u_{0}^{A}$ is harmonic in all of $\mathbb{R}^{n}$, and we complete the proof as in [48].

In order to study the singular set, in view of Proposition 1.10.1, we need to refine the growth estimate in Lemma 1.7 .1 by removing the logarithmic term in the case when $\kappa=$ $2 m<\kappa_{0}, m \in \mathbb{N}$. In the case $\kappa=3 / 2$ we were able to do so by proving a decay estimate for $W_{3 / 2}$ with the help of the epiperimetric inequality. In the case $\kappa=2 m$ we will use the so-called logarithmic epiperimetric inequality for the Weiss energy

$$
W_{\kappa}^{0}(w)=\int_{B_{1}}|\nabla w|^{2}-\kappa \int_{\partial B_{1}} w^{2}, \quad \kappa=2 m, m \in \mathbb{N}
$$

that first appeared in [27]. To state this result, we recall the notation

$$
\mathcal{A}=\left\{w \in W^{1,2}\left(B_{1}\right): w \geq 0 \text { on } B_{1}^{\prime}, w\left(x^{\prime}, x_{n}\right)=w\left(x^{\prime},-x_{n}\right)\right\} .
$$

Theorem 1.10.1 (Logarithmic epiperimetric inequality). Let $\kappa=2 m, m \in \mathbb{N}$ and $w \in \mathcal{A}$ be homogeneous of degree $\kappa$ in $B_{1}$ such that $w \in W^{1,2}\left(\partial B_{1}\right)$ and

$$
\int_{\partial B_{1}} w^{2} \leq 1, \quad\left|W_{\kappa}^{0}(w)\right| \leq 1
$$

There is constant $\varepsilon=\varepsilon(n, \kappa)>0$ and a function $v \in \mathcal{A}$ with $v=w$ on $\partial B_{1}$ such that

$$
W_{\kappa}^{0}(v) \leq W_{\kappa}^{0}(w)\left(1-\varepsilon\left|W_{\kappa}^{0}(w)\right|^{\gamma}\right), \quad \text { where } \gamma=\frac{n-2}{n} .
$$

To simplify the notations, in the results below all constants will depend on $n, \alpha, \kappa, \kappa_{0}$, as well as $\|u\|_{W^{1,2}\left(B_{1}\right)}$, unless stated otherwise, in addition to other quantities. Thus, when we write $C=C(\sigma)$, we mean $C=C\left(n, \alpha, \kappa, \kappa_{0},\|u\|_{W^{1,2}\left(B_{1}\right)}, \sigma\right)$.

The next lemma allows to apply the logarithmic epiperimetric inequality, without the constraints.

Lemma 1.10.2. Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$ such that $0 \in \Gamma(u)$ and $\widehat{N}(0+, u)=\kappa<\kappa_{0}, \kappa=2 m, m \in \mathbb{N}$. For $0<r<1$, let

$$
u_{r}(x)=u_{r}^{(\kappa)}(x)=\frac{u(r x)}{r^{\kappa}}, \quad w_{r}(x)=|x|^{\kappa} u_{r}\left(\frac{x}{|x|}\right) .
$$

Suppose that for a given $0 \leq \sigma \leq 1$, there is $C=C(\sigma)$ such that

$$
\int_{\partial B_{r}} u^{2} \leq C\left(\log \frac{1}{r}\right)^{\sigma} r^{n+2 \kappa-1}
$$

Then there is a constant $\varepsilon=\varepsilon(\sigma)>0$ and $h \in \mathcal{A}$ with $h=w_{r}$ on $\partial B_{1}$ such that
(i) If $\left|W_{\kappa}^{0}\left(w_{r}\right)\right| \geq \int_{\partial B_{1}} w_{r}^{2}$, then

$$
W_{\kappa}^{0}(h) \leq(1-\varepsilon) W_{\kappa}^{0}\left(w_{r}\right)
$$

(ii) If $\left|W_{\kappa}^{0}\left(w_{r}\right)\right| \leq 2 \int_{\partial B_{1}} w_{r}^{2}$, then

$$
W_{\kappa}^{0}(h) \leq W_{\kappa}^{0}\left(w_{r}\right)\left(1-\varepsilon\left(\log \frac{1}{r}\right)^{-\sigma \gamma}\left|W_{\kappa}^{0}\left(w_{r}\right)\right|^{\gamma}\right), \quad \text { where } \gamma=\frac{n-2}{n} .
$$

Proof. Let $A=\int_{\partial B_{1}} w_{r}^{2}+\left|W_{\kappa}^{0}\left(w_{r}\right)\right|$. Then by Theorem 1.10 .1 applied to $w_{r} / A^{1 / 2}$, there is $h \in \mathcal{A}$ such that $h=w_{r}$ on $\partial B_{1}$ and

$$
W_{\kappa}^{0}(h) \leq W_{\kappa}\left(w_{r}\right)^{0}\left(1-\varepsilon A^{-\gamma}\left|W_{\kappa}^{0}\left(w_{r}\right)\right|^{\gamma}\right)
$$

If $\left|W_{\kappa}^{0}\left(w_{r}\right)\right| \geq \int_{\partial B_{1}} w_{r}^{2}$, then $A \leq 2\left|W_{\kappa}^{0}\left(w_{r}\right)\right|$, implying

$$
W_{\kappa}^{0}(h) \leq W_{\kappa}^{0}\left(w_{r}\right)\left(1-\varepsilon 2^{-\gamma}\right) .
$$

If $\left|W_{\kappa}^{0}\left(w_{r}\right)\right| \leq 2 \int_{\partial B_{1}} w_{r}^{2}$, then

$$
A \leq 3 \int_{\partial B_{1}} w_{r}^{2}=\frac{3}{r^{n+2 \kappa-1}} \int_{\partial B_{r}} u^{2} \leq C(\sigma)\left(\log \frac{1}{r}\right)^{\sigma} .
$$

This completes the proof.

Now we show that the logarithmic epiperimetric inequality, combined with a growth estimate for $u$, implies a growth estimate on $W_{\kappa}(t, u)$. This is the first part of a bootstrapping argument that gradually decreases the power of $\log (1 / t)$ in the bound for $u$.

Lemma 1.10.3. Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$ such that $0 \in \Gamma(u)$ and $\widehat{N}(0+, u)=\kappa<\kappa_{0}, \kappa=2 m, m \in \mathbb{N}$. Suppose that for some $0 \leq \sigma \leq 1$

$$
\int_{\partial B_{r}} u^{2} \leq C(\sigma)\left(\log \frac{1}{r}\right)^{\sigma} r^{n+2 \kappa-1}, \quad 0<r<r_{0}(\sigma) .
$$

Then,

$$
0 \leq W_{\kappa}(t, u) \leq C(\sigma)\left(\log \frac{1}{t}\right)^{-\frac{1-\sigma \gamma}{\gamma}}, \quad 0<t<t_{0}(\sigma)
$$

Proof. We first observe that $W_{\kappa}(t, u) \geq 0$ for $0<t<t_{0}$, which follows easily from the condition $\widehat{N}(0+, u)=\kappa<\kappa_{0}$, see the beginning of the proof of Lemma 1.7.1.

Next, recall that in the proof of Lemma 1.7.3, we have used epiperimetric inequality to show that $0 \leq W_{3 / 2}(t, u) \leq C t^{\delta}$. This followed by obtaining a differential inequality for $W_{3 / 2}$. Thus, if for $0<t<t_{0}$, if alternative (i) holds in Lemma 1.10.2, i.e., $W_{\kappa}^{0}(h) \leq(1-\varepsilon) W_{\kappa}^{0}\left(w_{t}\right)$, by arguing in the same way, we can show that

$$
\begin{equation*}
\frac{d}{d t} W_{\kappa}(t, u) \geq \frac{\varepsilon / 4}{t} W_{\kappa}(t, u)-C t^{\alpha / 2-1} \tag{1.10.1}
\end{equation*}
$$

for $C=C(\sigma)$.
Suppose now the alternative (ii) holds in Lemma 1.10.2 for some $0<t<t_{0}$. Then, following the computations in Lemma 1.7.3, we have

$$
\frac{d}{d t} W_{\kappa}(t, u) \geq-\frac{(n+2 \kappa-2)\left(1-t^{\alpha}\right)}{t} W_{\kappa}(t, u)
$$

$$
\begin{aligned}
& +\frac{\mathrm{e}^{a t^{\alpha}}\left(1-b t^{\alpha}\right)}{t} \int_{\partial B_{1}}\left(\partial_{\nu} u_{t}-\kappa u_{t}\right)^{2}+\left(\partial_{\tau} u_{t}\right)^{2}-\kappa(n+\kappa-2) u_{t}^{2} \\
& +\left(2 \kappa_{0}+n\right) t^{\alpha-1} \int_{\partial B_{1}} u_{t}^{2}
\end{aligned}
$$

For $w_{t}$ as in the statement of Lemma 1.10.2, by following the computations in the proof of Theorem 1.5.1, we have the identity

$$
\int_{\partial B_{1}}\left(\partial_{\tau} u_{t}\right)^{2}-\kappa(n+\kappa-2) u_{t}^{2}=(n+2 \kappa-2) W_{\kappa}^{0}\left(w_{t}\right) .
$$

This gives

$$
\begin{align*}
& \frac{d}{d t} W_{\kappa}(t, u) \geq-\frac{(n+2 \kappa-2)\left(1-t^{\alpha}\right)}{t} W_{\kappa}(t, u) \\
& \quad+\frac{\mathrm{e}^{a t^{\alpha}}\left(1-b t^{\alpha}\right)}{t}(n+2 \kappa-2) W_{\kappa}^{0}\left(w_{t}\right)+\left(2 \kappa_{0}+n\right) t^{\alpha-1} \int_{\partial B_{1}} u_{t}^{2} . \tag{1.10.2}
\end{align*}
$$

Let now $v_{t}$ be the solution of the Signorini problem in $B_{1}$ with $v_{t}=u_{t}=w_{t}$ on $\partial B_{1}$. Then

$$
\begin{align*}
\left(1+t^{\alpha}\right) W_{\kappa}^{0}\left(w_{t}\right) & \geq\left(1+t^{\alpha}\right) W_{\kappa}^{0}\left(v_{t}\right) \geq \int_{B_{1}}\left|\nabla u_{t}\right|^{2}-\kappa\left(1+t^{\alpha}\right) \int_{\partial B_{1}} u_{t}^{2}  \tag{1.10.3}\\
& =W_{\kappa}^{0}\left(u_{t}\right)-\kappa t^{\alpha} \int_{\partial B_{1}} u_{t}^{2}=\mathrm{e}^{-a t^{\alpha}} W_{\kappa}(t, u)-\kappa(b+1) t^{\alpha} \int_{\partial B_{1}} u_{t}^{2}
\end{align*}
$$

Now, if

$$
\mathrm{e}^{-a t^{\alpha}} W_{\kappa}(t, u)-\kappa(b+1) t^{\alpha} \int_{\partial B_{1}} u_{t}^{2} \leq 0
$$

then by Lemma 1.7.1 we have

$$
\begin{equation*}
W_{\kappa}(t, u) \leq \mathrm{e}^{a t^{\alpha}} \kappa(b+1) t^{\alpha} \int_{\partial B_{1}} u_{t}^{2} \leq C t^{\alpha}\left(\log \frac{1}{t}\right) \leq C t^{\alpha / 2} \tag{1.10.4}
\end{equation*}
$$

We then proceed under the assumption

$$
\mathrm{e}^{-a t^{\alpha}} W_{\kappa}(t, u)-\kappa(b+1) t^{\alpha} \int_{\partial B_{1}} u_{t}^{2}>0
$$

which also implies

$$
W_{\kappa}^{0}\left(w_{t}\right)>0 .
$$

Now, applying Lemma 1.10.2, we have

$$
\begin{align*}
& W_{\kappa}^{0}\left(w_{t}\right) \geq W_{\kappa}^{0}\left(v_{t}\right)+\varepsilon\left(\log \frac{1}{t}\right)^{-\sigma \gamma} W_{\kappa}^{0}\left(w_{t}\right)^{\gamma+1} \\
& \geq \frac{1}{1+t^{\alpha}} {\left[\mathrm{e}^{-a t^{\alpha}} W_{\kappa}(t, u)-\kappa(b+1) t^{\alpha} \int_{\partial B_{1}} u_{t}^{2}\right] } \\
&+ \varepsilon\left(\log \frac{1}{t}\right)^{-\sigma \gamma}\left(\frac{1}{1+t^{\alpha}}\right)^{\gamma+1} \times \\
& \times\left[\mathrm{e}^{-a t^{\alpha}} W_{\kappa}(t, u)-\kappa(b+1) t^{\alpha} \int_{\partial B_{1}} u_{t}^{2}\right]^{\gamma+1} \\
& \geq\left(1-t^{\alpha}\right) {\left[\mathrm{e}^{-a t^{\alpha}} W_{\kappa}(t, u)-\kappa(b+1) t^{\alpha} \int_{\partial B_{1}} u_{t}^{2}\right] } \\
&+ \varepsilon\left(\log \frac{1}{t}\right)^{-\sigma \gamma}\left(1-t^{\alpha}\right)^{\gamma+1} \times  \tag{1.10.5}\\
& \times\left[\frac{\left(\mathrm{e}^{-a t^{\alpha}} W_{\kappa}(t, u)\right)^{\gamma+1}}{2^{\gamma}}-\left(\kappa(b+1) t^{\alpha} \int_{\partial B_{1}} u_{t}^{2}\right)^{\gamma+1}\right] \\
&=\left(1-t^{\alpha}\right) \mathrm{e}^{-a t^{\alpha}} W_{\kappa}(t, u) \\
& \quad+\frac{\varepsilon}{2^{\gamma}}\left(\log \frac{1}{t}\right)^{-\sigma \gamma}\left(1-t^{\alpha}\right)^{\gamma+1} \mathrm{e}^{-a(\gamma+1) t^{\alpha}} W_{\kappa}(t, u)^{\gamma+1} \\
& \quad-\left(1-t^{\alpha}\right) \kappa(b+1) t^{\alpha} \int_{\partial B_{1}} u_{t}^{2} \\
& \quad-\varepsilon\left(\log \frac{1}{t}\right)^{-\sigma \gamma}\left(1-t^{\alpha}\right)^{\gamma+1} \kappa^{\gamma+1}(b+1)^{\gamma+1} t^{\alpha(\gamma+1)}\left(\int_{\partial B_{1}} u_{t}^{2}\right)^{\gamma+1},
\end{align*}
$$

where we used (1.10.3) in the second inequality and the convexity of $x \mapsto x^{\gamma+1}$ on $\mathbb{R}_{+}$in the third inequality. Now (1.10.2) and (1.10.5), together with Lemma 1.7.1, yield

$$
\begin{equation*}
\frac{d}{d t} W_{\kappa}(t, u) \geq-C_{1} t^{\alpha-1} W_{\kappa}(t, u)+C_{2} t^{-1}\left(\log \frac{1}{t}\right)^{-\sigma \gamma} W_{\kappa}(t, u)^{\gamma+1}-C_{3} t^{\alpha / 2-1} \tag{1.10.6}
\end{equation*}
$$

where $C_{i}=C_{i}(\sigma)$. Summarizing, we have that at every $0<t<t_{0}(\sigma)$, either (1.10.1), (1.10.6), or the bound (1.10.4) holds. Further note that by the growth estimate in Lemma 1.7.1, the bound (1.10.1) implies (1.10.6) for sufficiently small $t$ and thus we may assume that (1.10.6) holds for all $0<t<t_{0}$ for which $W_{\kappa}(t, u)>C t^{\alpha / 2}$.

To proceed, let $0<t<t_{0}$ be such that $W_{\kappa}(t, u) \geq t^{\alpha / 8}$. Then the bound (1.10.6) holds and we can derive that for $C=\frac{\gamma C_{2}}{2(1-\sigma \gamma)}$, we have

$$
\begin{aligned}
& \frac{d}{d t}\left(-W_{\kappa}(t, u)^{-\gamma} \mathrm{e}^{-t^{\alpha / 4}}+C\left(\log \frac{1}{t}\right)^{1-\sigma \gamma}\right) \\
& =W_{\kappa}(t, u)^{-\gamma-1} \mathrm{e}^{-t^{\alpha / 4}}\left(\gamma \frac{d}{d t} W_{\kappa}(t, u)+\frac{\alpha}{4} W_{\kappa}(t, u) t^{\alpha / 4-1}\right)-C(1-\sigma \gamma) t^{-1}\left(\log \frac{1}{t}\right)^{-\sigma \gamma} \\
& \geq \\
& \quad W_{\kappa}(t, u)^{-\gamma} \mathrm{e}^{-t^{\alpha / 4}} t^{\alpha / 4-1}\left(\frac{\alpha}{4}-\gamma C_{1} t^{3 \alpha / 4}-\frac{\gamma C_{3} t^{\alpha / 4}}{W_{\kappa}(t, u)}\right) \\
& \quad \quad+\left(\log \frac{1}{t}\right)^{-\sigma \gamma} t^{-1}\left(\mathrm{e}^{-t^{\alpha / 4}} \gamma C_{2}-C(1-\sigma \gamma)\right) \\
& \geq 0
\end{aligned}
$$

$0<t<t_{0}=t_{0}(\sigma)$. Since also the function $-t^{-\gamma(\alpha / 8)} \mathrm{e}^{-t^{\alpha / 4}}+C\left(\log \frac{1}{t}\right)^{1-\sigma \gamma}$ is nondecreasing for small $t$, denoting

$$
\widehat{W}_{\kappa}(t, u)=\max \left\{W_{\kappa}(t, u), t^{\alpha / 8}\right\}
$$

we obtain that the function

$$
-\widehat{W}_{\kappa}(t, u)^{-\gamma} \mathrm{e}^{-t^{\alpha / 4}}+C\left(\log \frac{1}{t}\right)^{1-\sigma \gamma}
$$

is nondecreasing on $\left(0, t_{0}\right)$. Hence,

$$
\begin{aligned}
-\widehat{W}_{\kappa}(t, u)^{-\gamma} \mathrm{e}^{-t^{\alpha / 4}}+C\left(\log \frac{1}{t}\right)^{1-\sigma \gamma} & \leq-\widehat{W}_{\kappa}\left(t_{0}, u\right)^{-\gamma} \mathrm{e}^{-t_{0}^{\alpha / 4}}+C\left(\log \frac{1}{t_{0}}\right)^{1-\sigma \gamma} \\
& \leq C\left(\log \frac{1}{t_{0}}\right)^{1-\sigma \gamma}
\end{aligned}
$$

If $0<t<t_{0}^{2}$, then $\left(\log \frac{1}{t_{0}}\right)^{1-\sigma \gamma}<\left(\frac{1}{2}\right)^{1-\sigma \gamma}\left(\log \frac{1}{t}\right)^{1-\sigma \gamma}$, implying that

$$
-\widehat{W}_{\kappa}(t, u)^{-\gamma} \mathrm{e}^{-t^{\alpha / 4}} \leq C\left((1 / 2)^{1-\sigma \gamma}-1\right)\left(\log \frac{1}{t}\right)^{1-\sigma \gamma}
$$

and hence

$$
W_{\kappa}(t, u) \leq \widehat{W}_{\kappa}(t, u) \leq C\left(1-(1 / 2)^{1-\sigma \gamma}\right)^{-\frac{1}{\gamma}}\left(\log \frac{1}{t}\right)^{-\frac{1-\sigma \gamma}{\gamma}} .
$$

Lemma 1.10.4. If $u$ is as in Lemma 1.10 .3 with $\frac{2}{n-2}<\sigma \leq 1$, then there exist positive $C=C(\sigma), t_{0}=t_{0}(\sigma)$ such that

$$
\int_{\partial B_{t}} u^{2} \leq C\left(\log \frac{1}{t}\right)^{\sigma-\frac{2}{n-2}} t^{n+2 \kappa-1}, \quad 0<t<t_{0}
$$

Proof. Going back to the proof and notations of Lemma 1.7.1, we have that for $0<s<t<t_{0}$

$$
|m(t)-m(s)| \leq C\left(\log \frac{t}{s}\right)^{1 / 2}\left(W_{\kappa}(t)-W_{\kappa}(s)\right)^{1 / 2}
$$

Let now $0 \leq \mathrm{j} \leq i$ be such that $2^{-2^{i+1}}<t \leq 2^{-2^{i}}, 2^{-2^{j+1}}<t_{0} \leq 2^{-2^{j}}$. Then

$$
\begin{aligned}
\left|m\left(t_{0}\right)-m(t)\right| & \leq\left|m\left(t_{0}\right)-m\left(2^{-2^{j+1}}\right)\right|+\left|m\left(2^{-2^{i}}\right)-m(t)\right|+\sum_{k=\mathrm{j}+1}^{i-1}\left|m\left(2^{-2^{k}}\right)-m\left(2^{-2^{k+1}}\right)\right| \\
& \leq \sum_{k=0}^{i} C\left[\log \left(2^{-2^{k}}\right)-\log \left(2^{-2^{k+1}}\right)\right]^{1 / 2}\left[W_{\kappa}\left(2^{-2^{k}}\right)-W_{\kappa}\left(2^{-2^{k+1}}\right)\right]^{1 / 2} \\
& \leq C \sum_{k=0}^{i} 2^{k / 2} W_{\kappa}\left(2^{-2^{k}}\right)^{1 / 2} \leq C \sum_{k=0}^{i} 2^{\left(1-\frac{1-\sigma \gamma}{\gamma}\right) k / 2} \\
& \leq C 2^{\left(\sigma-\frac{2}{n-2}\right) i / 2} \leq C\left(\log \frac{1}{t}\right)^{\frac{1}{2}\left(\sigma-\frac{2}{n-2}\right)}
\end{aligned}
$$

Note that in the fifth inequality we have used that $1-\frac{1-\sigma \gamma}{\gamma}=\sigma-\frac{2}{n-2}>0$. Thus

$$
m(t) \leq m\left(t_{0}\right)+C\left(\log \frac{1}{t}\right)^{\frac{1}{2}\left(\sigma-\frac{2}{n-2}\right)} \leq C\left(\log \frac{1}{t}\right)^{\frac{1}{2}\left(\sigma-\frac{2}{n-2}\right)}
$$

This implies the desired result.
Lemma 1.10.3 and Lemma 1.10.4 imply the following.

Corollary 1.10.5 (Bootstraping). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$ such that $0 \in \Gamma(u)$ and $\widehat{N}(0+, u)=\kappa<\kappa_{0}, \kappa=2 m, m \in \mathbb{N}$. Suppose that for $\frac{2}{n-2}<\leq 1$

$$
\int_{\partial B_{t}} u^{2} \leq C(\sigma)\left(\log \frac{1}{t}\right)^{\sigma} t^{n+2 \kappa-1}, \quad 0<t<t_{0}(\sigma) .
$$

Then

$$
\int_{\partial B_{t}} u^{2} \leq C^{\prime}(\sigma)\left(\log \frac{1}{t}\right)^{\sigma-\frac{2}{n-2}} t^{n+2 \kappa-1}, \quad 0<t<t_{0}^{\prime}(\sigma) .
$$

Lemma 1.10.6 (Optimal growth estimate at sigular points). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$ such that $0 \in \Gamma(u)$ and $\widehat{N}(0+, u)=\kappa<\kappa_{0}, \kappa=2 m, m \in \mathbb{N}$. Then, for $0<t<t_{0}$,

$$
\begin{aligned}
& \int_{\partial B_{t}} u^{2} \leq C t^{n+2 \kappa-1}, \\
& \int_{B_{t}}|\nabla u|^{2} \leq C t^{n+2 \kappa-2} .
\end{aligned}
$$

Proof. Starting with $\sigma=1$ in Lemma 1.7.1 and repeatedly applying Corollary 1.10.5, we find $0<\leq \min \left\{\frac{2}{n-2}, 1\right\}$ such that

$$
\int_{\partial B_{t}} u^{2} \leq C\left(\log \frac{1}{t}\right)^{\sigma} t^{n+2 \kappa-1}, \quad 0<t<t_{0}
$$

In fact, we can make $\sigma$ to be strictly less than $\frac{2}{n-2}$ by noticing that in Lemma 1.10.4 we can replace $\frac{2}{n-2}$ by any smaller positive number. Then by Lemma 1.10.3

$$
0 \leq W_{\kappa}(t, u) \leq C\left(\log \frac{1}{t}\right)^{-\frac{1-\sigma \gamma}{\gamma}}
$$

Recall also that for $0<s<t<t_{0}$

$$
|m(t)-m(s)| \leq C\left(\log \frac{t}{s}\right)^{1 / 2}\left(W_{\kappa}(t)-W_{\kappa}(s)\right)^{1 / 2}
$$

We then again consider the exponentially dyadic decomposition as in the proof of Lemma 1.10.4.
Let $0 \leq \mathrm{j} \leq i$ be such that $2^{-2^{i+1}} \leq s / t_{0}<2^{-2^{i}}$ and $2^{-2^{j+1}} \leq t / t_{0}<2^{-2^{j}}$. Then,

$$
\begin{align*}
|m(t)-m(s)| & \leq C \sum_{k=\mathrm{j}}^{i} 2^{k / 2} W_{\kappa}\left(2^{-2^{k}} t_{0}\right)^{1 / 2} \leq C \sum_{k=\mathrm{j}}^{\infty} 2^{\left(1-\frac{1-\sigma \gamma}{\gamma}\right) k / 2}  \tag{1.10.7}\\
& \leq C 2^{\left(\sigma-\frac{2}{n-2}\right) \mathrm{j} / 2} \leq C\left(\log \frac{1}{t}\right)^{\left(\sigma-\frac{2}{n-2}\right) / 2}
\end{align*}
$$

Particularly,

$$
m(t) \leq m\left(t_{0}\right)+C\left(\log \frac{1}{t_{0}}\right)^{\left(\sigma-\frac{2}{n-2}\right) / 2}
$$

This gives the first bound. The second bound is obtained from the first one by arguing as at the end of Lemma 1.7.1.

Remark 1.10.7. The growth estimates in Lemma 1.10.6 enable us to consider $\kappa$-homogeneous blowups

$$
u_{t_{\mathrm{j}}}^{\phi} \rightarrow u_{0}^{\phi} \quad \text { in } \quad C_{\mathrm{loc}}^{1}\left(\mathbb{R}_{ \pm}^{n} \cup \mathbb{R}^{n-1}\right)
$$

for $t=t_{\mathrm{j}} \rightarrow 0+$, similar to 3/2-homogeneous blowups, defined at the beginning of Section 1.7, see Remark 1.8.1.

Proposition 1.10.2. Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$ such that $0 \in \Gamma(u)$ and $\widehat{N}(0+, u)=\kappa<\kappa_{0}, \kappa=2 m, m \in \mathbb{N}$. Then there exist $C>0$ and $t_{0}>0$ such that

$$
\int_{\partial B_{1}}\left|u_{t}^{\phi}-u_{s}^{\phi}\right| \leq C\left(\log \frac{1}{t}\right)^{-\frac{1-\gamma}{2 \gamma}}, \quad 0<t<t_{0}
$$

In particular the blowup $u_{0}^{\phi}$ is unique.
Proof. Using Lemma 1.10.6, we apply Lemma 1.10 .3 with $=0$ to obtain

$$
0 \leq W_{\kappa}(t, u) \leq C\left(\log \frac{1}{t}\right)^{-\frac{1}{\gamma}}
$$

Recall now the estimate

$$
\int_{\partial B_{1}}\left|u_{t}^{\phi}-u_{s}^{\phi}\right| \leq C\left(\log \frac{t}{s}\right)^{1 / 2}\left(W_{\kappa}(t)-W_{\kappa}(s)\right)^{1 / 2}
$$

for $0<s<t<t_{0}$, that we proved in Lemma 1.8.2 in the case $\kappa=3 / 2$ - the proof actually works for any $0<\kappa<\kappa_{0}$. Then, applying the exponentially dyadic argument as in the proof of Lemma 1.10.6, we obtain

$$
\int_{\partial B_{1}}\left|u_{t}^{\phi}-u_{s}^{\phi}\right| \leq C\left(\log \frac{1}{t}\right)^{-\frac{1-\gamma}{2 \gamma}}
$$

Lemma 1.10.8 (Nondegeneracy). Let 0 be a free boundary point of $u$ such that $\widehat{N}(0+, u)=\kappa$, $\kappa=2 m, m \in \mathbb{N}$. Then

$$
\liminf _{t \rightarrow 0} \int_{\partial B_{1}}\left(u_{t}^{\phi}\right)^{2}=\liminf _{t \rightarrow 0} \frac{1}{t^{n+2 \kappa-1}} \int_{\partial B_{t}} u^{2}>0
$$

Proof. We use the approach of Lemma 7.2 in [27]. Assume to the contrary that for some $r_{\mathrm{j}} \searrow 0+$

$$
\lim _{\mathrm{j} \rightarrow \infty} \frac{1}{r_{\mathrm{j}}^{n+2 \kappa-1}} \int_{\partial B_{r_{\mathrm{j}}}} u^{2}=0
$$

Consider then the corresponding Almgren rescalings $u_{r_{\mathrm{j}}}^{A}(x)$. By Proposition 1.6.1, over a subsequence, $u_{r_{\mathrm{j}}}^{A} \rightarrow q$ for some blowup $q$. By a characterization of singular points in Proposition 1.10.1, $q$ is $\kappa$-homogeneous and is normalized by $\|q\|_{L^{2}\left(\partial B_{1}\right)}=1$. Next, for each Almgren rescaling $u_{r_{\mathrm{j}}}^{A}$ consider its $\kappa$-almost homogeneous rescalings

$$
\left[u_{r_{\mathrm{j}}}^{A}\right]_{t}^{\phi}:=\frac{u_{r_{\mathrm{j}}}^{A}(t x)}{\phi(t)} .
$$

Since $u_{r_{\mathrm{j}}}^{A}$ is an almost minimizer in $B_{1 / r_{\mathrm{j}}}$ with gauge function $\omega(t)=\left(r_{\mathrm{j}} t\right)^{\alpha}$, we have

$$
N\left(0+, u_{r_{\mathrm{j}}}^{A}\right)=\lim _{s \rightarrow 0+} N\left(s, u_{r_{\mathrm{j}}}^{A}\right)=\lim _{s \rightarrow 0+} N\left(r_{\mathrm{j}} s, u\right)=N(0+, u)=\kappa .
$$

Thus, by Proposition 1.10.2, over subsequences, $\left[u_{r_{\mathrm{j}}}^{A}\right]_{t}^{\phi}$ converges to a unique blowup $q_{r_{\mathrm{j}}}$ and

$$
\int_{\partial B_{1}}\left|\left[u_{r_{\mathrm{j}}}^{A}\right]_{t}^{\phi}-q_{r_{\mathrm{j}}}\right| \leq C\left(\log \frac{1}{t}\right)^{-\frac{1-\gamma}{2 \gamma}}, \quad 0<t<t_{0}
$$

Notice that since $\left\|u_{r_{\mathrm{j}}}^{A}\right\|_{W^{1,2}\left(B_{1}\right)}$ is uniformly bounded, the constant $C$ is independent of $r_{\mathrm{j}}, t$. Now we fix $r_{\mathrm{j}}$, and consider a sequence $\left\{\rho_{i}\right\}_{i=1}^{\infty}=\left\{r_{i} / r_{\mathrm{j}}\right\}_{i=1}^{\infty}$. Note that up to subsequence, $\left[u_{r_{\mathrm{j}}}^{A}\right]_{\rho_{i}}^{\phi} \rightarrow q_{r_{\mathrm{j}}}$ as $\rho_{i} \rightarrow 0$, by the uniqueness. Then

$$
\begin{aligned}
\int_{\partial B_{1}} q_{r_{\mathrm{j}}}^{2} & =\lim _{\rho_{i} \rightarrow 0} \frac{1}{\rho_{i}^{n+2 \kappa-1}} \int_{\partial B_{\rho_{i}}}\left(u_{r_{\mathrm{j}}}^{A}\right)^{2} \frac{r_{\mathrm{j}}^{n+2 \kappa-1}}{\int_{\partial B_{r_{\mathrm{r}}}} u^{2}} \lim _{i \rightarrow \infty} \frac{1}{\left(r_{\mathrm{j}} \rho_{i}\right)^{n+2 \kappa-1}} \int_{\partial{r_{r_{\mathrm{j}} \rho_{i}}} u^{2}} \\
& =\frac{r_{\mathrm{j}}^{n+2 \kappa-1}}{\int_{\partial B_{r_{\mathrm{j}}}} u^{2}} \lim _{i \rightarrow \infty} \frac{1}{r_{i}^{n+2 \kappa-1}} \int_{\partial B_{r_{i}}} u^{2}=0
\end{aligned}
$$

by the contradiction assumption. Thus, $q_{r_{\mathrm{j}}}=0$ on $\partial B_{1}$, and hence

$$
\int_{\partial B_{1}} \left\lvert\,\left[u_{r_{j}, t}^{A}{ }^{\phi} \left\lvert\, \leq C\left(\log \frac{1}{t}\right)^{-\frac{1-\gamma}{2 \gamma}} .\right.\right.\right.
$$

Now for any $\rho>0$ and $r_{\mathrm{j}}$,

$$
\begin{aligned}
1 & =\frac{1}{\rho^{n+2 \kappa-1}} \int_{\partial B_{\rho}} q^{2} \\
& \leq \frac{\|q\|_{L^{\infty}\left(\partial B_{\rho}\right)}}{\rho^{\kappa}} \frac{1}{\rho^{n+\kappa-1}} \int_{\partial B_{\rho}}|q| \\
& \leq\|q\|_{L^{\infty}\left(\partial B_{1}\right)}\left[\frac{1}{\rho^{n+\kappa-1}} \int_{\partial B_{\rho}}\left|q-u_{r_{\mathrm{j}}}^{A}\right|+\frac{1}{\rho^{n+\kappa-1}} \int_{\partial B_{\rho}}\left|u_{r_{\mathrm{j}}}^{A}\right|\right] \\
& \leq\|q\|_{L^{\infty}\left(\partial B_{1}\right)}\left[\frac{1}{\rho^{n+\kappa-1}} C_{n} \rho^{\frac{n-1}{2}}\left(\int_{\partial B_{\rho}}\left|q-u_{r_{\mathrm{j}}}^{A}\right|^{2}\right)^{1 / 2}+\mathrm{e}^{-\left(\frac{\kappa b}{\alpha}\right) \rho^{\alpha}} \int_{\partial B_{1}}\left|\left[u_{r_{\mathrm{j}}}^{A}\right]_{\rho}^{\phi}\right|\right] \\
& \leq C\|q\|_{L^{\infty}\left(\partial B_{1}\right)}\left[\left(\frac{1}{\rho^{n+2 \kappa-1}} \int_{\partial B_{\rho}}\left|q-u_{r_{\mathrm{j}}}^{A}\right|^{2}\right)^{1 / 2}+\left(\log \frac{1}{\rho}\right)^{-\frac{1-\gamma}{2 \gamma}}\right]
\end{aligned}
$$

Note that $u_{r_{\mathrm{j}}}^{A} \rightarrow q$ in $C_{\text {loc }}^{1}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)$. We choose first $\rho>0$ small and then $r_{\mathrm{j}}=r_{\mathrm{j}}(\rho)>0$ small to reach a contradiction.

The nondegeneracy implies the following important fact, which enables the use of the Whitney Extension Theorem in the proof of the structural theorem on the singular set (Theorem 1.10.10 below).

For $\kappa=2 m<\kappa_{0}, m \in \mathbb{N}$, we denote

$$
\Sigma_{\kappa}(u):=\left\{x_{0} \in \Sigma(u): N\left(0+, u, x_{0}\right)=\kappa\right\} .
$$

Lemma 1.10.9. The set $\Sigma_{\kappa}(u)$ is of topological type $F_{\sigma}$; i.e., it is a countable union of closed sets.

Proof. For $\mathrm{j} \in \mathbb{N}, \mathrm{j} \geq 2$, let

$$
E_{\mathrm{j}}:=\left\{x_{0} \in \Sigma_{\kappa}(u) \cap \overline{B_{1-1 / \mathrm{j}}}: \frac{1}{\mathrm{j}} \leq \frac{1}{\rho^{n+2 \kappa-1}} \int_{\partial B_{\rho}\left(x_{0}\right)} u^{2} \leq \mathrm{j} \text { for } 0<\rho<\frac{1}{2 \mathrm{j}}\right\}
$$

Then by Lemma 1.10.6 and Lemma 1.10.8, $\Sigma_{\kappa}(u)=\bigcup_{\mathrm{j}=2}^{\infty} E_{\mathrm{j}}$. We now claim that $E_{\mathrm{j}}$ is closed for any $\mathrm{j} \geq 2$. Indeed, take a sequence $x_{i} \in E_{\mathrm{j}}$ such that $x_{i} \rightarrow x_{0}$ as $i \rightarrow \infty$. Then $x_{0} \in \overline{B_{1-1 / \mathrm{j}}}$ and for every $0<\rho<1 /(2 \mathrm{j})$, by the local uniform continuity of $u$,

$$
\begin{equation*}
\frac{1}{\rho^{n+2 \kappa-1}} \int_{\partial B_{\rho}\left(x_{0}\right)} u^{2}=\lim _{i \rightarrow \infty} \frac{1}{\rho^{n+2 \kappa-1}} \int_{\partial B_{\rho}\left(x_{i}\right)} u^{2} \in\left[\frac{1}{\mathrm{j}}, \mathrm{j}\right] . \tag{1.10.8}
\end{equation*}
$$

Next, since $\Gamma(u)$ is relatively closed in $B_{1}^{\prime}$, we also know that $x_{0} \in \Gamma(u)$. Moreover, since $N\left(0+, u, x_{i}\right)=\kappa$ and the function $x \mapsto \widehat{N}(0+, u, x)$ is upper semicontinuous, we have

$$
\kappa=\limsup _{i \rightarrow \infty} \widehat{N}\left(0+, u, x_{i}\right) \leq \widehat{N}\left(0+, u, x_{0}\right) .
$$

If $\widehat{N}\left(0+, u, x_{0}\right)=\kappa^{\prime}>\kappa$, then by Lemma 1.7.1,

$$
\frac{1}{\rho^{n+2 \kappa-1}} \int_{\partial B_{\rho}\left(x_{0}\right)} u^{2} \leq C \rho^{2\left(\kappa^{\prime}-\kappa\right)}\left(\log \frac{1}{\rho}\right) \rightarrow 0 \quad \text { as } \rho \rightarrow 0,
$$

which contradicts (1.10.8). Therefore, $\widehat{N}\left(0+, u, x_{0}\right)=\kappa$ and consequently $x_{0} \in E_{j}$. Hence, $E_{\mathrm{j}}$ is closed, $\mathrm{j}=2,3, \ldots$, implying that $\Sigma_{\kappa}(u)$ is $F_{\sigma}$.

To state the main result of this chapter concerning the singular points, we need to introduce the following notations. For $\kappa=2 m<\kappa_{0}, m \in \mathbb{N}$ and $x_{0} \in \Sigma_{\kappa}(u)$, we define

$$
d_{x_{0}}^{(\kappa)}:=\operatorname{dim}\left\{\xi \in \mathbb{R}^{n-1}: \xi \cdot \nabla_{x} u_{x_{0}}^{\phi}(x, 0) \equiv 0 \text { on } \mathbb{R}^{n-1}\right\},
$$

which has the meaning of the dimension of $\Sigma_{\kappa}(u)$ at $x_{0}$, and where $u_{x_{0}}^{\phi}$ is the unique $\kappa$ homogeneous blowup at $x_{0}$. In fact, $d_{x_{0}}^{(\kappa)}$ is the dimension of the linear subspace $\Sigma_{\kappa}\left(u_{x_{0}}^{\phi}\right) \subset$ $\mathbb{R}^{n-1}$. Since $u_{x_{0}}^{\phi}$ is a nonzero solution of the Signorini problem, it cannot vanish identically on $\mathbb{R}^{n-1}$ (see [14]) and therefore $d_{x_{0}}^{(\kappa)}<n-1$.

For $d=0,1, \ldots, n-2$, we denote

$$
\Sigma_{\kappa}^{d}(u):=\left\{x_{0} \in \Sigma_{\kappa}(u): d_{x_{0}}^{(\kappa)}=d\right\}
$$

Theorem 1.10.10 (Structure of the singular set). Let $u$ be an almost minimizer for the Signorini problem in $B_{1}$. Then for every $\kappa=2 m<\kappa_{0}, m \in \mathbb{N}$, and $d=0,1, \ldots, n-2$, the set $\Sigma_{\kappa}^{d}(u)$ is contained in the union of countably many submanifolds of dimension $d$ and class $C^{1, \log }$.

Proof. Let $\kappa=2 m, m \in \mathbb{N}$. For $x \in \Sigma_{\kappa}(u) \cap B_{1 / 2}^{\prime}$, let $q_{x} \in \mathcal{Q}_{\kappa}$ denote the unique $\kappa$ homogeneous blowup of $u$. By the optimal growth (Lemma 1.10.6) and the nondegeneracy (Lemma 1.10.8), we can write

$$
q_{x}=\lambda_{x} q_{x}^{A}, \quad \lambda_{x}>0, \quad\left\|q_{x}^{A}\right\|_{L^{2}\left(\partial B_{1}\right)}=1
$$

where $q_{x}^{A} \in \mathcal{Q}_{\kappa}$ is the corresponding Almgren blowup. We want to show that the $q_{x}, q_{x}^{A}, \lambda_{x}$ depend continuously on $x \in \Sigma_{\kappa}$, with a logarithmic modulus of continuity.

Let $x_{1}, x_{2} \in \Sigma_{\kappa}(u) \cap B_{1 / 2}^{\prime}$. Then for $t>0$, to be chosen below, we can write

$$
\begin{equation*}
\left\|q_{x_{1}}-q_{x_{2}}\right\|_{L^{1}\left(\partial B_{1}\right)} \leq\left\|q_{x_{1}}-u_{x_{1}, t}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)}+\left\|u_{x_{1}, t}^{\phi}-u_{x_{2}, t}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)}+\left\|u_{x_{2}, t}^{\phi}-q_{x_{2}}\right\|_{L^{1}\left(\partial B_{1}\right)} . \tag{1.10.9}
\end{equation*}
$$

By Proposition 1.10.2, we have

$$
\begin{equation*}
\left\|q_{x}-u_{x, t}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)} \leq C\left(\log \frac{1}{t}\right)^{-\frac{1}{n-2}} \tag{1.10.10}
\end{equation*}
$$

for $x \in \Sigma_{\kappa}(u) \cap B_{1 / 2}^{\prime}$. This controls the first and third term on the right hand side of (1.10.9) To estimate the middle term, we observe that

$$
\left\|u_{x_{1}, t}^{\phi}-u_{x_{2}, t}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)} \leq \frac{\mathrm{e}^{\left(\frac{\kappa b}{\alpha}\right) t^{\alpha}}}{t^{\kappa}} \int_{\partial B_{1}} \int_{0}^{1}\left|\nabla u\left(x_{1}+t z+r\left(x_{2}-x_{1}\right)\right)\right|\left|x_{1}-x_{2}\right| d r d S_{z}
$$

for any $0<t<1 / 2$. Recalling that $\nabla u\left(x_{1}\right)=0$ and $u \in C^{1, \beta}\left(B_{1}^{ \pm} \cup B_{1}\right)$, we have

$$
\left\lvert\, \nabla u\left(x_{1}+t z+r\left(x_{2}-x_{1}\right)|\leq C| t z+\left.r\left(x_{2}-x_{1}\right)\right|^{\beta} \leq C\left(t+\left|x_{1}-x_{2}\right|\right)^{\beta} \leq C\left|x_{1}-x_{2}\right|^{\frac{\beta}{2(\kappa-\beta)}}\right.\right.
$$

if we choose $t=\left|x_{1}-x_{2}\right|^{\frac{1}{2(\kappa-\beta)}}$ and have $\left|x_{1}-x_{2}\right|<(1 / 2)^{2(\kappa-\beta)}$. This gives

$$
\begin{equation*}
\left\|u_{x_{1}, t}^{\phi}-u_{x_{2}, t}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)} \leq \frac{C}{t^{\kappa}}\left|x_{1}-x_{2}\right|^{\frac{\beta}{2(\kappa-\beta)}}\left|x_{1}-x_{2}\right| \leq C\left|x_{1}-x_{2}\right|^{1 / 2} . \tag{1.10.11}
\end{equation*}
$$

Combining (1.10.9), (1.10.11), and (1.10.10), we obtain

$$
\begin{equation*}
\left\|q_{x_{1}}-q_{x_{2}}\right\|_{L^{1}\left(\partial B_{1}\right)} \leq C\left(\log \frac{1}{\left|x_{1}-x_{2}\right|}\right)^{-\frac{1}{n-2}} \tag{1.10.12}
\end{equation*}
$$

Next, by Lemma 1.10.6, for any $x \in \Sigma_{\kappa}(u) \cap B_{1 / 2}^{\prime}$ and small $t$

$$
\int_{\partial B_{1}}\left(u_{x, t}^{\phi}\right)^{2} \leq C
$$

with $C$ independent of $x$, and passing to the limit as $t \rightarrow \infty$ obtain the bound

$$
\lambda_{x}^{2}=\int_{\partial B_{1}} q_{x}^{2} \leq C
$$

Moreover, since $q_{x}$ is a $\kappa$-homogeneous harmonic polynomial, we also have

$$
\begin{equation*}
\left\|q_{x}\right\|_{L^{\infty}\left(B_{1}\right)} \leq C(n, \kappa)\left\|q_{x}\right\|_{L^{2}\left(\partial B_{1}\right)} \leq C \tag{1.10.13}
\end{equation*}
$$

Then, by combining (1.10.12) and (1.10.13), we have

$$
\begin{align*}
\left|\lambda_{x_{1}}-\lambda_{x_{2}}\right| & \leq\left|\lambda_{x_{1}}^{2}-\lambda_{x_{2}}^{2}\right|^{1 / 2} \leq\left(\int_{\partial B_{1}}\left|q_{x_{1}}^{2}-q_{x_{2}}^{2}\right|\right)^{1 / 2} \\
& \leq\left\|q_{x_{1}}+q_{x_{2}}\right\|_{L^{\infty}\left(B_{1}\right)}^{1 / 2} \mid q_{x_{1}}-q_{x_{1}} \|_{L^{1}\left(\partial B_{1}\right)}^{1 / 2}  \tag{1.10.14}\\
& \leq C\left(\log \frac{1}{\left|x_{1}-x_{2}\right|}\right)^{-\frac{1}{2(n-2)}}
\end{align*}
$$

Finally, we want to estimate $q_{x_{1}}^{A}-q_{x_{2}}^{A}$. By writing

$$
\begin{aligned}
\left\|q_{x_{1}}-q_{x_{2}}\right\|_{L^{1}\left(\partial B_{1}\right)} & =\int_{\partial B_{1}}\left|\lambda_{x_{1}} q_{x_{1}}^{A}-\lambda_{x_{2}} q_{x_{2}}^{A}\right| \\
& =\int_{\partial B_{1}}\left|\lambda_{x_{1}}\left(q_{x_{1}}^{A}-q_{x_{2}}^{A}\right)+\left(\lambda_{x_{1}}-\lambda_{x_{2}}\right) q_{x_{2}}^{A}\right|
\end{aligned}
$$

$$
\geq \lambda_{x_{1}} \int_{\partial B_{1}}\left|q_{x_{1}}^{A}-q_{x_{2}}^{A}\right|-\left|\lambda_{x_{1}}-\lambda_{x_{2}}\right| \int_{\partial B_{1}}\left|q_{x_{2}}^{A}\right|,
$$

we estimate

$$
\begin{align*}
\lambda_{x_{1}} \int_{\partial B_{1}}\left|q_{x_{1}}^{A}-q_{x_{2}}^{A}\right| & \leq\left\|q_{x_{1}}-q_{x_{2}}\right\|_{L^{1}\left(\partial B_{1}\right)}+\left|\lambda_{x_{1}}-\lambda_{x_{2}}\right| \int_{\partial B_{1}}\left|q_{x_{2}}^{A}\right| \\
& \leq\left\|q_{x_{1}}-q_{x_{2}}\right\|_{L^{1}\left(\partial B_{1}\right)}+C(n)\left|\lambda_{x_{1}}-\lambda_{x_{2}}\right|  \tag{1.10.15}\\
& \leq C\left(\log \frac{1}{\left|x_{1}-x_{2}\right|}\right)^{-\frac{1}{2(n-2)}},
\end{align*}
$$

where we used $\left\|q_{x_{2}}^{A}\right\|_{L^{2}\left(\partial B_{1}\right)}=1$ in the second inequality and (1.10.12) and the bound (1.10.14) in the third inequality. Next, using that $q_{x}^{A}$ are $\kappa$-homogeneous harmonic polynomials, we have

$$
\left\|q_{x_{1}}^{A}-q_{x_{2}}^{A}\right\|_{L^{\infty}\left(B_{1}\right)} \leq C\left\|q_{x_{1}}^{A}-q_{x_{2}}^{A}\right\|_{L^{1}\left(\partial B_{1}\right)}
$$

which combined with (1.10.15) gives

$$
\begin{equation*}
\lambda_{x_{1}}\left\|q_{x_{1}}^{A}-q_{x_{2}}^{A}\right\|_{L^{\infty}\left(B_{1}\right)} \leq C\left(\log \frac{1}{\left|x_{1}-x_{2}\right|}\right)^{-\frac{1}{2(n-2)}} \tag{1.10.16}
\end{equation*}
$$

Now we fix $x_{0} \in \Sigma_{\kappa}(u) \cap B_{1 / 4}^{\prime}$. Then by (1.10.14), there exists $\delta=\delta\left(x_{0}\right) \in\left(0,(1 / 2)^{2(\kappa-\beta)+1}\right)$ such that $\lambda_{x} \geq 1 / 2 \lambda_{x_{0}}$ if $x \in \Sigma_{\kappa}(u) \cap B_{\delta}^{\prime}\left(x_{0}\right)$. Then by (1.10.16), we conclude that

$$
\begin{equation*}
\left\|q_{x_{1}}^{A}-q_{x_{2}}^{A}\right\|_{L^{\infty}\left(B_{1}\right)} \leq C\left(\log \frac{1}{\left|x_{1}-x_{2}\right|}\right)^{-\frac{1}{2(n-2)}}, \quad x_{1}, x_{2} \in \Sigma_{\kappa}(u) \cap B_{\delta}^{\prime}\left(x_{0}\right) \tag{1.10.17}
\end{equation*}
$$

Notice that the constant $C$ does not depend on $x_{1}, x_{2}$, but both $C$ and $\delta$ do depend on $x_{0}$.
Once we have the estimates (1.10.14) and (1.10.17), as well as Lemma 1.10.9, we can apply the Whitney Extension Theorem of Fefferman's [56], to complete the proof, see e.g., the proof of Theorem 5 in [27].

## 1.A Some examples of almost minimizers

Example 1.A.1. If $u$ is a minimizer of the functional

$$
\int_{D} a(x)|\nabla u|^{2}
$$

over the set $\mathfrak{K}_{\psi, g}(D, \mathcal{M})$ with strictly positive $a \in C^{0, \alpha}(\bar{D}), 0<\alpha \leq 1$, then $u$ is an almost minimizer for the Signorini problem with a gauge function $\omega(r)=C r^{\alpha}$.

Proof. This is rather immediate.
Example 1.A.2. Let $u$ be a solution of the Signorini problem for the Laplacian with drift with the velocity field $b \in L^{p}\left(B_{1}\right), p>n$ :

$$
\begin{aligned}
-\Delta u+b(x) \nabla u=0 & \text { in } B_{1}^{ \pm} \\
-\partial_{x_{n}} u \geq 0, \quad u \geq 0, \quad u \partial_{x_{n}} u=0 & \text { on } B_{1}^{\prime},
\end{aligned}
$$

even in $x_{n}$-variable. We understand this in the weak sense that $u$ satisfies the variational inequality

$$
\int_{B_{1}} \nabla u \nabla(w-u)+(b(x) \nabla u)(w-u) \geq 0
$$

for any competitor $w \in \mathfrak{K}_{0, u}\left(B_{1}, B_{1}^{\prime}\right)$, i.e. $w \in u+W_{0}^{1,2}\left(B_{1}\right)$ such that $w \geq 0$ on $B_{1}^{\prime}$ in the sense of traces. Then $u$ is an almost minimizer for the Signorini problem with $\psi=0$ on $\mathcal{M}=\mathbb{R}^{n-1} \times\{0\}$ and a gauge function $\omega(r)=C r^{1-n / p}$.

Proof. This example corresponds to Example 2.A.1 when $A=I$. Thus, for this proof, we refer to the proof of Example 2.A. 1

## 2. ALMOST MINIMIZERS FOR THE THIN OBSTACLE PROBLEM WITH VARIABLE COEFFICIENTS

### 2.1 Introduction and Main Results

### 2.1.1 The thin obstacle (or Signorini) problem with variable coefficients

Let $D$ be a domain in $\mathbb{R}^{n}, n \geq 2$, and $\Pi$ a smooth hypersurface (the thin space), that splits $D$ into two subdomains $D^{ \pm}: D \backslash \Pi=D^{+} \cup D^{-}$. Let $\psi: \Pi \rightarrow \mathbb{R}$ be a certain (smooth) function (the thin obstacle) and $g: \partial D \rightarrow \mathbb{R}$ (the boundary values). Let also $A(x)=\left(a_{\mathrm{ij}}(x)\right)$ be an $n \times n$ symmetric uniformly elliptic matrix, $\alpha$-Hölder continuous as a function of $x \in D$, for some $0<\alpha<1$, with ellipticity constants $0<\lambda \leq 1 \leq \Lambda<\infty$ :

$$
\lambda|\xi|^{2} \leq\langle A(x) \xi, \xi\rangle \leq \Lambda|\xi|^{2}, \quad x \in D, \xi \in \mathbb{R}^{n} .
$$

Then consider the minimizer $U$ of the energy functional

$$
\mathcal{J}_{A, D}(V)=\int_{D}\langle A(x) \nabla V, \nabla V\rangle d x
$$

over a closed convex set $\mathfrak{K}_{\psi, g}(D, \Pi) \subset W^{1,2}(D)$ defined by

$$
\mathfrak{K}_{\psi, g}(D, \Pi):=\left\{V \in W^{1,2}(D): V=g \text { on } \partial D, V \geq \psi \text { on } \Pi \cap D\right\} .
$$

Because of the unilateral constraint on the thin space $\Pi$, the problem is known as the thin obstacle problem. Away from $\Pi$, the minimizer solves a uniformly elliptic divergence form equation with variable coefficients

$$
\operatorname{div}(A(x) \nabla U)=0 \quad \text { in } D^{+} \cup D^{-}
$$

On the thin space, the minimizers satisfy

$$
\begin{aligned}
& U \geq \psi, \quad\left\langle A \nabla U, \nu^{+}\right\rangle+\left\langle A \nabla U, \nu^{-}\right\rangle \geq 0 \\
& \quad(U-\psi)\left(\left\langle A \nabla U, \nu^{+}\right\rangle+\left\langle A \nabla U, \nu^{-}\right\rangle\right)=0 \quad \text { on } D \cap \Pi
\end{aligned}
$$

in a certain weak sense, where $\nu^{ \pm}$are the exterior normals to $D^{ \pm}$on $\Pi$ and $\left\langle A \nabla U, \nu^{ \pm}\right\rangle$are understood as the limits from inside $D^{ \pm}$. These are known as the Signorini complementarity conditions and therefore the problem is often referred to as the Signorini problem with variable coefficients (or $A$-Signorini problem, for short). One of the main objects of the study is the free boundary

$$
\Gamma(U)=\partial_{\Pi}\{x \in \Pi: U(x)=\psi(x)\} \cap D,
$$

which separates the coincidence set $\{U=\psi\}$ from the noncoincidence set $\{U>\psi\}$ in $D \cap \Pi$. The set $\Gamma(U)$ is also called a thin free boundary as it lives in $\Pi$ and is expected to be of codimension two with respect to the domain $D$.

These types of problems go back to the original Signorini problem in elastostatics [4], but also appear in many applications ranging from math biology (semipermeable membranes) to boundary heat control [5] or more recently in math finance, with connection to the obstacle problem for the fractional Laplacian, through the Caffarelli-Silvestre extension [10]. The presence of the free boundary makes the problem particularly challenging and while the $C^{1, \beta}$ regularity of the minimizers (on the either side of the thin space) was known already in [6][8], the study of the free boundary became possible only after the breakthrough work of [9] on the optimal $C^{1,1 / 2}$ regularity of the minimizers. Since then there has been a significant effort in the literature to understand the structure and regularity properties of the free boundary in many different settings including equations with variable coefficients, problems for the fractional Laplacian, as well as the time-dependent problems, see e.g. [11]-[30], [55], [57], [58], and many others.

### 2.1.2 Almost minimizers

The approach we take in this chapter is by considering almost minimizers of the functional $\mathcal{J}_{A, D}$ in the sense of Anzellotti [31]. For this we need a gauge function $\omega:\left(0, r_{0}\right) \rightarrow[0, \infty)$, $r_{0}>0$, which is a nondecreasing function with $\omega(0+)=0$, as well as a family $\left\{E_{r}\left(x_{0}\right)\right\}_{0<r<r_{0}}$ of open sets for any $x_{0} \in D$, comparable to balls centered at $x_{0}$ (in what comes next, we will take it to be a family of ellipsoids).

Definition 2.1.1 (Almost minimizers). We say $U$ is an almost minimizer for the $A$-Signorini problem in $D$ if $U \in W_{\mathrm{loc}}^{1,2}(D), U \geq \psi$ on $D \cap \Pi$, and for any $E_{r}\left(x_{0}\right) \Subset D$ with $0<r<r_{0}$, we have

$$
\begin{equation*}
\int_{E_{r}\left(x_{0}\right)}\langle A \nabla U, \nabla U\rangle \leq(1+\omega(r)) \int_{E_{r}\left(x_{0}\right)}\langle A \nabla V, \nabla V\rangle, \tag{2.1.1}
\end{equation*}
$$

for any competitor function $V \in \mathfrak{K}_{\psi, U}\left(E_{r}\left(x_{0}\right)\right.$, $)$, i.e., $V$ satisfying

$$
V=U \quad \text { on } \partial E_{r}\left(x_{0}\right), \quad V \geq \psi \quad \text { on } E_{r}\left(x_{0}\right) \cap \Pi .
$$

In fact, observing that for $x, x_{0} \in D$, and $\xi \in \mathbb{R}^{n}, \xi \neq 0$

$$
\left(1-C\left|x-x_{0}\right|^{\alpha}\right) \leq \frac{\left\langle A\left(x_{0}\right) \xi, \xi\right\rangle}{\langle A(x) \xi, \xi\rangle} \leq\left(1+C\left|x-x_{0}\right|^{\alpha}\right)
$$

with $C$ depending on the ellipticity of $A$ and $\|A\|_{C^{0, \alpha}(D)}$, we can rewrite (2.1.1) in the form with frozen coefficients

$$
\begin{equation*}
\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle \leq(1+\omega(r)) \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla V, \nabla V\right\rangle, \tag{2.1.2}
\end{equation*}
$$

by replacing the gauge $\omega(r)$ with $C\left(\omega(r)+r^{\alpha}\right)$ if necessary.
An example of an almost minimizer is given in Example 2.A.1. Generally, we view almost minimizers as perturbations of minimizers in a certain sense, but in the case of variable coefficients there are even some advantages of treating minimizers themselves as almost minimizers, particularly in the sense of frozen coefficients (2.1.2).

Almost minimizers for the Signorini problem have already been studied in Chapter 1 in the case $A(x) \equiv I$, where their $C^{1, \beta}$-regularity (on the either side of the thin space) has been established and a number of technical tools such as Weiss- and Almgren-type monotonicity formulas were proved. In combination with the epiperimetric and log-epiperimetric inequalities these tools allowed to establish the optimal growth and prove the $C^{1, \gamma}$-regularity of the regular set and a structural theorem on the singular set. The aim of this chapter is to extend these results to the variable coefficient case. It is noteworthy that the results that we obtain (see Theorems F-I below) for almost minimizers improve even on some of the results
available for the minimizers. For example, we only need the coefficients $A(x)$ to be $C^{0, \alpha}$ with arbitrary $0<\alpha<1$ in order to study the free boundary, compared to $W^{1, p}$, $p>n$, in [26] or $C^{0, \alpha}, 1 / 2<\alpha<1$, in [55] for the regular part of the free boundary and $C^{0,1}$ in [58] for the singular set.

### 2.1.3 Main results

Since we are interested in local regularity results, we will assume that $D=B_{1}$, the unit ball in $\mathbb{R}^{n}$, and that

$$
\Pi=\mathbb{R}^{n-1} \times\{0\}
$$

after a local diffeomorphism. In this chapter, we will consider only the case when the thin obstacle is identically zero: $\psi \equiv 0$.

Further, we will assume $r_{0}=1$ in Definition 2.1.1 and take $\left\{E_{r}\left(x_{0}\right)\right\}$ to be the family of ellipsoids associated with the positive symmetric matrix $A\left(x_{0}\right)$ :

$$
E_{r}\left(x_{0}\right):=A^{1 / 2}\left(x_{0}\right)\left(B_{r}\right)+x_{0} .
$$

By the ellipticity of $A\left(x_{0}\right)$, we have

$$
B_{\lambda^{1 / 2} r}\left(x_{0}\right) \subset E_{r}\left(x_{0}\right) \subset B_{\Lambda^{1 / 2} r}\left(x_{0}\right)
$$

To simplify the tracking of the constants, we will assume that there is $M>0$ such that

$$
\begin{equation*}
\|A\|_{C^{0, \alpha}\left(B_{1}\right)} \leq M, \quad \lambda^{-1}, \Lambda \leq M, \quad \omega(r) \leq M r^{\alpha}, \quad 0<\alpha<1 . \tag{2.1.3}
\end{equation*}
$$

Then we can go between almost minimizing properties (2.1.1) and (2.1.2) by changing $M$ if necessary.

Then our first result is as follows.

Theorem $\mathbf{F}$ ( $C^{1, \beta}$-regularity of almost minimizers). Let $U \in W^{1,2}\left(B_{1}\right)$ be an almost minimizer for the $A$-Signorini problem in $B_{1}$, under the assumptions above. Then, $U \in$ $C_{\mathrm{loc}}^{1, \beta}\left(B_{1}^{ \pm} \cup B_{1}\right)$ for $\beta=\beta(\alpha, n) \in(0,1)$ and

$$
\|U\|_{C^{1, \beta}(K)} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}
$$

for any $K \Subset B_{1}^{ \pm} \cup B_{1}^{\prime}$ and $C=C(n, \alpha, M, K)$.
The proof is obtained by using Morrey and Campanato space estimates, following the original idea of Anzellotti [31] that was successfully used in the constant coefficient case of our problem in Chapter 1. We explicitly mention, however, that in the above theorem we do not require the even symmetry of the almost minimizer in the $x_{n}$-variable, so Theorem F extends the corresponding result in Chapter 1 also in that respect.

To state our results related to the free boundary, we need to assume the following quasisymmetry condition. For $x_{0} \in B_{1}^{\prime}=B_{1} \cap \Pi$, let

$$
P_{x_{0}}=I-2 \frac{A\left(x_{0}\right) \mathrm{e}_{n} \otimes \mathrm{e}_{n}}{a_{n n}\left(x_{0}\right)}
$$

be a matrix corresponding to the reflection with respect to $\Pi$ in the conormal direction $A\left(x_{0}\right) \mathrm{e}_{n}$ at $x_{0}$. Note that $P_{x_{0}} x=x$ for any $x \in \Pi$ and $P_{x_{0}} E_{r}\left(x_{0}\right)=E_{r}\left(x_{0}\right)$. Then, for a function $U$ in $B_{1}$ define

$$
U_{x_{0}}^{*}(x):=\frac{U(x)+U\left(P_{x_{0}} x\right)}{2} .
$$

Note that $U_{x_{0}}^{*}$ may not be defined in all of $B_{1}$, but is defined in any ellipsoid $E_{r}\left(x_{0}\right)$ as long as it is contained in $B_{1}$. Note also that $U=U_{x_{0}}^{*}$ on $\Pi$.

Definition 2.1.2 (Quasisymmetry). We say that $U \in W^{1,2}\left(B_{1}\right)$ is $A$-quasisymmetric with respect to $\Pi$, if there is a constant $Q$ such that

$$
\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle \leq Q \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U_{x_{0}}^{*}, \nabla U_{x_{0}}^{*}\right\rangle,
$$

for any ellipsoid $E_{r}\left(x_{0}\right) \Subset B_{1}$ centered at any $x_{0} \in B_{1}^{\prime}$.
We will assume $Q \leq M$ throughout the chapter, in addition to (2.1.3).

Note that when $A(x) \equiv I$ and $U$ is even in $x_{n}$, then it is automatically quasisymmetric in the sense of the above definition. The quasisymmetry condition will also hold for even minimizers if $\mathrm{e}_{n}$ is an eigenvector of $A\left(x_{0}\right)$ for any $x_{0} \in B_{1}^{\prime}$, i.e., when

$$
a_{i n}\left(x_{0}\right)=0, \quad \text { for } i=1, \ldots, n-1, x_{0} \in B_{1}^{\prime} .
$$

This condition is typically imposed in the existing literature and can be satisfied with an application of a local $C^{1, \alpha}$-diffeomorphism that preserves $\Pi$, see [15], [55], [59]. The reason for a quasisymmetry condition is that the growth rate of the symmetrization $U_{x_{0}}^{*}$ over the ellipsoids $E_{r}\left(x_{0}\right)$ captures that of $U=U_{x_{0}}^{*}$ on the thin space $\Pi$ at $x_{0} \in \Gamma(U)$, while in the nonsymmetric case there could be a mismatch in those rates caused by the odd component of $U$, vanishing on $\Pi$.

More specifically, the growth rate of $U$ on $\Pi$ at $x_{0} \in \Gamma(U)$ is determined by the following quantity

$$
N^{A}\left(r, U_{x_{0}}^{*}, x_{0}\right):=\frac{r \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U_{x_{0}}^{*}, \nabla U_{x_{0}}^{*}\right\rangle}{\int_{\partial E_{r}\left(x_{0}\right)}\left(U_{x_{0}}^{*}\right)^{2} \mu_{x_{0}}\left(x-x_{0}\right)},
$$

which is a version of Almgren's frequency functional [47] written in the geometric terms determined by $A\left(x_{0}\right)$, where $\mu_{x_{0}}(z)=\frac{\left|A^{-1 / 2}\left(x_{0}\right) z\right|}{\left|A^{-1}\left(x_{0}\right) z\right|}$ is the conformal factor. As in the constant coefficient case, this quantity is of paramount importance for the classification of free boundary points.

Theorem G (Monotonicity of the truncated frequency). Let $U$ be as in Theorem $F$ and assume additionally that $U$ is $A$-quasisymmetric with respect to $\Pi$. Then for any $\kappa_{0} \geq 2$, there is $b=b\left(n, \alpha, M, \kappa_{0}\right)$ such that the truncated frequency

$$
r \mapsto \widehat{N}_{\kappa_{0}}^{A}\left(r, U_{x_{0}}^{*}, x_{0}\right):=\min \left\{\frac{1}{1-b r^{\alpha}} N^{A}\left(r, U_{x_{0}}^{*}, x_{0}\right), \kappa_{0}\right\}
$$

is monotone increasing for $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(u)$, and $0<r<r_{0}\left(n, \alpha, M, \kappa_{0}\right)$. Moreover, if we define

$$
\kappa\left(x_{0}\right):=\widehat{N}_{\kappa_{0}}^{A}\left(0+, U_{x_{0}}^{*}, x_{0}\right),
$$

the frequency of $U$ at $x_{0}$, then we have that either

$$
\kappa\left(x_{0}\right)=3 / 2 \quad \text { or } \quad \kappa\left(x_{0}\right) \geq 2
$$

The monotonicity of the truncated frequency follows from that of an one-parametric family of so-called Weiss-type energy functionals $\left\{W_{\kappa}^{A}\right\}_{0<\kappa<\kappa_{0}}$, see Section 2.7, which also play a fundamental role in the analysis of the free boundary.

The theorem above gives the following decomposition of the free boundary

$$
\Gamma(U)=\Gamma_{3 / 2}(U) \cup \bigcup_{\kappa \geq 2} \Gamma_{\kappa}(U)
$$

where

$$
\Gamma_{\kappa}(U):=\left\{x_{0} \in \Gamma(U): \kappa\left(x_{0}\right)=\kappa\right\} .
$$

The set $\Gamma_{3 / 2}(U)$, where the frequency is minimal is known as the regular set and is also denoted $\mathcal{R}(U)$.

Theorem H (Regularity of the regular set). Let $U$ be as in Theorem $G$. Then $\mathcal{R}(U)$ is a relatively open subset of the free boundary $\Gamma(U)$ and is an $(n-2)$-dimensional manifold of class $C^{1, \gamma}$.

Finally, we state our main result for the so-called singular set. A free boundary point $x_{0} \in \Gamma(U)$ is called singular if the coincidence set $\Lambda(U):=\left\{x \in B_{1}^{\prime}: U(x)=0\right\}$ has $H^{n-1}$-density zero at $x_{0}$, i.e.,

$$
\lim _{r \rightarrow 0+} \frac{H^{n-1}\left(\Lambda(U) \cap B_{r}^{\prime}\left(x_{0}\right)\right)}{H^{n-1}\left(B_{r}^{\prime}\right)}=0
$$

We denote the set of all singular points by $\Sigma(U)$ and call it the singular set. It can be shown that if $\kappa\left(x_{0}\right)<\kappa_{0}$, then $x_{0} \in \Sigma(U)$ if and only if $\kappa\left(x_{0}\right)=2 m, m \in \mathbb{N}$ (see Proposition 2.12.1). For such values of $\kappa$, we then define

$$
\Sigma_{\kappa}(U):=\Gamma_{\kappa}(U) .
$$

Theorem I (Structure of the singular set). Let $U$ be as in Theorem G. Then, for any $\kappa=2 m<\kappa_{0}, m \in \mathbb{N}, \Sigma_{\kappa}(U)$ is contained in a countable union of $(n-2)$-dimensional manifolds of class $C^{1, \log }$.

A more refined version of this result is given in Theorem 2.12.4.
Theorems H and I follow by establishing the uniqueness and continuous dependence of almost homogeneous blowups with Hölder modulus of continuity in the case of regular free boundary points and a logarithmic one in the case of the singular points. These follow from optimal growth and rotation estimates which are based on the use of Weiss-type monotonicity formulas in conjunction with so-called epiperimetric [20] and log-epiperimetric [27] inequalities for the solutions of the Signorini problem.

## Proofs of Theorems F-I

While we don't give formal proofs of the theorems above in the main body of the chapter, they are contained in the following results proved there:

- Theorem F is essentially the same as Theorem 2.5.1.
- Theorem G follows by combining Theorem 2.7.2 and Corollary 2.11.2.
- The statement of Theorem H is contained in that of Theorem 2.11.5.
- The statement of Theorem I is contained in that of Theorem 2.12.4.


### 2.2 Coordinate transformations

In order to use the results available for almost minimizers in the case of $A \equiv I$, proved in Chapter 1, in this section we describe a "deskewing procedure" or coordinate transformations to straighten $A\left(x_{0}\right), x_{0} \in B_{1}$.

For the notational convenience, we will denote

$$
\mathfrak{a}_{x_{0}}=A^{1 / 2}\left(x_{0}\right), \quad x_{0} \in B_{1}
$$

so that

$$
\left\langle A\left(x_{0}\right) \xi, \xi\right\rangle=\left|\mathfrak{a}_{x_{0}} \xi\right|^{2}, \quad \xi \in \mathbb{R}^{n}
$$

Then $\mathfrak{a}_{x_{0}}$ is a symmetric positive definite matrix, with eigenvalues between $\lambda^{1 / 2}$ and $\Lambda^{1 / 2}$ and the mapping $x_{0} \mapsto \mathfrak{a}_{x_{0}}$ is $\alpha$-Hölder continuous for $x_{0} \in B_{1}$. For every $x_{0} \in B_{1}$, we define an affine transformation $T_{x_{0}}$ by

$$
T_{x_{0}}(x)=\mathfrak{a}_{x_{0}}^{-1}\left(x-x_{0}\right)
$$

Note that $T_{x_{0}}^{-1}(y)=\mathfrak{a}_{x_{0}} y+x_{0}$. Then for the ellipsoids $E_{r}\left(x_{0}\right)$, we have

$$
E_{r}\left(x_{0}\right)=T_{x_{0}}^{-1}\left(B_{r}\right)=\mathfrak{a}_{x_{0}} B_{r}+x_{0}, \quad T_{x_{0}}\left(E_{r}\left(x_{0}\right)\right)=B_{r} .
$$

Further, we let

$$
\Pi_{x_{0}}:=T_{x_{0}}(\Pi) .
$$

Then $\Pi_{x_{0}}$ is a hyperplane parallel to a linear subspace $\mathfrak{a}_{x_{0}}^{-1} \Pi$ spanned by the vectors $\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{1}$, $\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{2}, \ldots, \mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{n-1}$ and with a normal $\mathfrak{a}_{x_{0}} \mathrm{e}_{n}$. Generally, this hyperplane will be tilted with respect to $\Pi$, unless $\mathfrak{a}_{x_{0}} \mathrm{e}_{n}$ is a multiple of $\mathrm{e}_{n}$, or equivalently that $\mathrm{e}_{n}$ is an eigenvector of the matrix $A\left(x_{0}\right)$, or that $a_{i n}\left(x_{0}\right)=0$ for $i=1, \ldots, n-1$ for its entries. To rectify that, we construct a family of orthogonal transformations $O_{x_{0}}, x_{0} \in B_{1}$, by applying the GramSchmidt process to the ordered basis $\left\{\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{1}, \mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{2}, \ldots, \mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{n-1}\right\}$ of $\mathfrak{a}_{x_{0}}^{-1} \Pi$. Namely, let

$$
\begin{aligned}
\mathrm{e}_{1}^{x_{0}} & :=\frac{\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{1}}{\left|\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{1}\right|}, \\
\mathrm{e}_{2}^{x_{0}} & :=\frac{\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{2}-\left\langle\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{2}, \mathrm{e}_{1}^{x_{0}}\right\rangle \mathrm{e}_{1}^{x_{0}}}{\left|\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{2}-\left\langle\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{2}, \mathrm{e}_{1}^{x_{0}}\right\rangle \mathrm{e}_{1}^{x_{0}}\right|}, \\
\mathrm{e}_{3}^{x_{0}} & :=\frac{\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{3}-\left\langle\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{3}, \mathrm{e}_{1}^{x_{0}}\right\rangle \mathrm{e}_{1}^{x_{0}}-\left\langle\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{3}, \mathrm{e}_{2}^{x_{0}}\right\rangle \mathrm{e}_{2}^{x_{0}}}{\left|\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{3}-\left\langle\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{3}, \mathrm{e}_{1}^{x_{0}}\right\rangle \mathrm{e}_{1}^{x_{0}}-\left\langle\mathfrak{a}_{x_{0}}^{-1} \mathrm{e}_{3}, \mathrm{e}_{2}^{x_{0}}\right\rangle \mathrm{e}_{2}^{x_{0}}\right|} \\
\quad &
\end{aligned}
$$

Moreover, letting

$$
\mathrm{e}_{n}^{x_{0}}:=\frac{\mathfrak{a}_{x_{0}} \mathrm{e}_{n}}{\left|\mathfrak{a}_{x_{0}} \mathrm{e}_{n}\right|},
$$

we obtain an ordered orthonormal basis $\left\{\mathrm{e}_{1}^{x_{0}}, \ldots, \mathrm{e}_{n-1}^{x_{0}}, \mathrm{e}_{n}^{x_{0}}\right\}$ of $\mathbb{R}^{n}$. Then consider the rotation $O_{x_{0}}$ of $\mathbb{R}^{n}$ that takes the standard basis $\left\{\mathrm{e}_{1}, \mathrm{e}_{2}, \ldots, \mathrm{e}_{n}\right\}$ to the one above, i.e.,

$$
O_{x_{0}}: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}, \quad O_{x_{0}}\left(\mathrm{e}_{i}\right)=\mathrm{e}_{i}^{x_{0}}, \quad i=1,2, \ldots, n
$$

Note that the Gram-Schmidt process above guarantees that $x_{0} \mapsto O_{x_{0}}$ is $\alpha$-Hölder continuous. We also have that by construction

$$
O_{x_{0}}^{-1} \mathfrak{a}_{x_{0}}^{-1} \Pi=\Pi .
$$

In particular, when $x_{0} \in \Pi$, we have $\Pi_{x_{0}}=\mathfrak{a}_{x_{0}}^{-1} \Pi$ and therefore

$$
O_{x_{0}}^{-1}\left(\Pi_{x_{0}}\right)=\Pi .
$$

Because of this property, we also define the modifications of the matrices $\mathfrak{a}_{x_{0}}$ and the transformations $T_{x_{0}}$ as follows:

$$
\overline{\mathfrak{a}}_{x_{0}}=\mathfrak{a}_{x_{0}} O_{x_{0}}, \quad \bar{T}_{x_{0}}=O_{x_{0}}^{-1} \circ T_{x_{0}}
$$

so that $\bar{T}_{x_{0}}(x)=\overline{\mathfrak{a}}_{x_{0}}^{-1}\left(x-x_{0}\right)$. Since $O_{x_{0}}$ is a rotation, we still have

$$
E_{r}\left(x_{0}\right)=\bar{T}_{x_{0}}^{-1}\left(B_{r}\right), \quad \bar{T}_{x_{0}}\left(E_{r}\left(x_{0}\right)\right)=B_{r},
$$

see Fig. 2.1.
Next, for a function $U: B_{1} \rightarrow \mathbb{R}$ and a point $x_{0} \in B_{1}$, we define its "deskewed" version at $x_{0}$ by

$$
u_{x_{0}}=U \circ \bar{T}_{x_{0}}^{-1} .
$$



Figure 2.1. Deskewing: coordinate transformations $T_{x_{0}}, O_{x_{0}}^{-1}, \bar{T}_{x_{0}}$.

As we will see, if $U$ is an almost minimizer, the transformed function $u_{x_{0}}$ will satisfy an almost minimizing property with the identity matrix $I$ at the origin. Before we state and prove that fact, we need the following basic change of variable formulas:

$$
\begin{align*}
\int_{E_{r}\left(x_{0}\right)} U^{2} & =\operatorname{det} \mathfrak{a}_{x_{0}} \int_{B_{r}} u_{x_{0}}^{2}  \tag{2.2.1}\\
\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle & =\operatorname{det} \mathfrak{a}_{x_{0}} \int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2}  \tag{2.2.2}\\
\int_{\partial E_{r}\left(x_{0}\right)} U^{2} \mu_{x_{0}}\left(x-x_{0}\right) & =\operatorname{det} \mathfrak{a}_{x_{0}} \int_{\partial B_{r}} u_{x_{0}}^{2}, \tag{2.2.3}
\end{align*}
$$

with the conformal factor

$$
\begin{equation*}
\mu_{x_{0}}(z):=\frac{\left|\mathfrak{a}_{x_{0}}^{-1} z\right|}{\left|A^{-1}\left(x_{0}\right) z\right|} . \tag{2.2.4}
\end{equation*}
$$

We also have the following modified version of (2.2.2).

$$
\begin{equation*}
\int_{E_{r}\left(x_{0}\right)}\left|\mathfrak{a}_{x_{0}} \nabla U-\left\langle\mathfrak{a}_{x_{0}} \nabla U\right\rangle_{E_{r}\left(x_{0}\right)}\right|^{2}=\operatorname{det} \mathfrak{a}_{x_{0}} \int_{B_{r}}\left|\nabla u_{x_{0}}-\left\langle\nabla u_{x_{0}}\right\rangle_{B_{r}}\right|^{2} . \tag{2.2.5}
\end{equation*}
$$

While (2.2.1)-(2.2.2) and (2.2.5) are clear, let us give more details on (2.2.3). If we let $f(x):=\left|\mathfrak{a}_{x_{0}}^{-1}\left(x-x_{0}\right)\right|$, then $\{f=t\}=\partial E_{t}\left(x_{0}\right), t>0$, and by the coarea formula

$$
\int_{E_{r}\left(x_{0}\right)} U^{2} d x=\int_{0}^{r} \int_{\partial E_{t}\left(x_{0}\right)} \frac{U^{2}}{|\nabla f(x)|} d S_{x} d t .
$$

Using now that $1 /|\nabla f(x)|=\frac{\left|\mathfrak{a}_{x_{0}}^{-1}\left(x-x_{0}\right)\right|}{\left|A^{-1}\left(x_{0}\right)\left(x-x_{0}\right)\right|}=\mu_{x_{0}}\left(x-x_{0}\right)$ and then differentiating (2.2.1), we obtain (2.2.3).

We will also need the following estimate for the conformal factor $\mu_{x_{0}}$ :

$$
\begin{equation*}
\lambda^{1 / 2} \leq \mu_{x_{0}}(z) \leq \Lambda^{1 / 2} \tag{2.2.6}
\end{equation*}
$$

Indeed, if $y=A^{-1}\left(x_{0}\right) z$, then

$$
\mu_{x_{0}}(z)=\frac{\left|\mathfrak{a}_{x_{0}} y\right|}{|y|} \in\left[\lambda^{1 / 2}, \Lambda^{1 / 2}\right]
$$

Definition 2.2.1 (Almost Signorini property at a point). We say that a function $u \in$ $W^{1,2}\left(B_{R}\right)$ satisfies the almost Signorini property at 0 in $B_{R}$ if

$$
\int_{B_{r}}|\nabla u|^{2} \leq(1+\omega(r)) \int_{B_{r}}|\nabla v|^{2}
$$

for all $0<r<R$ and $v \in \mathfrak{K}_{0, u}\left(B_{r}, \Pi\right)$.
Lemma 2.2.1. Suppose $U$ is an almost minimizer of the $A$-Signorini problem in $B_{1}$. Let $x_{0} \in B_{1}^{\prime}$ be such that $E_{R}\left(x_{0}\right) \subset B_{1}$. Then $u_{x_{0}}=U \circ \bar{T}_{x_{0}}^{-1}$ satisfies the almost Signorini property at 0 in $B_{R}$.

Proof. Let $V$ be the energy minimizer of $\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla V, \nabla V\right\rangle$ on $\mathfrak{K}_{0, U}\left(E_{r}\left(x_{0}\right), \Pi\right), 0<r<$ $R$. Then $v_{x_{0}}=V \circ \bar{T}_{x_{0}}^{-1}$ is the energy minimizer of $\int_{B_{r}}\left|\nabla v_{x_{0}}\right|^{2}$ on $\mathfrak{K}_{0, u_{x_{0}}}\left(B_{r}, \Pi\right)$. Moreover, by (2.2.2),

$$
\begin{aligned}
\int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2} & =\operatorname{det} \mathfrak{a}_{x_{0}}^{-1} \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle \\
& \leq(1+\omega(r)) \operatorname{det} \mathfrak{a}_{x_{0}}^{-1} \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla V, \nabla V\right\rangle \\
& =(1+\omega(r)) \int_{B_{r}}\left|\nabla v_{x_{0}}\right|^{2} .
\end{aligned}
$$

This completes the proof.

### 2.3 Almost $A$-harmonic functions

We start our analysis of almost minimizers in the absence of the thin obstacle. We call such functions almost $A$-harmonic functions. In this section, we establish their $C^{1, \alpha / 2}$ regularity (Theorem 2.3.2). A similar result has already been proved by Anzellotti [31], but for almost minimizers over balls $\left\{B_{r}\left(x_{0}\right)\right\}$ rather than ellipsoids $\left\{E_{r}\left(x_{0}\right)\right\}$; nevertheless, the proofs are similar. The proofs in this section also illustrate how we are going to use the results available for "deskewed" functions $u_{x_{0}}=U \circ \bar{T}_{x_{0}}^{-1}$ to infer the corresponding results for almost minimizers $U$.

Definition 2.3.1 (Almost $A$-harmonic functions). We say that $U$ is an almost $A$-harmonic function in $B_{1}$ if $U \in W^{1,2}\left(B_{1}\right)$ and

$$
\int_{E_{r}\left(x_{0}\right)}\langle A \nabla U, \nabla U\rangle \leq(1+\omega(r)) \int_{E_{r}\left(x_{0}\right)}\langle A \nabla V, \nabla V\rangle,
$$

whenever $E_{r}\left(x_{0}\right) \Subset B_{1}$ and $V \in \mathfrak{K}_{U}\left(E_{r}\left(x_{0}\right)\right):=U+W_{0}^{1,2}\left(E_{r}\left(x_{0}\right)\right)$.

Note that similarly to the case of $A$-Signorini problem, we can write the almost minimizing property above in the form with frozen coefficients

$$
\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle \leq(1+\omega(r)) \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla V, \nabla V\right\rangle .
$$

Definition 2.3.2 (Almost harmonic property at a point). We say that a function $u \in$ $W^{1,2}\left(B_{R}\right)$ satisfies almost harmonic property at 0 in $B_{R}$ if

$$
\int_{B_{r}}|\nabla u|^{2} \leq(1+\omega(r)) \int_{B_{r}}|\nabla v|^{2},
$$

for all $0<r<R$ and $v \in \mathfrak{K}_{u}\left(B_{r}\right)$.

Lemma 2.3.1. If $U$ is an almost $A$-harmonic function in $B_{1}$ and $x_{0} \in B_{1}$ with $E_{R}\left(x_{0}\right) \subset B_{1}$, then $u_{x_{0}}$ satisfies the almost harmonic property at 0 in $B_{R}$.

Proof. The proof is similar to that of Lemma 2.2.1.

Proposition 2.3.1 (cf.Proposition 1.2.2). Let $U$ be an almost $A$-harmonic function in $B_{1}$. Then for any $B_{r}\left(x_{0}\right) \Subset B_{1}$ and $0<\rho<r$, we have

$$
\begin{align*}
& \int_{B_{\rho}\left(x_{0}\right)}|\nabla U|^{2} \leq C\left[\left(\frac{\rho}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2},  \tag{2.3.1}\\
& \int_{B_{\rho}\left(x_{0}\right)}\left|\nabla U-\langle\nabla U\rangle_{B_{\rho}\left(x_{0}\right)}\right|^{2} \leq C\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}\left(x_{0}\right)}\left|\nabla U-\langle\nabla U\rangle_{B_{r}\left(x_{0}\right)}\right|^{2}  \tag{2.3.2}\\
&+C r^{\alpha} \int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2},
\end{align*}
$$

with $C=C(n, \alpha, M)$.

Proof. Since $u_{x_{0}}$ satisfies the almost harmonic property at 0 , if $h$ is the harmonic replacement of $u_{x_{0}}$ in $B_{r}$ (i.e., $h$ is harmonic in $B_{r}$ with $h=u_{x_{0}}$ on $\partial B_{r}$ ), then

$$
\int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2} \leq\left(1+M r^{\alpha}\right) \int_{B_{r}}|\nabla h|^{2} .
$$

This is enough to repeat the arguments in Proposition 1.2.2, to obtain

$$
\begin{gathered}
\int_{B_{\rho}}\left|\nabla u_{x_{0}}\right|^{2} \leq 2\left[\left(\frac{\rho}{r}\right)^{n}+M r^{\alpha}\right] \int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2}, \\
\int_{B_{\rho}}\left|\nabla u_{x_{0}}-\left\langle\nabla u_{x_{0}}\right\rangle_{B_{\rho}}\right|^{2} \leq 9\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\nabla u_{x_{0}}-\left\langle\nabla u_{x_{0}}\right\rangle_{B_{r}}\right|^{2}+24 M r^{\alpha} \int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2} .
\end{gathered}
$$

Then, by the change of variables formulas (2.2.2) and (2.2.5), we have

$$
\begin{align*}
& \int_{E_{\rho}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle \leq 2\left[\left(\frac{\rho}{r}\right)^{n}+M r^{\alpha}\right] \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle,  \tag{2.3.3}\\
& \int_{E_{\rho}\left(x_{0}\right)}\left|\mathfrak{a}_{x_{0}} \nabla U-\left\langle\mathfrak{a}_{x_{0}} \nabla U\right\rangle_{E_{\rho}\left(x_{0}\right)}\right|^{2} \leq 9\left(\frac{\rho}{r}\right)^{n+2} \int_{E_{r}\left(x_{0}\right)}\left|\mathfrak{a}_{x_{0}} \nabla U-\left\langle\mathfrak{a}_{x_{0}} \nabla U\right\rangle_{E_{r}\left(x_{0}\right)}\right|^{2}  \tag{2.3.4}\\
& +24 M r^{\alpha} \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle .
\end{align*}
$$

To show now that (2.3.3)-(2.3.4) imply (2.3.1)-(2.3.2), we first consider the case

$$
0<\rho<(\lambda / \Lambda)^{1 / 2} r
$$

Then, using the inclusions

$$
B_{\rho}\left(x_{0}\right) \subset E_{\lambda^{-1 / 2} \rho}\left(x_{0}\right) \subset E_{\Lambda^{-1 / 2} r}\left(x_{0}\right) \subset B_{r}\left(x_{0}\right)
$$

applying (2.3.3)-(2.3.4) with $\lambda^{-1 / 2} \rho$ and $\Lambda^{-1 / 2} r$ in place of $\rho$ and $r$, and using the ellipticity of $A\left(x_{0}\right)$, we obtain (2.3.1)-(2.3.2) in this case.

In the remaining case

$$
(\lambda / \Lambda)^{1 / 2} r \leq \rho \leq r
$$

the inequalities (2.3.1)-(2.3.2) hold readily, as

$$
\begin{aligned}
\int_{B_{\rho}\left(x_{0}\right)}|\nabla U|^{2} & \leq\left(\frac{\Lambda}{\lambda}\right)^{n / 2}\left(\frac{\rho}{r}\right)^{n} \int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2}, \\
\int_{B_{\rho}\left(x_{0}\right)}\left|\nabla U-\langle\nabla U\rangle_{B_{\rho}\left(x_{0}\right)}\right|^{2} & \leq \int_{B_{\rho}\left(x_{0}\right)}\left|\nabla U-\langle\nabla U\rangle_{B_{r}\left(x_{0}\right)}\right|^{2} \\
& \leq\left(\frac{\Lambda}{\lambda}\right)^{\frac{n+2}{2}}\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}\left(x_{0}\right)}\left|\nabla U-\langle\nabla U\rangle_{B_{r}\left(x_{0}\right)}\right|^{2} .
\end{aligned}
$$

Theorem 2.3.2. Let $U$ be an almost A-harmonic function in $B_{1}$. Then $U \in C^{1, \alpha / 2}\left(B_{1}\right)$ with

$$
\|U\|_{C^{1, \alpha / 2}(K)} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}
$$

for any $K \Subset B_{1}$, with $C=C(n, \alpha, M, K)$.
Proof. Let $K \Subset B_{1}$ and $x_{0} \in \widetilde{K}:=\left\{y \in B_{1}: \operatorname{dist}\left(y, \partial B_{1}\right) \geq r_{0}\right\}$, where $r_{0}=\frac{1}{2} \operatorname{dist}\left(K, \partial B_{1}\right)$. For in $(0,1)$, a direct application of Lemma 1.2.2 to (2.3.1) gives

$$
\int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2} \leq C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n-2+2 \sigma}
$$

for any $0<r<r_{0}$, with $C$ depending on $n, \alpha, \sigma, M, K$. Combining this with (2.3.2) also gives,

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\nabla U-\langle\nabla U\rangle_{B_{\rho}\left(x_{0}\right)}\right|^{2} \leq C\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}\left(x_{0}\right)}\left|\nabla U-\langle\nabla U\rangle_{B_{r}\left(x_{0}\right)}\right|^{2}
$$

$$
\begin{equation*}
+C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n-2+2 \sigma+\alpha} \tag{2.3.5}
\end{equation*}
$$

for any $0<\rho<r<r_{0}$. If we take $\sigma \in(0,1)$ such that $\alpha^{\prime}:=\frac{-2+2 \sigma+\alpha}{2}>0$, then Lemma 1.2.2 produces

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\nabla U-\langle\nabla U\rangle_{B_{\rho}\left(x_{0}\right)}\right|^{2} \leq C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} \rho^{n+2 \alpha^{\prime}}
$$

and this readily implies $\nabla U \in C^{0, \alpha^{\prime}}(\widetilde{K})$. Now we know that $\nabla U$ is bounded in $\widetilde{K}$, and thus $\int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2} \leq C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n}$. Plugging this in the last term of (2.3.2) and repeating the arguments above, we conclude that $U \in C^{1, \alpha / 2}$.

### 2.4 Almost Lipschitz regularity of almost minimizers

In this section, we make the first step towards the regularity of almost minimizers for the $A$-Signorini problem and show that they are almost Lipschitz, i.e., $C^{0, \sigma}$ for every $0<\sigma<1$ (Theorem 2.4.1). The proof is based on the Morrey space embedding, similar to the case of almost $A$-harmonic functions, as well as the case of almost minimizers with $A=I$, treated in Chapter 1. We want to emphasize, however, that the results on almost Lipschitz and $C^{1, \beta}$ regularity of almost minimizers (in the next section) do not require any symmetry condition that was imposed in Chapter 1.

We start with an auxiliary result on the solutions of the Signorini problem.

Proposition 2.4.1. Let $h$ be a solution of the Signorini problem in $B_{1}$. Then

$$
\begin{equation*}
\int_{B_{\rho}}|\nabla h|^{2} \leq\left(\frac{\rho}{R}\right)^{n} \int_{B_{R}}|\nabla h|^{2}, \quad 0<\rho<R<1 \tag{2.4.1}
\end{equation*}
$$

Proof. The difference of this proposition from Proposition 1.3.1 is that $h(y)$ is not assumed to be even symmetric in $y_{n}$-variable. To circumvent that, we decompose $h$ into the sum of even and odd functions in $y_{n}$, i.e.,

$$
\begin{align*}
h\left(y^{\prime}, y_{n}\right) & =\frac{h\left(y^{\prime}, y_{n}\right)+h\left(y^{\prime},-y_{n}\right)}{2}+\frac{h\left(y^{\prime}, y_{n}\right)-h\left(y^{\prime},-y_{n}\right)}{2}  \tag{2.4.2}\\
& =: h^{*}\left(y^{\prime}, y_{n}\right)+h^{\sharp}\left(y^{\prime}, y_{n}\right) .
\end{align*}
$$

It is easy to see that $h^{*}$ is a solution of the Signorini problem, even in $y_{n}$-variable, and $h^{\sharp}$ is a harmonic function, odd in $y_{n}$-variable.

Then both $\left|\nabla h^{*}\right|^{2}$ and $\left|\nabla h^{\sharp}\right|^{2}$ are subharmonic functions in $B_{1}$ (see Proposition 1.3.1 for $h^{*}$ ), which implies that for $0<\rho<R<1$

$$
\begin{aligned}
\int_{B_{\rho}}\left|\nabla h^{*}\right|^{2} & \leq\left(\frac{\rho}{R}\right)^{n} \int_{B_{R}}\left|\nabla h^{*}\right|^{2}, \\
\int_{B_{\rho}}\left|\nabla h^{\sharp}\right|^{2} & \leq\left(\frac{\rho}{R}\right)^{n} \int_{B_{R}}\left|\nabla h^{\sharp}\right|^{2} .
\end{aligned}
$$

Now observing that $\int_{B_{t}}|\nabla h|^{2}=\int_{B_{t}}\left(\left|\nabla h^{*}\right|^{2}+\left|\nabla h^{\sharp}\right|^{2}\right)$, for $0<t \leq R$, we obtain (2.4.1).
Proposition 2.4.2 (cf. Proposition 1.3.2). Let $U$ be an almost minimizer for the $A$-Signorini problem in $B_{1}$, and $B_{R}\left(x_{0}\right) \Subset B_{1}$. Then, there is $C_{1}=C_{1}(n, M)>1$ such that

$$
\begin{equation*}
\int_{B_{\rho}\left(x_{0}\right)}|\nabla U|^{2} \leq C_{1}\left[\left(\frac{\rho}{R}\right)^{n}+R^{\alpha}\right] \int_{B_{R}\left(x_{0}\right)}|\nabla U|^{2}, \quad 0<\rho<R . \tag{2.4.3}
\end{equation*}
$$

Proof. Case 1. Suppose $x_{0} \in B_{1}^{\prime}$. Note that $u_{x_{0}}$ satisfies the Signorini property at 0 in $B_{r}$ with $r=\Lambda^{-1 / 2} R$. If $h$ is the Signorini replacement of $u_{x_{0}}$ in $B_{r}$ (that is, $h$ solves the Signorini problem in $B_{r}$ with thin obstacle 0 on $\Pi$ and boundary values $h=u_{x_{0}}$ on $\partial B_{r}$ ), then $h$ satisfies

$$
\int_{B_{r}}\langle\nabla h, \nabla(v-h)\rangle \geq 0,
$$

for any $v \in \mathfrak{K}_{0, u_{x_{0}}}\left(B_{r}, \Pi\right)$, which easily follows from the standard first variation argument. Plugging in $v=u_{x_{0}}$, we obtain

$$
\int_{B_{r}}\left\langle\nabla h, \nabla u_{x_{0}}\right\rangle \geq \int_{B_{r}}|\nabla h|^{2} .
$$

Then it follows that

$$
\begin{aligned}
\int_{B_{r}}\left|\nabla\left(u_{x_{0}}-h\right)\right|^{2} & =\int_{B_{r}}\left(\left|\nabla u_{x_{0}}\right|^{2}+|\nabla h|^{2}-2\left\langle\nabla u_{x_{0}}, \nabla h\right\rangle\right) \leq \int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2}-\int_{B_{r}}|\nabla h|^{2} \\
& \leq\left(1+M r^{\alpha}\right) \int_{B_{r}}|\nabla h|^{2}-\int_{B_{r}}|\nabla h|^{2} \leq M r^{\alpha} \int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2},
\end{aligned}
$$

where in the last inequality we have used that $h$ is the energy minimizer of the Dirichlet integral in $\mathfrak{K}_{0, u_{x_{0}}}\left(B_{r}, \Pi\right)$. Then, for $\rho \leq r$, we have

$$
\begin{aligned}
\int_{B_{\rho}}\left|\nabla u_{x_{0}}\right|^{2} & \leq 2 \int_{B_{\rho}}|\nabla h|^{2}+2 \int_{B_{\rho}}\left|\nabla\left(u_{x_{0}}-h\right)\right|^{2} \leq 2\left(\frac{\rho}{r}\right)^{n} \int_{B_{r}}|\nabla h|^{2}+2 M r^{\alpha} \int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2} \\
& \leq C\left[\left(\frac{\rho}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2} .
\end{aligned}
$$

Now, we transform back from $u_{x_{0}}$ to $U$ as we did in Proposition 2.3.1 to obtain (2.4.3) in this case.

Case 2. Now consider the case $x_{0} \in B_{1}^{+}$. If $\rho \geq r / 4$, then we simply have

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla U|^{2} \leq 4^{n}\left(\frac{\rho}{r}\right)^{n} \int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2}
$$

Thus, we may assume $\rho<r / 4$. Then, let $d:=\operatorname{dist}\left(x_{0}, B_{1}^{\prime}\right)>0$ and choose $x_{1} \in \partial B_{d}\left(x_{0}\right) \cap B_{1}^{\prime}$.
Case 2.1. If $\rho \geq d$, then we use $B_{\rho}\left(x_{0}\right) \subset B_{2 \rho}\left(x_{1}\right) \subset B_{r / 2}\left(x_{1}\right) \subset B_{r}\left(x_{0}\right)$ and the result of Case 1 to write

$$
\begin{aligned}
\int_{B_{\rho}\left(x_{0}\right)}|\nabla U|^{2} & \leq \int_{B_{2 \rho}\left(x_{1}\right)}|\nabla U|^{2} \leq C\left[\left(\frac{2 \rho}{r / 2}\right)^{n}+(r / 2)^{\alpha}\right] \int_{B_{r / 2}\left(x_{1}\right)}|\nabla U|^{2} \\
& \leq C\left[\left(\frac{\rho}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2}
\end{aligned}
$$

Case 2.2. Suppose now $d>\rho$. If $d>r$, then $B_{r}\left(x_{0}\right) \Subset B_{1}^{+}$. Since $U$ is almost harmonic in $B_{1}^{+}$, we can apply Proposition 2.3.1 to obtain

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla U|^{2} \leq C\left[\left(\frac{\rho}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2}
$$

Thus, we may assume $d \leq r$. Then we note that $B_{d}\left(x_{0}\right) \subset B_{1}^{+}$and by a limiting argument from the previous estimate, we obtain

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla U|^{2} \leq C\left[\left(\frac{\rho}{d}\right)^{n}+r^{\alpha}\right] \int_{B_{d}\left(x_{0}\right)}|\nabla U|^{2} .
$$

To estimate $\int_{B_{d}\left(x_{0}\right)}|\nabla U|^{2}$ in the right-hand side of the above inequality, we further consider the two subcases.

Case 2.2.1. If $r / 4 \leq d$, then

$$
\int_{B_{d}\left(x_{0}\right)}|\nabla U|^{2} \leq 4^{n}\left(\frac{d}{r}\right)^{n} \int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2}
$$

which immediately implies (2.4.3).
Case 2.2.2. It remains to consider the case $\rho<d<r / 4$. Using Case 1 again, we have

$$
\begin{aligned}
\int_{B_{d}\left(x_{0}\right)}|\nabla U|^{2} & \leq \int_{B_{2 d}\left(x_{1}\right)}|\nabla U|^{2} \leq C\left[\left(\frac{2 d}{r / 2}\right)^{n}+(r / 2)^{\alpha}\right] \int_{B_{r / 2}\left(x_{1}\right)}|\nabla U|^{2} \\
& \leq C\left[\left(\frac{d}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2}
\end{aligned}
$$

which also implies (2.4.3). This concludes the proof of the proposition.
As we have seen in Chapter 1, Proposition 2.4.2 implies the almost Lipschitz regularity of almost minimizers.

Theorem 2.4.1. Let $U$ be an almost minimizer for the $A$-Signorini problem in $B_{1}$. Then $U \in C^{0, \sigma}\left(B_{1}\right)$ for all $0<\sigma<1$. Moreover, for any $K \Subset B_{1}$,

$$
\|U\|_{C^{0, \sigma}(K)} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}
$$

with $C=C(n, \alpha, M, \sigma, K)$.

Proof. The proof is essentially identical to that of Theorem 1.3.1. Let $K \Subset B_{1}$ and $x_{0} \in$ $K$. Take $r_{0}=r_{0}(n, \alpha, M, \sigma, K)>0$ such that $r_{0}<\operatorname{dist}\left(K, \partial B_{1}\right)$ and $r_{0}^{\alpha} \leq \varepsilon\left(C_{1}, n, n+\right.$ $2 \sigma-2)$, where $\varepsilon=\varepsilon\left(C_{1}, n, n+2 \sigma-2\right)$ is as in Lemma 1.2.2 and $C_{1}=C_{1}(n, M)$ is as in Proposition 2.4.2. Then for all $0<\rho<r<r_{0}$, by Proposition 2.4.2,

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla U|^{2} \leq C_{1}\left[\left(\frac{\rho}{r}\right)^{n}+r^{\alpha}\right] \int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2} .
$$

By Lemma 1.2.2, we get

$$
\int_{B_{\rho}\left(x_{0}\right)}|\nabla U|^{2} \leq C(n, M, \sigma)\left(\frac{\rho}{r}\right)^{n+2 \sigma-2} \int_{B_{r}\left(x_{0}\right)}|\nabla U|^{2}
$$

Taking $r \nearrow r_{0}$, we conclude that

$$
\begin{equation*}
\int_{B_{\rho}\left(x_{0}\right)}|\nabla U|^{2} \leq C(n, \alpha, M, \sigma, K)\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} \rho^{n+2 \sigma-2} . \tag{2.4.4}
\end{equation*}
$$

By the Morrey space embedding (Corollary 3.2 in [50]), we obtain $U \in C^{0, \sigma}(K)$ with

$$
\begin{equation*}
\|U\|_{C^{0, \sigma}(K)} \leq C(n, \alpha, M, \sigma, K)\|U\|_{W^{1,2}\left(B_{1}\right)} \tag{2.4.5}
\end{equation*}
$$

## $2.5 C^{1, \beta}$ regularity of almost minimizers

In this section we prove $C^{1, \beta}$ regularity of the almost minimizers for the $A$-Signorini problem (Theorem 2.5.1). While we take advantage of the results available for the even symmetric almost minimizers with $A=I$ in Chapter 1 , removing the symmetry condition requires new additional steps, combined with "deskewing" arguments to generalize to the variable coefficient case.

We start again with an auxiliary result for the solutions of the Signorini problem.
Proposition 2.5.1. Let $h$ be a solution of the Signorini problem in $B_{r}, 0<r<1$. Define

$$
\widehat{\nabla h}:= \begin{cases}\nabla h\left(y^{\prime}, y_{n}\right), & y_{n} \geq 0 \\ \nabla h\left(y^{\prime},-y_{n}\right), & y_{n}<0\end{cases}
$$

the even extension of $\nabla h$ from $B_{r}^{+}$to $B_{r}$. Then for $0<\alpha<1$, there are $C_{1}=C_{1}(n, \alpha)$, $C_{2}=C_{2}(n, \alpha)$ such that for all $0<\rho \leq s \leq(3 / 4) r$,

$$
\begin{equation*}
\int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{\rho}}\right|^{2} \leq C_{1}\left(\frac{\rho}{s}\right)^{n+\alpha} \int_{B_{s}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{s}}\right|^{2}+C_{2} \frac{s^{n+1}}{r^{n+3}} \int_{B_{r}} h^{2} \tag{2.5.1}
\end{equation*}
$$

Proof. This proposition differs from Proposition 1.4.1 only by not requiring $h(y)$ to be even in the $y_{n}$-variable. As in the proof of Proposition 2.4 .1 we split $h$ into its even and odd parts

$$
h(y)=h^{*}(y)+h^{\sharp}(y), \quad y \in B_{r} .
$$

Recall that $h^{*}$ is still a solution of the Signorini problem in $B_{r}$, but now even in $y_{n}$ and $h^{\sharp}$ is a harmonic function in $B_{r}$, odd in $y_{n}$. Then, by Proposition 2.4.1 we have

$$
\begin{equation*}
\int_{B_{\rho}}\left|\widehat{\nabla h^{*}}-\left\langle\widehat{\nabla h^{*}}\right\rangle_{B_{\rho}}\right|^{2} \leq C_{1}\left(\frac{\rho}{s}\right)^{n+\alpha} \int_{B_{s}}\left|\widehat{\nabla h^{*}}-\left\langle\widehat{\nabla h^{*}}\right\rangle_{B_{s}}\right|^{2}+C_{2} \frac{s^{n+1}}{r^{n+3}} \int_{B_{r}}\left(h^{*}\right)^{2} \tag{2.5.2}
\end{equation*}
$$

Now we need a similar estimate for $h^{\sharp}$. Since $h^{\sharp}$ is harmonic, by the standard interior estimates, we have

$$
\sup _{B_{(3 / 4) r}}\left|D^{2} h^{\sharp}\right| \leq \frac{C(n)}{r^{2}}\left(\frac{1}{r^{n}} \int_{B_{r}}\left(h^{\sharp}\right)^{2}\right)^{1 / 2} .
$$

Thus, taking the averages on $B_{\rho}^{+}$, we will therefore have

$$
\begin{aligned}
\int_{B_{\rho}^{+}}\left|\nabla h^{\sharp}-\left\langle\nabla h^{\sharp}\right\rangle_{B_{\rho}^{+}}\right|^{2} & \leq C(n) \rho^{n+2}\left(\sup _{B_{\rho}}\left|D^{2} h^{\sharp}\right|\right)^{2} \leq C(n) \frac{\rho^{n+2}}{r^{n+4}} \int_{B_{r}}\left(h^{\sharp}\right)^{2} \\
& \leq C(n) \frac{s^{n+1}}{r^{n+3}} \int_{B_{r}}\left(h^{\sharp}\right)^{2}, \quad 0<\rho<s \leq(3 / 4) r,
\end{aligned}
$$

which can be rewritten as

$$
\begin{equation*}
\int_{B_{\rho}}\left|\widehat{\nabla h^{\sharp}}-\left\langle\widehat{\nabla h^{\sharp}}\right\rangle_{B_{\rho}}\right|^{2} \leq C(n) \frac{s^{n+1}}{r^{n+3}} \int_{B_{r}}\left(h^{\sharp}\right)^{2} . \tag{2.5.3}
\end{equation*}
$$

Now using that $\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{\rho}}=\left[\widehat{\nabla h^{*}}-\left\langle\widehat{\nabla h^{*}}\right\rangle_{B_{\rho}}\right]+\left[\widehat{\nabla h^{\sharp}}-\left\langle\widehat{\nabla h^{\sharp}}\right\rangle_{B_{\rho}}\right]$ in $B_{\rho}$, we deduce from (2.5.3) that

$$
\begin{align*}
\int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{\rho}}\right|^{2} & \leq 2 \int_{B_{\rho}}\left|\widehat{\nabla h^{*}}-\left\langle\widehat{\nabla h^{*}}\right\rangle_{B_{\rho}}\right|^{2}+2 \int_{B_{\rho}} \mid \widehat{\nabla h^{\sharp}}-\left\langle\left.\widehat{\left.\nabla h^{\sharp}\right\rangle_{B_{\rho}}}\right|^{2}\right.  \tag{2.5.4}\\
& \leq 2 \int_{B_{\rho}}\left|\widehat{\nabla h^{*}}-\left\langle\widehat{\nabla h^{*}}\right\rangle_{B_{\rho}}\right|^{2}+C(n) \frac{s^{n+1}}{r^{n+3}} \int_{B_{r}}\left(h^{\sharp}\right)^{2} .
\end{align*}
$$

Similarly, representing $\widehat{\nabla h^{*}}-\left\langle\widehat{\nabla h^{*}}\right\rangle_{B_{s}}=\left[\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{s}}\right]-\left[\widehat{\nabla h^{\sharp}}-\left\langle\widehat{\left.\nabla h^{\sharp}\right\rangle_{B_{s}}}\right]\right.$ in $B_{s}$, we deduce from (2.5.3) (by taking $\rho=s$ ) that

$$
\begin{equation*}
\int_{B_{s}}\left|\widehat{\nabla h^{*}}-\left\langle\widehat{\nabla h^{*}}\right\rangle_{B_{s}}\right|^{2} \leq 2 \int_{B_{s}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{s}}\right|^{2}+C(n) \frac{s^{n+1}}{r^{n+3}} \int_{B_{r}}\left(h^{\sharp}\right)^{2} . \tag{2.5.5}
\end{equation*}
$$

Hence, combining (2.5.2)-(2.5.5), and using that both $\int_{B_{r}}\left(h^{*}\right)^{2}$ and $\int_{B_{r}}\left(h^{\sharp}\right)^{2}$ cannot exceed $\int_{B_{r}} h^{2}$, we obtain the claimed estimate (2.5.1).

Theorem 2.5.1. Let $U$ be an almost minimizer of the $A$-Signorini problem in $B_{1}$. Then

$$
U \in C^{1, \beta}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right) \quad \text { with } \beta=\frac{\alpha}{4(2 n+\alpha)} .
$$

Moreover, for any $K \Subset B_{1}^{ \pm} \cup B_{1}^{\prime}$, we have

$$
\begin{equation*}
\|U\|_{C^{1, \beta}(K)} \leq C(n, \alpha, M, K)\|U\|_{W^{1,2}\left(B_{1}\right)} . \tag{2.5.6}
\end{equation*}
$$

Proof. Let $K$ be a ball centered at 0 . Fix a small $r_{0}=r_{0}(n, \alpha, M, K)>0$ to be determined later. In particular, we will ask $r_{1}:=r_{0}^{\frac{2 n}{2 n+\alpha}} \Lambda^{1 / 2} \leq(1 / 2) \operatorname{dist}\left(K, \partial B_{1}\right)$, which implies that

$$
\widetilde{K}:=\left\{y \in B_{1}: \operatorname{dist}(y, K) \leq r_{1}\right\} \Subset B_{1} .
$$

Define

$$
\widehat{\nabla U}\left(y^{\prime}, y_{n}\right):= \begin{cases}\nabla U\left(y^{\prime}, y_{n}\right), & y_{n} \geq 0 \\ \nabla U\left(y^{\prime},-y_{n}\right), & y_{n}<0\end{cases}
$$

Our goal is to show that for $x_{0} \in K, 0<\rho<r<r_{0}$,

$$
\begin{align*}
\int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla U}-\langle\widehat{\nabla U}\rangle_{B_{\rho}\left(x_{0}\right)}\right|^{2} \leq C(n, \alpha, M)\left(\frac{\rho}{r}\right)^{n+\alpha} & \int_{B_{r}\left(x_{0}\right)}\left|\widehat{\nabla U}-\langle\widehat{\nabla U}\rangle_{B_{r}\left(x_{0}\right)}\right|^{2} \\
& +C(n, \alpha, M, K)\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+2 \beta} \tag{2.5.7}
\end{align*}
$$

Case 1. Suppose $x_{0} \in K \cap B_{1}^{\prime}$. For given $0<r<r_{0}$, we denote $\alpha^{\prime}:=1-\frac{\alpha}{8 n} \in(0,1)$, $R:=r^{\frac{2 n}{2 n+\alpha}}$. We then consider two cases:

$$
\sup _{\partial E_{R}\left(x_{0}\right)}|U| \leq C_{3}\left(\Lambda^{1 / 2} R\right)^{\alpha^{\prime}} \quad \text { and } \quad \sup _{\partial E_{R}\left(x_{0}\right)}|U|>C_{3}\left(\Lambda^{1 / 2} R\right)^{\alpha^{\prime}}
$$

where $C_{3}=2[U]_{0, \alpha^{\prime}, \widetilde{K}}=2 \sup _{\substack{y, z \in \widetilde{K} \\ y \neq z}} \frac{|U(y)-U(z)|}{|y-z|^{\alpha^{\prime}}}$.
Case 1.1. Assume that $\sup _{\partial E_{R}\left(x_{0}\right)}|U| \leq C_{3}\left(\Lambda^{1 / 2} R\right)^{\alpha^{\prime}}$. Then $u_{x_{0}}$ satisfies almost Signorini property at 0 in $B_{R}$ with

$$
\sup _{\partial B_{R}}\left|u_{x_{0}}\right| \leq C_{3}\left(\Lambda^{1 / 2} R\right)^{\alpha^{\prime}}
$$

Let $h$ be the Signorini replacement of $u_{x_{0}}$ in $B_{R}$. If we define

$$
\widehat{\nabla u_{x_{0}}}\left(y^{\prime}, y_{n}\right):= \begin{cases}\nabla u_{x_{0}}\left(y^{\prime}, y_{n}\right), & y_{n} \geq 0 \\ \nabla u_{x_{0}}\left(y^{\prime},-y_{n}\right), & y_{n}<0\end{cases}
$$

and

$$
\widehat{\nabla h}\left(y^{\prime}, y_{n}\right):= \begin{cases}\nabla h\left(y^{\prime}, y_{n}\right), & y_{n} \geq 0 \\ \nabla h\left(y^{\prime},-y_{n}\right), & y_{n}<0\end{cases}
$$

then we have

$$
\begin{align*}
& \int_{B_{\rho}}\left|\widehat{\nabla u_{x_{0}}}-\left\langle\widehat{\nabla u_{x_{0}}}\right\rangle_{B_{\rho}}\right|^{2} \leq 3 \int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{\rho}}\right|^{2}+6 \int_{B_{\rho}}\left|\widehat{\nabla u_{x_{0}}}-\widehat{\nabla h}\right|^{2},  \tag{2.5.8}\\
& \int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{r}}\right|^{2} \leq 3 \int_{B_{r}}\left|\widehat{\nabla u_{x_{0}}}-\left\langle\widehat{\nabla u_{x_{0}}}\right\rangle_{B_{r}}\right|^{2}+6 \int_{B_{r}}\left|\widehat{\nabla u_{x_{0}}}-\widehat{\nabla h}\right|^{2} . \tag{2.5.9}
\end{align*}
$$

Note that if $r_{0} \leq(3 / 4)^{\frac{2 n+\alpha}{\alpha}}$, then $r<(3 / 4) R$, thus by Proposition 2.5.1, the Signorini replacement $h$ satisfies, for $0<\rho<r$,

$$
\int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{\rho}}\right|^{2} \leq C(n, \alpha)\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{r}}\right|^{2}+C(n, \alpha) \frac{r^{n+1}}{R^{3}} \sup _{\partial B_{R}} h^{2} .
$$

Combining the above three inequalities, we obtain

$$
\begin{align*}
& \int_{B_{\rho}}\left|\widehat{\nabla u_{x_{0}}}-\left\langle\widehat{\nabla u_{x_{0}}}\right\rangle_{B_{\rho}}\right|^{2} \leq C(n, \alpha)\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{B_{r}}\left|\widehat{\nabla u_{x_{0}}}-\left\langle\widehat{\nabla u_{x_{0}}}\right\rangle_{B_{r}}\right|^{2} \\
&+C(n, \alpha) \frac{r^{n+1}}{R^{3}} \sup _{\partial B_{R}} h^{2}+C(n, \alpha) \int_{B_{r}}\left|\widehat{\nabla u_{x_{0}}}-\widehat{\nabla h}\right|^{2} . \tag{2.5.10}
\end{align*}
$$

Let us estimate the last term in the right-hand side of (2.5.10). Take $\delta=\delta(n, \alpha, M, K)>0$ such that $\delta<\operatorname{dist}\left(K, \partial B_{1}\right)$ and $\delta^{\alpha} \leq \varepsilon=\varepsilon\left(C_{1}, n, n+2 \alpha^{\prime}-2\right)$, where $C_{1}=C_{1}(n, M)$ is as in Proposition 2.4.2 and $\varepsilon$ is as in Lemma 1.2.2. If $r_{0} \leq\left(\Lambda^{-1 / 2} \delta\right)^{\frac{2 n+\alpha}{2 n}}$, then $\Lambda^{1 / 2} R<\delta$, thus, by following the proof of Theorem 2.4.1 up to (2.4.4), we have

$$
\int_{B_{\Lambda^{1 / 2} R}\left(x_{0}\right)}|\nabla U|^{2} \leq C(n, \alpha, M, K)\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2}\left(\Lambda^{1 / 2} R\right)^{n+2 \alpha^{\prime}-2}
$$

It follows that

$$
\int_{E_{R}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle \leq \Lambda \int_{B_{\Lambda^{1 / 2}}\left(x_{0}\right)}|\nabla U|^{2} \leq C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} R^{n+2 \alpha^{\prime}-2}
$$

Then by the change of variables (2.2.2), we have

$$
\begin{equation*}
\int_{B_{R}}\left|\nabla u_{x_{0}}\right|^{2} \leq C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} R^{n+2 \alpha^{\prime}-2} . \tag{2.5.11}
\end{equation*}
$$

Now we can estimate the third term in the right-hand side of (2.5.10):

$$
\begin{align*}
& \int_{B_{r}}\left|\widehat{\nabla u_{x_{0}}}-\widehat{\nabla h}\right|^{2}=2 \int_{B_{r}^{+}}\left|\nabla u_{x_{0}}-\nabla h\right|^{2} \\
& \leq 2 \int_{B_{R}}\left|\nabla u_{x_{0}}-\nabla h\right|^{2} \leq 2\left(\int_{B_{R}}\left|\nabla u_{x_{0}}\right|^{2}-\int_{B_{R}}|\nabla h|^{2}\right) \\
& \leq 2 M R^{\alpha} \int_{B_{R}}|\nabla h|^{2} \leq 2 M R^{\alpha} \int_{B_{R}}\left|\nabla u_{x_{0}}\right|^{2}  \tag{2.5.12}\\
& \leq C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} R^{n+\alpha+2 \alpha^{\prime}-2}=C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2 n+\alpha}\left(n-\frac{1}{2}\right)} .
\end{align*}
$$

To estimate the second term in the right-hand side of (2.5.10), we observe that

$$
\sup _{\partial B_{R}} h^{2}=\sup _{\partial B_{R}} u_{x_{0}}^{2}=\sup _{\partial E_{R}\left(x_{0}\right)} U^{2} \leq C_{3}^{2}\left(\Lambda^{1 / 2} R\right)^{2 \alpha^{\prime}} .
$$

Note that by (2.4.5), $C_{3} \leq C(n, \alpha, M, K)\|U\|_{W^{1,2}\left(B_{1}\right)}$. Thus,

$$
\frac{r^{n+1}}{R^{3}} \sup _{\partial B_{R}} h^{2} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} .
$$

Now (2.5.10) becomes

$$
\begin{align*}
\left.\int_{B_{\rho}}\left|\widehat{\nabla u_{x_{0}}}-\left\langle\widehat{\nabla u_{x_{0}}}\right\rangle_{B_{\rho}}\right|^{2} \leq C(n, \alpha)\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{B_{r}} \right\rvert\, \widehat{\nabla u_{x_{0}}} & -\left.\left\langle\widehat{\nabla u_{x_{0}}}\right\rangle_{B_{r}}\right|^{2} \\
& +C\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} . \tag{2.5.13}
\end{align*}
$$

We now want to deduce (2.5.7) from (2.5.13). The complication here is that the mapping $\bar{T}_{x_{0}}^{-1}$ does not preserve the even symmetry with respect to the thin plane, since the conormal direction $A\left(x_{0}\right) \mathrm{e}_{n}$ might be different from the normal direction $\mathrm{e}_{n}$ to $\Pi$ at $x_{0}$. To address this issue, by using the even symmetry of $\widehat{\nabla u_{x_{0}}}$, we rewrite (2.5.13) in terms of halfballs $B_{r}^{+}=B_{r} \cap \mathbb{R}_{+}^{n}$

$$
\begin{align*}
\left.\int_{B_{\rho}^{+}}\left|\nabla u_{x_{0}}-\left\langle\nabla u_{x_{0}}\right\rangle_{B_{\rho}^{+}}\right|^{2} \leq C(n, \alpha)\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{B_{r}^{+}} \right\rvert\, \nabla u_{x_{0}} & -\left.\left\langle\nabla u_{x_{0}}\right\rangle_{B_{r}^{+}}\right|^{2} \\
& +C\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} . \tag{2.5.14}
\end{align*}
$$

Similarly, if we denote $E_{r}^{+}\left(x_{0}\right)=E_{r}\left(x_{0}\right) \cap \mathbb{R}_{+}^{n}$, then using that $\bar{T}_{x_{0}}\left(E_{t}^{+}\left(x_{0}\right)\right)=B_{t}^{+}, t>0$, (2.5.14) becomes

$$
\begin{aligned}
\int_{E_{\rho}^{+}\left(x_{0}\right)}\left|\mathfrak{a}_{x_{0}} \nabla U-\left\langle\mathfrak{a}_{x_{0}} \nabla U\right\rangle_{E_{\rho}^{+}\left(x_{0}\right)}\right|^{2} \leq C(n, \alpha)\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{E_{r}^{+}\left(x_{0}\right)}\left|\mathfrak{a}_{x_{0}} \nabla U-\left\langle\mathfrak{a}_{x_{0}} \nabla U\right\rangle_{E_{r}^{+}\left(x_{0}\right)}\right|^{2} \\
+C \operatorname{det} \mathfrak{a}_{x_{0}}\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}}
\end{aligned}
$$

Repeating the argument that (2.3.4) implies (2.3.2) in the proof of Proposition 2.3.1, we have

$$
\begin{align*}
& \int_{B_{\rho}^{+}\left(x_{0}\right)}\left|\nabla U-\langle\nabla U\rangle_{B_{\rho}^{+}\left(x_{0}\right)}\right|^{2} \leq C\left(\frac{\rho}{r}\right)^{n+\alpha} \int_{B_{r}^{+}\left(x_{0}\right)}\left|\nabla U-\langle\nabla U\rangle_{B_{r}^{+}\left(x_{0}\right)}\right|^{2} \\
&+C\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} . \tag{2.5.15}
\end{align*}
$$

Then by the even symmetry of $\widehat{\nabla U}$, (2.5.15) implies (2.5.7).
Case 1.2. Now we assume that $\sup _{\partial E_{R}\left(x_{0}\right)}|U|>C_{3}\left(\Lambda^{1 / 2} R\right)^{\alpha^{\prime}}$. By the choice of $C_{3}=$ $2[U]_{0, \alpha^{\prime}, \widetilde{K}}$, we have either

$$
\begin{aligned}
& U \geq\left(C_{3} / 2\right)\left(\Lambda^{1 / 2} R\right)^{\alpha^{\prime}} \quad \text { in } E_{R}\left(x_{0}\right), \text { or } \\
& U \leq-\left(C_{3} / 2\right)\left(\Lambda^{1 / 2} R\right)^{\alpha^{\prime}} \quad \text { in } E_{R}\left(x_{0}\right)
\end{aligned}
$$

However, from $U \geq 0$ on $B_{1}^{\prime}$, the only possibility is

$$
U \geq\left(C_{3} / 2\right)\left(\Lambda^{1 / 2} R\right)^{\alpha^{\prime}} \quad \text { in } E_{R}\left(x_{0}\right)
$$

Consequently,

$$
u_{x_{0}} \geq\left(C_{3} / 2\right)\left(\Lambda^{1 / 2} R\right)^{\alpha^{\prime}} \quad \text { in } B_{R} .
$$

If we let $h$ again be the Signorini replacement of $u_{x_{0}}$ in $B_{R}$, then the positivity of $h=u_{x_{0}}>0$ on $\partial B_{R}$ and superharmonicity of $h$ in $B_{R}$ give that $h>0$ in $B_{R}$, and hence $h$ is harmonic in $B_{R}$. Thus,

$$
\int_{B_{\rho}}\left|\nabla h-\langle\nabla h\rangle_{B_{\rho}}\right|^{2} \leq\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\nabla h-\langle\nabla h\rangle_{B_{r}}\right|^{2}, \quad 0<\rho<r .
$$

We next decompose $h=h^{*}+h^{\sharp}$ in $B_{R}$ as in (2.4.2). Note that since both $h$ and $h^{\sharp}$ are harmonic, $h^{*}$ must be harmonic as well. Then we have

$$
\begin{aligned}
\int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{\rho}}\right|^{2} & \leq 3 \int_{B_{\rho}}\left|\nabla h-\langle\nabla h\rangle_{B_{\rho}}\right|^{2}+6 \int_{B_{\rho}}|\widehat{\nabla h}-\nabla h|^{2} \\
& =3 \int_{B_{\rho}}\left|\nabla h-\langle\nabla h\rangle_{B_{\rho}}\right|^{2}+6 \int_{B_{\rho}^{-}}\left(\left|2 \nabla_{y^{\prime}} h^{\sharp}\right|^{2}+\left|2 \partial_{y_{n}} h^{*}\right|^{2}\right) \\
& =3 \int_{B_{\rho}}\left|\nabla h-\langle\nabla h\rangle_{B_{\rho}}\right|^{2}+12 \int_{B_{\rho}}\left(\left|\nabla_{y^{\prime}} h^{\sharp}\right|^{2}+\left|\partial_{y_{n}} h^{*}\right|^{2}\right),
\end{aligned}
$$

and similarly,

$$
\int_{B_{r}}\left|\nabla h-\langle\nabla h\rangle_{B_{r}}\right|^{2} \leq 3 \int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{r}}\right|^{2}+12 \int_{B_{r}}\left(\left|\nabla_{y^{\prime}} h^{\sharp}\right|^{2}+\left|\partial_{y_{n}} h^{*}\right|^{2}\right) .
$$

Combining the above three inequalities, we have that for all $0<\rho<r$

$$
\begin{equation*}
\int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{\rho}}\right|^{2} \leq 3\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{r}}\right|^{2}+48 \int_{B_{r}}\left(\left|\nabla_{y^{\prime}} h^{\sharp}\right|^{2}+\left|\partial_{y_{n}} h^{*}\right|^{2}\right) . \tag{2.5.16}
\end{equation*}
$$

Now, note that if $r_{0} \leq(1 / 2)^{\frac{2 n+\alpha}{\alpha}}$, then $r \leq R / 2$. By the harmonicity of both $h^{*}$ and $h^{\sharp}$ in $B_{R}$, we have

$$
\begin{aligned}
\sup _{B_{R / 2}}\left|D^{2} h^{*}\right|+\sup _{B_{R / 2}}\left|D^{2} h^{\sharp}\right| & \leq \frac{C(n)}{R}\left(\sup _{B_{(3 / 4) R}}\left|\nabla h^{*}\right|+\sup _{B_{(3 / 4) R}}\left|\nabla h^{\sharp}\right|\right) \\
& \leq \frac{C(n)}{R^{1+\frac{n}{2}}}\left(\int_{B_{R}}\left|\nabla h^{*}\right|^{2}+\int_{B_{R}}\left|\nabla h^{\sharp}\right|^{2}\right)^{1 / 2} \\
& =\frac{C(n)}{R^{1+\frac{n}{2}}}\left(\int_{B_{R}}|\nabla h|^{2}\right)^{1 / 2} \leq \frac{C(n)}{R^{1+\frac{n}{2}}}\left(\int_{B_{R}}\left|\nabla u_{x_{0}}\right|^{2}\right)^{1 / 2} \\
& \leq C(n, \alpha, M, K)\|\nabla U\|_{L^{2}\left(B_{1}\right)} R^{\alpha^{\prime}-2}
\end{aligned}
$$

where the last inequality follows from (2.5.11). Also, note that $\nabla_{y^{\prime}} h^{\sharp}=\partial_{y_{n}} h^{*}=0$ on $B_{R / 2}^{\prime}$. Thus, for $y=\left(y^{\prime}, y_{n}\right) \in B_{r}$, we have

$$
\begin{aligned}
\left|\nabla_{y^{\prime}} h^{\sharp}\right|+\left|\partial_{y_{n}} h^{*}\right| & \leq\left|y_{n}\right|\left(\sup _{B_{R / 2}}\left|D^{2} h^{*}\right|+\sup _{B_{R / 2}}\left|D^{2} h^{\sharp}\right|\right) \\
& \leq C\|\nabla U\|_{L^{2}\left(B_{1}\right)} r R^{\alpha^{\prime}-2} \\
& =C\|\nabla U\|_{L^{2}\left(B_{1}\right)} r^{1+\frac{2 n}{2 n+\alpha}\left(\alpha^{\prime}-2\right)},
\end{aligned}
$$

with $C=(n, \alpha, M, K)$. Hence, it follows that

$$
\begin{equation*}
\int_{B_{r}}\left|\nabla_{y^{\prime}} h^{\sharp}\right|^{2}+\left|\partial_{y_{n}} h^{*}\right|^{2} \leq C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+2+\frac{4 n}{2 n+\alpha}\left(\alpha^{\prime}-2\right)} \leq C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} . \tag{2.5.17}
\end{equation*}
$$

Combining (2.5.16) and (2.5.17), we obtain

$$
\begin{equation*}
\int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{\rho}}\right|^{2} \leq 3\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{r}}\right|^{2}+C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} . \tag{2.5.18}
\end{equation*}
$$

Note that (2.5.12) was induced in Case 1.1 without the use of the assumption $\sup _{\partial E_{r}\left(x_{0}\right)}|U| \leq$ $C_{3}\left(\Lambda^{1 / 2} R\right)^{\alpha^{\prime}}$, so it is also valid in this case. Finally, (2.5.8), (2.5.9), (2.5.12) and (2.5.18) give

$$
\begin{aligned}
\int_{B_{\rho}}\left|\widehat{\nabla u_{x_{0}}}-\left\langle\widehat{\nabla u_{x_{0}}}\right\rangle_{B_{\rho}}\right|^{2} \leq & 3 \int_{B_{\rho}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{\rho}}\right|^{2}+6 \int_{B_{\rho}}\left|\widehat{\nabla u_{x_{0}}}-\widehat{\nabla h}\right|^{2} \\
\leq & 9\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla h}-\langle\widehat{\nabla h}\rangle_{B_{r}}\right|^{2}+C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} \\
& +6 \int_{B_{\rho}}\left|\widehat{\nabla u_{x_{0}}}-\widehat{\nabla h}\right|^{2} \\
\leq & 27\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla u_{x_{0}}}-\left\langle\widehat{\nabla u_{x_{0}}}\right\rangle_{B_{r}}\right|^{2}+C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} \\
& +60 \int_{B_{r}}\left|\widehat{\nabla u_{x_{0}}}-\widehat{\nabla h}\right|^{2} \\
\leq & 27\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla u_{x_{0}}}-\left\langle\widehat{\nabla u_{x_{0}}}\right\rangle_{B_{r}}\right|^{2}+C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)}} \\
& +C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2 n+\alpha}(n-1 / 2)} \\
\leq & 27\left(\frac{\rho}{r}\right)^{n+2} \int_{B_{r}}\left|\widehat{\nabla u_{x_{0}}}-\left\langle\widehat{\nabla u_{x_{0}}}\right\rangle_{B_{r}}\right|^{2}+C\|\nabla U\|_{L^{2}\left(B_{1}\right)}^{2} r^{n+\frac{\alpha}{2(2 n+\alpha)} .}
\end{aligned}
$$

As we have seen in Case 1.1, this implies (2.5.7). This completes the proof of (2.5.7) when $x_{0} \in K \cap B_{1}^{\prime}$.

Case 2. The extension of (2.5.7) to general $x_{0} \in K$ follows from the combination of Case 1 and (2.3.5). The argument is the same as Case 2 in the proof of Theorem 1.4.1.

Thus, the estimate (2.5.7) holds in all possible cases.
To complete the proof of the theorem, we now apply Lemma 1.2.2 to the estimate (2.5.7) to obtain for $0<\rho<r<r_{0}$

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla U}-\langle\widehat{\nabla U}\rangle_{B_{\rho}\left(x_{0}\right)}\right|^{2} \leq C\left[\left(\frac{\rho}{r}\right)^{n+2 \beta} \int_{B_{r}\left(x_{0}\right)}\left|\widehat{\nabla U}-\langle\widehat{\nabla U}\rangle_{B_{r}\left(x_{0}\right)}\right|^{2}+\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} \rho^{n+2 \beta}\right] .
$$

Taking $r \nearrow r_{0}=r_{0}(n, \alpha, M, K)$, we have

$$
\int_{B_{\rho}\left(x_{0}\right)}\left|\widehat{\nabla U}-\langle\widehat{\nabla U}\rangle_{B_{\rho}\left(x_{0}\right)}\right|^{2} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} \rho^{n+2 \beta}
$$

with $C=C(n, \alpha, M, K)$. Then by the Campanato space embedding this readily implies that $\widehat{\nabla U} \in C^{0, \beta}(K)$ with

$$
\|\widehat{\nabla U}\|_{C^{0, \beta}(K)} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}
$$

Since $\widehat{\nabla U}=\nabla U$ in $B_{1}^{+} \cup B_{1}^{\prime}$, we therefore conclude that

$$
U \in C^{1, \beta}\left(K \cap\left(B_{1}^{+} \cup B_{1}^{\prime}\right)\right),
$$

and combining with the bound in Theorem 2.4.1, we also deduce that

$$
\|U\|_{C^{1, \beta}\left(K \cap\left(B_{1}^{+} \cup B_{1}^{\prime}\right)\right)} \leq C(n, \alpha, M, K)\|U\|_{W^{1,2}\left(B_{1}\right)}
$$

To see the $C^{1, \beta}$ regularity of $U$ in $B_{1}^{-} \cup B_{1}^{\prime}$, we simply observe that the function $U\left(y^{\prime},-y_{n}\right)$ is also an almost minimizer of the Signorini problem with the appropriately modified coefficient matrix $A$.

### 2.6 Quasisymmetric almost minimizers

In the study of the free boundary in the Signorini problem, the even symmetry of the minimizer with respect to the thin space plays a crucial role. The even symmetry guarantees that the growth rate of the minimizer $u$ over "thick" balls $B_{r}\left(x_{0}\right) \subset \mathbb{R}^{n}$ matches the growth rate over thin balls $B_{r}^{\prime}\left(x_{0}\right) \subset \Pi$. This allows to use tools such as Almgren's monotonicity formula (see the next section) to classify the free boundary points. Without even symmetry, minimizers may have an odd component, vanishing on the thin space $\Pi$ that may create a mismatch of growth rates on the thick and thin spaces.

In the case of minimizers of the Signorini problem (with $A=I$ ) or harmonic functions, it is easy to see that the even symmetrization

$$
u^{*}(x)=\frac{u\left(x^{\prime}, x_{n}\right)+u\left(x^{\prime},-x_{n}\right)}{2}
$$

is still a minimizer. Unfortunately, the even symmetrization may destroy the almost minimizing property, as well as the minimizing property with variable coefficients, as can be seen from the following simple example.

Example 2.6.1. Let $u:(-1,1) \rightarrow \mathbb{R}$ be defined by $u(x)=x+x^{2} / 4$. Then $u$ is an almost harmonic function in $(-1,1)$ with a gauge function $\omega(r)=C(\alpha) r^{\alpha}$ for $0<\alpha<1$. In fact, $u$ is a minimizer of the energy functional

$$
\int(1+x / 2)^{-1}\left(v^{\prime}\right)^{2}
$$

with a Lipschitz function $A(x)=(1+x / 2)^{-1}$ in $(-1,1)$. On the other hand, the even symmetrization

$$
u^{*}(x)=\frac{u(x)+u(-x)}{2}=\frac{x^{2}}{4}
$$

is not almost harmonic for any gauge function $\omega(r)$. Indeed, for any small $\delta>0$, if we take a competitor $v=\delta^{2} / 4$ in $(-\delta, \delta)$, then it satisfies $\int_{-\delta}^{\delta}\left|v^{\prime}\right|^{2}=0$ and if $u^{*}$ were almost harmonic, we would have that $\int_{-\delta}^{\delta}\left|\left(u^{*}\right)^{\prime}\right|^{2}=0$ as well, implying that $u^{*}$ is constant in $(-\delta, \delta)$, a contradiction.

To overcome this difficulty, we need to impose the $A$-quasisymmetry condition on almost minimizers $U$, that we have already stated in Definition 2.1.2. In this section, we give more details on quasisymmetric almost minimizers.

Recall that for each $x_{0} \in B_{1}^{\prime}$, we defined a reflection matrix $P_{x_{0}}$ by

$$
P_{x_{0}}=I-2 \frac{A\left(x_{0}\right) \mathrm{e}_{n} \otimes \mathrm{e}_{n}}{a_{n n}\left(x_{0}\right)} .
$$

From the ellipticity of $A$, we have $a_{n n}\left(x_{0}\right) \geq \lambda$, thus $P_{x_{0}}$ is well-defined. Note that $P_{x_{0}}^{2}=I$. Besides, $\left.P_{x_{0}}\right|_{\Pi}=\left.I\right|_{\Pi}$ and $P_{x_{0}} E_{r}\left(x_{0}\right)=E_{r}\left(x_{0}\right)$. We then define the "skewed" even/odd symmetrizations of the almost minimizer $U$ in $B_{1}$ by

$$
\begin{aligned}
U_{x_{0}}^{*}(x) & :=\frac{U(x)+U\left(P_{x_{0}} x\right)}{2}, \\
U_{x_{0}}^{\sharp}(x) & :=\frac{U(x)-U\left(P_{x_{0}} x\right)}{2} .
\end{aligned}
$$



Figure 2.2. Reflection $P_{x_{0}}$ : here $\bar{x}=P_{x_{0}} x, y=\bar{T}_{x_{0}}(x)$, and $\bar{y}=\left(y^{\prime},-y_{n}\right)=\bar{T}_{x_{0}}(\bar{x})$

Note that $U_{x_{0}}^{*}$ and $U_{x_{0}}^{\sharp}$ may not be defined in all of $B_{1}$, but are defined in any ellipsoid $E_{r}\left(x_{0}\right)$ as long as it is contained in $B_{1}$. Note also that $U=U_{x_{0}}^{*}$ and $U_{x_{0}}^{\sharp}=0$ on $\Pi$. Further, we note that transformed with $\bar{T}_{x_{0}}, P_{x_{0}}$ becomes an even reflection with respect to $\Pi$, i.e.,

$$
\bar{T}_{x_{0}} \circ P_{x_{0}} \circ \bar{T}_{x_{0}}^{-1}(y)=\left(y^{\prime},-y_{n}\right),
$$

see Fig 2.2. Therefore, denoting

$$
\begin{aligned}
u_{x_{0}}^{*}(y) & :=\frac{u_{x_{0}}\left(y^{\prime}, y_{n}\right)+u_{x_{0}}\left(y^{\prime},-y_{n}\right)}{2}, \\
u_{x_{0}}^{\sharp}(y) & :=\frac{u_{x_{0}}\left(y^{\prime}, y_{n}\right)-u_{x_{0}}\left(y^{\prime},-y_{n}\right)}{2},
\end{aligned}
$$

the even/odd symmetrizations of $u_{x_{0}}$ about $\Pi$, we will have

$$
U_{x_{0}}^{*} \circ \bar{T}_{x_{0}}^{-1}=u_{x_{0}}^{*}, \quad U_{x_{0}}^{\sharp} \circ \bar{T}_{x_{0}}^{-1}=u_{x_{0}}^{\sharp} .
$$

We also observe that the symmetries of $u_{x_{0}}^{*}$ and $u_{x_{0}}^{\sharp}$ imply the following decompositions

$$
\begin{align*}
\int_{B_{r}} u_{x_{0}}^{2} & =\int_{B_{r}}\left(u_{x_{0}}^{*}\right)^{2}+\int_{B_{r}}\left(u_{x_{0}}^{\sharp}\right)^{2},  \tag{2.6.1}\\
\int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2} & =\int_{B_{r}}\left|\nabla u_{x_{0}}^{*}\right|^{2}+\int_{B_{r}}\left|\nabla u_{x_{0}}^{\sharp}\right|^{2}, \tag{2.6.2}
\end{align*}
$$

which after a change of variables, can also be written as

$$
\begin{align*}
\int_{E_{r}\left(x_{0}\right)} U^{2} & =\int_{E_{r}\left(x_{0}\right)}\left(U_{x_{0}}^{*}\right)^{2}+\int_{E_{r}\left(x_{0}\right)}\left(U_{x_{0}}^{\sharp}\right)^{2},  \tag{2.6.3}\\
\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle & =\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U_{x_{0}}^{*}, \nabla U_{x_{0}}^{*}\right\rangle+\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U_{x_{0}}^{\sharp}, \nabla U_{x_{0}}^{\sharp}\right\rangle . \tag{2.6.4}
\end{align*}
$$

We now recall that by Definition 2.1.2, $U \in W^{1,2}\left(B_{1}\right)$ is called $A$-quasisymmetric if there is a constant $Q>0$ such that

$$
\begin{equation*}
\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U, \nabla U\right\rangle \leq Q \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U_{x_{0}}^{*}, \nabla U_{x_{0}}^{*}\right\rangle, \tag{2.6.5}
\end{equation*}
$$

whenever $E_{r}\left(x_{0}\right) \Subset B_{1}$ and $x_{0} \in B_{1}$. By the uniform ellipticity of $A$, (2.6.5) is equivalent to

$$
\int_{E_{r}\left(x_{0}\right)}|\nabla U|^{2} \leq Q \int_{E_{r}\left(x_{0}\right)}\left|\nabla U_{x_{0}}^{*}\right|^{2}
$$

by changing $Q$ to $Q(\Lambda / \lambda)$, if necessary. Besides, using (2.6.4), (2.6.5) is also equivalent to

$$
\begin{equation*}
\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U_{x_{0}}^{\sharp}, \nabla U_{x_{0}}^{\sharp}\right\rangle \leq C \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U_{x_{0}}^{*}, \nabla U_{x_{0}}^{*}\right\rangle, \tag{2.6.6}
\end{equation*}
$$

with some $C=C(Q)$.
Lemma 2.6.2. Let $U$ be an A-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$, with constant $Q>0$. Then there are $r_{1}=r_{1}(n, \alpha, M, Q)>0$ and $M_{1}=M_{1}(n, M, Q)>$ 0 such that

$$
\begin{equation*}
\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla U_{x_{0}}^{*}, \nabla U_{x_{0}}^{*}\right\rangle \leq\left(1+M_{1} r^{\alpha}\right) \int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla W, \nabla W\right\rangle, \tag{2.6.7}
\end{equation*}
$$

whenever $E_{r}\left(x_{0}\right) \Subset B_{1}, x_{0} \in B_{1}, 0<r<r_{1}$, and $W \in \mathfrak{K}_{0, U_{x_{0}}^{*}}\left(E_{r}\left(x_{0}\right), \Pi\right)$.
Remark 2.6.3. Since we are interested in local results, in what follows, we will assume without loss of generality that $r_{1}=1$ and $M_{1}=M$.

Proof. Let $V$ be the energy minimizer of

$$
\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla V, \nabla V\right\rangle \quad \text { on } \mathfrak{K}_{0, U}\left(E_{r}\left(x_{0}\right), \Pi\right) .
$$

Then $v_{x_{0}}=V \circ \bar{T}_{x_{0}}^{-1}$ is the energy minimizer of

$$
\int_{B_{r}}\left|\nabla v_{x_{0}}\right|^{2} \quad \text { on } \mathfrak{K}_{0, u_{x_{0}}}\left(B_{r}, \Pi\right) .
$$

Note that $v_{x_{0}}^{*}$ is a solution of the Signorini problem, even in $y_{n}$, with $v_{x_{0}}^{*}=u_{x_{0}}^{*}$ on $\partial B_{r}$. Similarly, $v_{x_{0}}^{\sharp}$ is a harmonic function, odd in $y_{n}$, with $v_{x_{0}}^{\sharp}=u_{x_{0}}^{\sharp}$ on $\partial B_{r}$. Thus, $v_{x_{0}}^{*}$ is the energy minimizer of

$$
\int_{B_{r}}\left|\nabla v_{x_{0}}^{*}\right|^{2} \quad \text { on } \mathfrak{K}_{0, u_{x_{0}}^{*}}\left(B_{r}, \Pi\right)
$$

and so $V_{x_{0}}^{*}$ is the energy minimizer of

$$
\int_{E_{r}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla V_{x_{0}}^{*}, \nabla V_{x_{0}}^{*}\right\rangle \quad \text { on } \mathfrak{K}_{0, U_{x_{0}}^{*}}\left(E_{r}\left(x_{0}\right), \Pi\right) .
$$

Thus, to show (2.6.7), it is enough to show

$$
\int_{B_{r}}\left|\nabla u_{x_{0}}^{*}\right|^{2} \leq\left(1+M_{1} r^{\alpha}\right) \int_{B_{r}}\left|\nabla v_{x_{0}}^{*}\right|^{2}
$$

To this end, we first observe that the quasisymmetry of $U$ implies the quasisymmetry of $u_{x_{0}}$ :

$$
\int_{B_{r}}\left|\nabla u_{x_{0}}^{\sharp}\right|^{2} \leq C \int_{B_{r}}\left|\nabla u_{x_{0}}^{*}\right|^{2} .
$$

Using this, together with the symmetry of $u_{x_{0}}^{*}, u_{x_{0}}^{\sharp}, v_{x_{0}}^{*}$ and $v_{x_{0}}^{\sharp}$, we have

$$
\begin{aligned}
& \int_{B_{r}}\left|\nabla u_{x_{0}}^{*}\right|^{2}=\int_{B_{r}}\left|\nabla u_{x_{0}}\right|^{2}-\int_{B_{r}}\left|\nabla u_{x_{0}}^{\sharp}\right|^{2} \\
& \leq\left(1+M r^{\alpha}\right) \int_{B_{r}}\left|\nabla v_{x_{0}}\right|^{2}-\int_{B_{r}}\left|\nabla u_{x_{0}}^{\sharp}\right|^{2} \\
&=\left(1+M r^{\alpha}\right) \int_{B_{r}}\left|\nabla v_{x_{0}}^{*}\right|^{2}+\left(1+M r^{\alpha}\right) \int_{B_{r}}\left|\nabla v_{x_{0}}^{\sharp}\right|^{2}-\int_{B_{r}}\left|\nabla u_{x_{0}}^{\sharp}\right|^{2} \\
& \leq\left(1+M r^{\alpha}\right) \int_{B_{r}}\left|\nabla v_{x_{0}}^{*}\right|^{2}+M r^{\alpha} \int_{B_{r}}\left|\nabla u_{x_{0}}^{\sharp}\right|^{2}
\end{aligned}
$$

$$
\leq\left(1+M r^{\alpha}\right) \int_{B_{r}}\left|\nabla v_{x_{0}}^{*}\right|^{2}+C M r^{\alpha} \int_{B_{r}}\left|\nabla u_{x_{0}}^{*}\right|^{2} .
$$

Therefore,

$$
\int_{B_{r}}\left|\nabla u_{x_{0}}^{*}\right|^{2} \leq \frac{1+M r^{\alpha}}{1-C M r^{\alpha}} \int_{B_{r}}\left|\nabla v_{x_{0}}^{*}\right|^{2} \leq\left(1+M_{1} r^{\alpha}\right) \int_{B_{r}}\left|\nabla v_{x_{0}}^{*}\right|^{2},
$$

for $0<r<r_{1}=(2 C M)^{-1 / \alpha}$, as desired.

Remark 2.6.4. If $U$ satisfies the following weak quasisymmetry with order $-\gamma$ :

$$
\int_{E_{r}\left(x_{0}\right)}|\nabla U|^{2} \leq Q r^{-\gamma} \int_{E_{r}\left(x_{0}\right)}\left|\nabla U_{x_{0}}^{*}\right|^{2},
$$

whenever $E_{r}\left(x_{0}\right) \Subset B_{1}, x_{0} \in B_{1}^{\prime}$ for some $0<\gamma<\alpha$, then it is easy to see from the proof of Lemma 2.6.2 that $U_{x_{0}}^{*}$ satisfies (2.6.7), but with $\alpha-\gamma>0$ instead of $\alpha$.

Theorem 2.6.5. Let $U$ be an $A$-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$. Then for $x_{0} \in B_{1 / 2}^{\prime}$ and $0<r \leq(1 / 2) \Lambda^{-1 / 2}$, we have $U_{x_{0}}^{*} \in C^{1, \beta}\left(E_{r}^{ \pm}\left(x_{0}\right) \cup E_{r}^{\prime}\left(x_{0}\right)\right)$ with $\beta=\frac{\alpha}{4(2 n+\alpha)}$. Moreover,

$$
\left\|U_{x_{0}}^{*}\right\|_{C^{1, \beta}(K)} \leq C(n, \alpha, M, K, r)\left\|U_{x_{0}}^{*}\right\|_{W^{1,2}\left(E_{r}\left(x_{0}\right)\right)}
$$

for any $K \Subset E_{r}^{ \pm}\left(x_{0}\right) \cup E_{r}^{\prime}\left(x_{0}\right)$. Similarly, $u_{x_{0}}^{*} \in C^{1, \beta}\left(B_{r}^{ \pm} \cup B_{r}^{\prime}\right)$ with

$$
\left\|u_{x_{0}}^{*}\right\|_{C^{1, \beta}(K)} \leq C(n, \alpha, M, K, r)\left\|u_{x_{0}}^{*}\right\|_{W^{1,2}\left(B_{r}\right)},
$$

for any $K \Subset B_{r}^{ \pm} \cup B_{r}^{\prime}$.
Proof. From Theorem 2.5.1, we have $U \in C^{1, \beta}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)$, which immediately gives $U_{x_{0}}^{*} \in$ $C^{1, \beta}\left(E_{r}^{ \pm}\left(x_{0}\right) \cup E_{r}^{\prime}\left(x_{0}\right)\right)$, by using the inclusion $E_{r}\left(x_{0}\right) \subset B_{\Lambda^{1 / 2} r}\left(x_{0}\right) \subset B_{1}$. Thus, for

$$
\widehat{\nabla U_{x_{0}}^{*}}\left(x^{\prime}, x_{n}\right):= \begin{cases}\nabla U_{x_{0}}^{*}\left(x^{\prime}, x_{n}\right), & x_{n} \geq 0 \\ \nabla U_{x_{0}}^{*}\left(x^{\prime},-x_{n}\right), & x_{n}<0\end{cases}
$$

we have $\widehat{\nabla U_{x_{0}}^{*}} \in C^{0, \beta}\left(E_{r}\left(x_{0}\right)\right)$ with

$$
\left\|\widehat{\nabla U_{x_{0}}^{*}}\right\|_{C^{0, \beta}(K)} \leq C(n, \alpha, M, K, r)\|U\|_{W^{1,2}\left(E_{r}\left(x_{0}\right)\right)}
$$

for any $K \Subset E_{r}\left(x_{0}\right)$. Hence, it is enough to show that

$$
\|U\|_{W^{1,2}\left(E_{r}\left(x_{0}\right)\right)} \leq C\left\|U_{x_{0}}^{*}\right\|_{W^{1,2}\left(E_{r}\left(x_{0}\right)\right)}
$$

Now, note that by (2.6.3)-(2.6.4), we readily have

$$
\|U\|_{W^{1,2}\left(E_{r}\left(x_{0}\right)\right)} \leq C\left(\left\|U_{x_{0}}^{*}\right\|_{W^{1,2}\left(E_{r}\left(x_{0}\right)\right)}+\left\|U_{x_{0}}^{\sharp}\right\|_{W^{1,2}\left(E_{r}\left(x_{0}\right)\right)}\right),
$$

and thus, it will suffice to show that

$$
\left\|U_{x_{0}}^{\sharp}\right\|_{W^{1,2}\left(E_{r}\left(x_{0}\right)\right)} \leq C\left\|U_{x_{0}}^{*}\right\|_{W^{1,2}\left(E_{r}\left(x_{0}\right)\right)} .
$$

By the symmetry again,

$$
\left\langle U_{x_{0}}^{\sharp}\right\rangle_{E_{r}\left(x_{0}\right)}=\left\langle u_{x_{0}}^{\sharp}\right\rangle_{B_{r}}=0,
$$

thus by Poincare's inequality,

$$
\begin{equation*}
\left\|U_{x_{0}}^{\sharp}\right\|_{L^{2}\left(E_{r}\left(x_{0}\right)\right)} \leq C(n, M) r\left\|\nabla U_{x_{0}}^{\sharp}\right\|_{L^{2}\left(E_{r}\left(x_{0}\right)\right)} . \tag{2.6.8}
\end{equation*}
$$

Finally, by the quasisymmetry of $U$, we have

$$
\left\|\nabla U_{x_{0}}^{\sharp}\right\|_{L^{2}\left(E_{r}\left(x_{0}\right)\right)} \leq C\left\|\nabla U_{x_{0}}^{*}\right\|_{L^{2}\left(E_{r}\left(x_{0}\right)\right)},
$$

see (2.6.6). This completes the proof of the theorem for $U_{x_{0}}^{*}$.
Applying now the affine transformation $\bar{T}_{x_{0}}$, we obtain the part of the theorem for $u_{x_{0}}^{*}$.
We complete this section with a version of Signorini's complementarity condition that will play an important role in the analysis of the free boundary.

Lemma 2.6.6 (Complementarity condition). Let $U$ be an $A$-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$, and $x_{0} \in B_{1 / 2}^{\prime}$. Then $u_{x_{0}}^{*}$ satisfies the following complementarity condition

$$
u_{x_{0}}^{*}\left(\partial_{y_{n}}^{+} u_{x_{0}}^{*}\right)=0 \quad \text { on } B_{R_{0}}^{\prime}, \quad R_{0}=(1 / 2) \Lambda^{-1 / 2}
$$

where $\partial_{y_{n}}^{+} u_{x_{0}}^{*}$ on $B_{R_{0}}^{\prime}$ is computed as the limit from inside $B_{R_{0}}^{+}$. Moreover, if $x_{0} \in \Gamma(U)$, then

$$
u_{x_{0}}^{*}(0)=0 \quad \text { and } \quad\left|\widehat{\nabla u_{x_{0}}^{*}}(0)\right|=0 .
$$

Proof. Let $y_{0} \in B_{R_{0}}^{\prime}$ be such that $u_{x_{0}}^{*}\left(y_{0}\right)>0$. Then we need to show that $\partial_{y_{n}}^{+} u_{x_{0}}^{*}\left(y_{0}\right)=0$. Since $u_{x_{0}}=u_{x_{0}}^{*}$ on $\Pi$, we have $u_{x_{0}}\left(y_{0}\right)>0$ and by continuity $u_{x_{0}}>0$ in a small ball $B_{\delta}\left(y_{0}\right)$. Then $U>0$ in $\Omega=\bar{T}_{x_{0}}^{-1}\left(B_{\delta}\left(y_{0}\right)\right)$. We claim now that $U$ is almost $A$-harmonic in $\Omega$. Indeed, if $E_{r}(y) \Subset \Omega$ (not necessarily with $y \in B_{1}^{\prime}$ ) and $V$ is $A(y)$-harmonic replacement of $U$ on $E_{r}(y)$ (i.e. $\operatorname{div}(A(y) \nabla V)=0$ in $E_{r}(y)$ with $V=U$ on $\left.\partial E_{r}(y)\right)$, then since $V=U>0$ on $\partial E_{r}(y)$, by the minimum principle $V>0$ on $\overline{E_{r}(y)}$. This means that $V \in \mathfrak{K}_{0, U}\left(E_{r}(y), \Pi\right)$ and therefore we must have

$$
\int_{E_{r}(y)}\langle A(y) \nabla U, \nabla U\rangle \leq(1+\omega(r)) \int_{E_{r}(y)}\langle A(y) \nabla V, \nabla V\rangle,
$$

which also implies that $U$ is an almost $A$-harmonic function in $\Omega$. Hence, $U \in C^{1, \alpha / 2}(\Omega)$ by Theorem 2.3.2, implying also that $u_{x_{0}} \in C^{1, \alpha / 2}\left(B_{\delta}\left(y_{0}\right)\right)$. Consequently, also $u_{x_{0}}^{*} \in$ $C^{1, \alpha / 2}\left(B_{\delta}\left(y_{0}\right)\right)$ and by even symmetry in the $y_{n}$-variable, we therefore conclude that $\partial_{y_{n}}^{+} u_{x_{0}}^{*}\left(y_{0}\right)=$ 0 .

The second part of the lemma now follows by the $C^{1, \beta}$ regularity and the complementarity condition.

### 2.7 Weiss- and Almgren-type monotonicity formulas

In this section we introduce two technical tools: Weiss- and Almgren-type monotonicity formulas, that will play a fundamental role in the analysis of the free boundary. In fact, the
proofs of these formulas follow immediately from the case $A \equiv I$, following the deskewing procedure.

To proceed, we fix a constant $\kappa_{0}>0$. We can take it as large as we want, however, some constants in what follows, will depend on $\kappa_{0}$. Then for $0<\kappa<\kappa_{0}$, we consider the Weiss-type energy functional introduced in Chapter 1:

$$
W_{\kappa}\left(t, v, x_{0}\right):=\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-2}}\left[\int_{B_{t}\left(x_{0}\right)}|\nabla v|^{2}-\kappa \frac{1-b t^{\alpha}}{t} \int_{\partial B_{t}\left(x_{0}\right)} v^{2}\right],
$$

with

$$
a=a_{\kappa}=\frac{M(n+2 \kappa-2)}{\alpha}, \quad b=\frac{M\left(n+2 \kappa_{0}\right)}{\alpha} .
$$

(The formula in Chapter 1 corresponds to the case $M=1$.) Based on that, we define an appropriate version of Weiss's functional for our problem. For a function $V$ in $E_{r}\left(x_{0}\right)$, let

$$
\begin{equation*}
W_{\kappa}^{A}\left(t, V, x_{0}\right):=\frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-2}}\left[\int_{E_{t}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla V, \nabla V\right\rangle-\kappa \frac{1-b t^{\alpha}}{t} \int_{\partial E_{t}\left(x_{0}\right)} V^{2} \mu_{x_{0}}\left(x-x_{0}\right)\right], \tag{2.7.1}
\end{equation*}
$$

for $0<t<r$, with $a, b$ same as above, where the weight $\mu_{x_{0}}$ is as in (2.2.4). Note that by the change of variables formulas (2.2.1)-(2.2.3), we have

$$
\begin{equation*}
W_{\kappa}^{A}\left(t, V, x_{0}\right):=\operatorname{det} \mathfrak{a}_{x_{0}} W_{\kappa}\left(t, v_{x_{0}}, 0\right), \quad v_{x_{0}}=V \circ \bar{T}_{x_{0}}^{-1} \tag{2.7.2}
\end{equation*}
$$

Let now $U$ be an $A$-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$ and $x_{0} \in B_{1 / 2}$. By Lemma 2.6.2, $U_{x_{0}}^{*}$ satisfies the almost $A$-Signorini property at $x_{0}$ in $E_{(1 / 2) \Lambda^{-1 / 2}}\left(x_{0}\right)$. Thus $u_{x_{0}}^{*}$ also satisfies the almost Signorini property at 0 in $B_{(1 / 2) \Lambda^{-1 / 2}}$. By using this observation, we then have the following Weiss-type monotonicity formulas for $U_{x_{0}}^{*}$ and $u_{x_{0}}^{*}$.

Theorem 2.7.1 (Weiss-type monotonicity formula). Let $U$ be an $A$-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$. Suppose $x_{0} \in B_{1 / 2}^{\prime}$. Let $0<\kappa<\kappa_{0}$ with a fixed $\kappa_{0}>0$. Then, for $0<t<t_{0}=t_{0}\left(n, \alpha, \kappa_{0}, M\right)$,

$$
\begin{aligned}
\frac{d}{d t} W_{\kappa}\left(t, u_{x_{0}}^{*}, 0\right) & \geq \frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-2}} \int_{\partial B_{t}}\left(\partial_{\nu} u_{x_{0}}^{*}-\frac{\kappa\left(1-b t^{\alpha}\right)}{t} u_{x_{0}}^{*}\right)^{2}, \\
\frac{d}{d t} W_{\kappa}^{A}\left(t, U_{x_{0}}^{*}, x_{0}\right) & \geq \frac{\mathrm{e}^{a t^{\alpha}}}{t^{n+2 \kappa-2}} \int_{\partial E_{t}\left(x_{0}\right)}\left(\left\langle\mathfrak{a}_{x_{0}} \nabla U_{x_{0}}^{*}, \nu\right\rangle-\frac{\kappa\left(1-b t^{\alpha}\right)}{t} U_{x_{0}}^{*}\right)^{2} \mu_{x_{0}}\left(x-x_{0}\right) .
\end{aligned}
$$

In particular, $W_{\kappa}\left(t, u_{x_{0}}^{*}, 0\right)$ and $W_{\kappa}^{A}\left(t, U_{x_{0}}^{*}, x_{0}\right)$ are nondecreasing in $t$ for $0<t<t_{0}$.
Proof. We note that the proof of Theorem 1.5.1 for the monotonicity of $W_{\kappa}\left(t, v, x_{0}\right)$ requires the function $v$ to be an almost minimizer for the Signorini problem for the monotonicity of its energy. However, it is not hard to see that the almost minimizing property of $v$ is used only when it is compared with the $\kappa$-homogeneous replacement $w$ of $v$ on balls centered at the given point $x_{0}$ to obtain

$$
\int_{B_{t}\left(x_{0}\right)}|\nabla w|^{2} \geq \frac{1}{1+t^{\alpha}} \int_{B_{t}\left(x_{0}\right)}|\nabla v|^{2},
$$

see (1.5.2). This means that the argument in the proof of Theorem 1.5.1 also works in our case and implies the part of the theorem for $u_{x_{0}}^{*}$. We note that the constants $a_{\kappa}$ and $b$ in our case will have an additional factor of $M$, as we work with $\omega(r)=M r^{\alpha}$ rather than $\omega(r)=r^{\alpha}$ in our case, but this change of the constants can be easily traced.

The part of the theorem for $U_{x_{0}}^{*}$ follows by a change of variables.
The families of monotonicity formulas $\left\{W_{\kappa}\right\}_{0<\kappa<\kappa_{0}}$ and $\left\{W_{\kappa}^{A}\right\}_{0<\kappa<\kappa_{0}}$ have an important feature that their intervals of monotonicity and the constant $b$ can be taken the same for all $0<\kappa<\kappa_{0}$. Because of that, their monotonicity indirectly implies that of another important quantity that we describe below. Namely, recall that for a function $v$ in $B_{r}\left(x_{0}\right)$, Almgren's frequency of $v$ at $x_{0}$ is defined as

$$
N\left(t, v, x_{0}\right):=\frac{t \int_{B_{t}\left(x_{0}\right)}|\nabla v|^{2}}{\int_{\partial B_{t}\left(x_{0}\right)} v^{2}}, \quad 0<t<r .
$$

Note that this quantity is well-defined when $v$ has an almost Signorini property at $x_{0}$ and $x_{0} \in \Gamma(v)$, since vanishing of $\int_{\partial B_{t}\left(x_{0}\right)} v^{2}$ for any $t>0$, would imply vanishing of $v$ in $B_{t}\left(x_{0}\right)$ by taking 0 as a competitor and consequently that $x_{0} \notin \Gamma(v)$.

Next consider a modification of $N$, which we call the truncated frequency:

$$
\widehat{N}_{\kappa_{0}}\left(t, v, x_{0}\right):=\min \left\{\frac{1}{1-b t^{\alpha}} N\left(t, v, x_{0}\right), \kappa_{0}\right\},
$$

where $b$ is as in Weiss-type monotonicity formulas for $\kappa<\kappa_{0}$. We next define the appropriate version of $N, \widehat{N}_{\kappa_{0}}$ in our setting. For a function $V$ in $E_{r}\left(x_{0}\right)$, we define

$$
\begin{aligned}
& N^{A}\left(t, V, x_{0}\right):=N\left(t, v_{x_{0}}, 0\right), \\
& \widehat{N}_{\kappa_{0}}^{A}\left(t, V, x_{0}\right):=\widehat{N}_{\kappa_{0}}\left(t, v_{x_{0}}, 0\right),
\end{aligned}
$$

for $0<t<r$, where $v_{x_{0}}=V \circ \bar{T}_{x_{0}}^{-1}$. More explicitly, we have

$$
\begin{aligned}
N^{A}\left(t, V, x_{0}\right) & :=\frac{t \int_{E_{t}\left(x_{0}\right)}\left\langle A\left(x_{0}\right) \nabla V, \nabla V\right\rangle}{\int_{\partial E_{t}\left(x_{0}\right)} V^{2} \mu_{x_{0}}\left(x-x_{0}\right)}, \\
\widehat{N}_{\kappa_{0}}^{A}\left(t, V, x_{0}\right) & :=\min \left\{\frac{1}{1-b t^{\alpha}} N^{A}\left(t, V, x_{0}\right), \kappa_{0}\right\} .
\end{aligned}
$$

As observed in Theorem 1.5.4, the Weiss-type monotonicity formula implies the following monotonicity of $\widehat{N}_{\kappa_{0}}^{A}$.

Theorem 2.7.2 (Almgren-type monotonicity formula). Let $U$, $\kappa_{0}$, and $t_{0}$ be as in Theorem 2.7.1, and $x_{0} \in B_{1 / 2}^{\prime}$ a free boundary point. Then

$$
t \mapsto \widehat{N}_{\kappa_{0}}^{A}\left(t, U_{x_{0}}^{*}, x_{0}\right)=\widehat{N}_{\kappa_{0}}\left(t, u_{x_{0}}^{*}, 0\right)
$$

is nondecreasing for $0<t<t_{0}$.

Definition 2.7.1 (Almgren's frequency at free boundary point). For an A-quasisymmetric almost minimizer $U$ of the $A$-Signorini problem in $B_{1}$ and $x_{0} \in \Gamma(U)$ let

$$
\kappa\left(x_{0}\right):=\widehat{N}_{\kappa_{0}}^{A}\left(0+, U_{x_{0}}^{*}, x_{0}\right)=\widehat{N}_{\kappa_{0}}\left(0+, u_{x_{0}}^{*}, 0\right) .
$$

We call $\kappa\left(x_{0}\right)$ Almgren's frequency at $x_{0}$.
Remark 2.7.3. Note that even though the monotonicity of the truncated frequency is stated in Theorem 2.7.2 only for $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(U)$, by a simple recentering and a scaling argument, it will be monotone also at all $x_{0} \in \Gamma(U)$, but for a possibly shorter interval of values $0<t<t_{0}\left(x_{0}\right)$ depending on $x_{0}$. Thus, $\kappa\left(x_{0}\right)$ exists at all $x_{0} \in \Gamma(U)$.

Further note that when $\kappa\left(x_{0}\right)<\kappa_{0}$, then $\widehat{N}_{\kappa_{0}}^{A}\left(t, U_{x_{0}}^{*}, x_{0}\right)=\frac{1}{1-b t^{\alpha}} N^{A}\left(t, U_{x_{0}}^{*}, x_{0}\right)$ for small $t$ and therefore

$$
\kappa\left(x_{0}\right)=N^{A}\left(0+, U_{x_{0}}^{*}, x_{0}\right),
$$

which means that it will not change if we replace $\kappa_{0}$ with a larger value.

### 2.8 Almgren rescalings and blowups

Our analysis of the free boundary is based on the analysis of blowups, which are the limits of rescalings of the solutions at free boundary points. In Signorini problem, there are a few types of rescalings that use different normalizations. In this section, we look at so-called Almgren rescalings and blowups that play well with the Almgren frequency formula.

Let $V \in W^{1,2}\left(B_{1}\right)$ and $x_{0} \in B_{1 / 2}^{\prime}$ be a free boundary point. For small $r>0$ define the Almgren rescaling of $V$ at $x_{0}$ by

$$
V_{x_{0}, r}^{A}(x):=\frac{V\left(r x+x_{0}\right)}{\left(\frac{1}{r^{n-1}} \int_{\partial E_{r}\left(x_{0}\right)} V^{2} \mu_{x_{0}}\left(x-x_{0}\right)\right)^{1 / 2}}
$$

The Almgren rescalings have the following normalization and scaling properties

$$
\begin{aligned}
& \left\|V_{x_{0}, r}^{A}\right\|_{L^{2}\left(\mathfrak{a}_{x_{0}} \partial B_{1}\right)}=1 \\
& N^{A\left(x_{0}\right)}\left(\rho, V_{x_{0}, r}^{A}, 0\right)=N^{A}\left(\rho r, V, x_{0}\right)
\end{aligned}
$$

Here $N^{A\left(x_{0}\right)}$ denotes Almgren's frequency for a constant matrix $A\left(x_{0}\right)$. Thus, we also have $N^{A}\left(r, V, x_{0}\right)=N^{A\left(x_{0}\right)}\left(r, V, x_{0}\right)$. Note that when $A=I$, then

$$
V_{x_{0}, r}^{I}=\frac{V\left(r x+x_{0}\right)}{\left(\frac{1}{r^{n-1}} \int_{\partial B_{r}\left(x_{0}\right)} V^{2}\right)^{1 / 2}}
$$

is same as the Almgren rescaling in Chapter 1, and satisfies

$$
\begin{aligned}
& \left\|V_{x_{0}, r}^{I}\right\|_{L^{2}\left(\partial B_{1}\right)}=1 \\
& N\left(\rho, V_{x_{0}, r}^{I}, 0\right)=N\left(\rho r, V, x_{0}\right)
\end{aligned}
$$

We will call the limits of $V_{x_{0}, r}^{A}$ over any subsequence $r=r_{\mathrm{j}} \rightarrow 0+$ Almgren blowups of $V$ at $x_{0}$ and denote them by $V_{x_{0}, 0}^{A}$.

By using a change of variables, we can express Almgren rescalings of $V$ in terms of those of $v_{x_{0}}=V \circ \bar{T}_{x_{0}}^{-1}$ and vice versa. Namely, we have

$$
\left(v_{x_{0}}\right)_{r}^{I}(y)=\left(\operatorname{det} \mathfrak{a}_{x_{0}}\right)^{1 / 2} V_{x_{0}, r}^{A}\left(\overline{\mathfrak{a}}_{x_{0}} y\right)
$$

wherever they are defined. Applied to the particular case $V=U_{x_{0}}^{*}$, we have

$$
\left(u_{x_{0}}^{*}\right)_{r}^{I}(y)=\left(\operatorname{det} \mathfrak{a}_{x_{0}}\right)^{1 / 2}\left(U_{x_{0}}^{*}\right)_{x_{0}, r}^{A}\left(\overline{\mathfrak{a}}_{x_{0}} y\right) .
$$

Proposition 2.8.1 (Existence of Almgren blowups). Let $U$ be an $A$-quasisymmetric almost minimizer for the A-Signorini problem in $B_{1}$, and $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(U)$ be such that $\kappa\left(x_{0}\right)<\kappa_{0}$. Then, every sequence of Almgren rescalings $\left(U_{x_{0}}^{*}\right)_{x_{0}, t_{\mathrm{j}}}^{A}$, with $t_{\mathrm{j}} \rightarrow 0+$, contains a subsequence, sill denoted $t_{\mathrm{j}}$ such that for a function $\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{A} \in C_{\mathrm{loc}}^{1}\left(\mathfrak{a}_{x_{0}}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)\right)$

$$
\left(U_{x_{0}}^{*}\right)_{x_{0}, t_{\mathrm{j}}}^{A} \rightarrow\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{A} \quad \text { in } C_{\mathrm{loc}}^{1}\left(\mathfrak{a}_{x_{0}}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)\right)
$$

Moreover, $\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{A}$ extends to a nonzero solution of the $A\left(x_{0}\right)$-Signorini problem in $\mathbb{R}^{n}$, $\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{A}(x)=\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{A}\left(P_{x_{0}} x\right)$, and it is homogeneous of degree $\kappa\left(x_{0}\right)$ in $\mathbb{R}^{n}$.

Similarly, every sequence of Almgren rescalings $\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I}$, with $t_{\mathrm{j}} \rightarrow 0+$ contains a subsequence, sill denoted $t_{\mathrm{j}}$ such that for a function $\left(u_{x_{0}}^{*}\right)_{0}^{I} \in C_{\mathrm{loc}}^{1}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)$

$$
\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I} \rightarrow\left(u_{x_{0}}^{*}\right)_{0}^{I} \quad \text { in } C_{\mathrm{loc}}^{1}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)
$$

Moreover, $\left(u_{x_{0}}^{*}\right)_{0}^{I}$ extends to a nonzero solution of the Signorini problem in $\mathbb{R}^{n}$, even in $y_{n}$, and it is homogeneous of degree $\kappa\left(x_{0}\right)$ in $\mathbb{R}^{n}$.

Proof. Step 1. Since $\kappa\left(x_{0}\right)<\kappa_{0}$, we must have $N\left(t, u_{x_{0}}^{*}, 0\right)<\kappa_{0}$ for small $t>0$. Then, for such $t$

$$
\int_{B_{1}}\left|\nabla\left(u_{x_{0}}^{*}\right)_{t}^{I}\right|^{2}=N\left(1,\left(u_{x_{0}}^{*}\right)_{t}^{I}, 0\right)=N\left(t, u_{x_{0}}^{*}, 0\right) \leq \kappa_{0},
$$

and combined with the normalization $\int_{\partial B_{1}}\left(\left(u_{x_{0}}^{*}\right)_{t}^{I}\right)^{2}=1$, we see that the family $\left(u_{x_{0}}^{*}\right)_{t}^{I}$ is bounded in $W^{1,2}\left(B_{1}\right)$, for small $t>0$. Hence, for any sequence $t_{\mathrm{j}} \rightarrow 0+$, there is a function $\left(u_{x_{0}}^{*}\right)_{0}^{I} \in W^{1,2}\left(B_{1}\right)$ such that, over a subsequence,

$$
\begin{array}{ll}
\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I} \rightarrow\left(u_{x_{0}}^{*}\right)_{0}^{I} & \text { weakly in } W^{1,2}\left(B_{1}\right), \\
\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I} \rightarrow\left(u_{x_{0}}^{*}\right)_{0}^{I} & \text { strongly in } L^{2}\left(\partial B_{1}\right) .
\end{array}
$$

In particular, $\int_{\partial B_{1}}\left(\left(u_{x_{0}}^{*}\right)_{0}^{I}\right)^{2}=1$, implying that $\left(u_{x_{0}}^{*}\right)_{0}^{I} \not \equiv 0$ in $B_{1}$.
Step 2. For $0<t<1$ and $x \in B_{1 /(2 t)}\left(x_{0}\right)$, let

$$
U_{x_{0}, t}(x)=U\left(x_{0}+t\left(x-x_{0}\right)\right), \quad A_{x_{0}, t}(x)=A\left(x_{0}+t\left(x-x_{0}\right)\right) .
$$

Then by a simple scaling argument, we have that $U_{x_{0}, t}$ is an almost minimizer of the $A_{x_{0}, t^{-}}$ Signorini problem in $B_{1 /(2 t)}\left(x_{0}\right)$ with a gauge function $\mu_{t}(r)=(t r)^{\alpha} \leq r^{\alpha}$. In particular, for any $R>0$, we will have that $U_{x_{0}, t} \in C^{1, \beta}\left(E_{R}^{ \pm}\left(x_{0}\right) \cup E_{R}^{\prime}\left(x_{0}\right)\right)$ for $0<t<t(R, M)$ with

$$
\left\|U_{x_{0}, t}\right\|_{C^{1, \beta}(K)} \leq C\left\|U_{x_{0}, t}\right\|_{W^{1,2}\left(E_{R}\left(x_{0}\right)\right)}
$$

with $C=C(n, \alpha, M, R, K)$, for any $K \Subset E_{R}^{ \pm}\left(x_{0}\right) \cup E_{R}^{\prime}\left(x_{0}\right)$. Then, arguing as in the proof of Theorem 2.6.5, by using the quasisymmetry of $U$, we obtain that

$$
\left\|\left(U_{x_{0}, t}\right)_{x_{0}}^{*}\right\|_{C^{1, \beta}(K)} \leq C\left\|\left(U_{x_{0}, t}\right)_{x_{0}}^{*}\right\|_{W^{1,2}\left(E_{R}\left(x_{0}\right)\right)}
$$

where

$$
\left(U_{x_{0}, t}\right)_{x_{0}}^{*}(x)=\frac{U_{x_{0}, t}(x)+U_{x_{0}, t}\left(P_{x_{0}} x\right)}{2} .
$$

Next, observing that $\left(u_{x_{0}}^{*}\right)_{t}^{I}$ is a positive constant multiple of $\left(U_{x_{0}, t}\right)_{x_{0}}^{*} \circ \bar{T}_{x_{0}}^{-1}$, we obtain that

$$
\left\|\left(u_{x_{0}}^{*}\right)_{t}^{I}\right\|_{C^{1, \beta}(K)} \leq C\left\|\left(u_{x_{0}}^{*}\right)_{t}^{I}\right\|_{W^{1,2}\left(B_{R}\right)},
$$

for any $K \Subset B_{R}^{ \pm} \cup B_{R}^{\prime}$. Taking $R=1$, combined with the boundedness of $\left(u_{x_{0}}^{*}\right)_{t}^{I}$ in $W^{1,2}\left(B_{1}\right)$ for small $t>0$, it follows that up to a subsequence,

$$
\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I} \rightarrow\left(u_{x_{0}}^{*}\right)_{0}^{I} \quad \text { in } C_{\mathrm{loc}}^{1}\left(B_{1}^{ \pm} \cup B_{1}^{\prime}\right)
$$

Step 3. Next, we claim that the blowup $\left(u_{x_{0}}^{*}\right)_{0}^{I}$ is a solution of the Signorini problem in $B_{1}$. Indeed, fix $0<R<1$, and for each $t_{\mathrm{j}}$ let $h_{t_{\mathrm{j}}}$ be the Signorini replacement of $\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I}$ in $B_{R}$. Then a first variation argument gives (see (1.3.2))

$$
\int_{B_{R}}\left\langle\nabla h_{t_{\mathrm{j}}}, \nabla\left(\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I}-h_{t_{\mathrm{j}}}\right)\right\rangle \geq 0
$$

Since $\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I}$ has an almost Signorini property at 0 with a gauge function $r \mapsto C\left(t_{\mathrm{j}} r\right)^{\alpha}$, it follows that

$$
\int_{B_{R}}\left|\nabla\left(\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I}-h_{t_{\mathrm{j}}}\right)\right|^{2} \leq C\left(R t_{\mathrm{j}}\right)^{\alpha} \int_{B_{R}}\left|\nabla\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I}\right|^{2}
$$

This implies that $h_{t_{\mathrm{j}}} \rightarrow\left(u_{x_{0}}^{*}\right)_{0}^{I}$ weakly in $W^{1,2}\left(B_{R}\right)$. On the other hand, by the boundedness of the sequence $h_{t_{\mathrm{j}}}$ in $W^{1,2}\left(B_{R}\right)$, we have also boundedness in $C^{1,1 / 2}$ norm locally in $\left(B_{R}^{ \pm} \cup B_{R}^{\prime}\right)$
and hence, over a subsequence, $h_{t_{\mathrm{j}}} \rightarrow\left(u_{x_{0}}^{*}\right)_{0}^{I}$ in $C_{\mathrm{loc}}^{1}\left(B_{R}^{ \pm} \cup B_{R}^{\prime}\right)$. By this convergence, we then conclude that $\left(u_{x_{0}}^{*}\right)_{0}^{I}$ satisfies

$$
\begin{aligned}
\Delta\left(u_{x_{0}}^{*}\right)_{0}^{I}=0 & \text { in } B_{R} \backslash B_{R}^{\prime} \\
\left(u_{x_{0}}^{*}\right)_{0}^{I} \geq 0, \quad-\partial_{y_{n}}^{+}\left(u_{x_{0}}^{*}\right)_{0}^{I} \geq 0, \quad\left(u_{x_{0}}^{*}\right)_{0}^{I} \partial_{y_{n}}^{+}\left(u_{x_{0}}^{*}\right)_{0}^{I}=0 & \text { on } B_{R}^{\prime}
\end{aligned}
$$

and hence, by letting $R \rightarrow 1,\left(u_{x_{0}}^{*}\right)_{0}^{I}$ itself solves the Signorini problem in $B_{1}$.
Step 4. Recall now that the blowup $\left(u_{x_{0}}^{*} I_{0}^{I}\right.$ is nonzero in $B_{1}$. In particular, $\int_{\partial B_{r}}\left(\left(u_{x_{0}}^{*}\right)_{0}^{I}\right)^{2}>0$ for any $0<r<1$, otherwise we would have that $\left(u_{x_{0}}^{*}\right)_{0}^{I}$ is identically zero on $\partial B_{r}$ and consequently also on $B_{r}$. Using this fact, combined with $C_{\text {loc }}^{1}$ convergence in $B_{1}^{ \pm} \cup B_{1}^{\prime}$, we have that for any $0<r<1$

$$
\begin{aligned}
N\left(r,\left(u_{x_{0}}^{*}\right)_{0}^{I}, 0\right) & =\lim _{t_{\mathrm{j}} \rightarrow 0} N\left(r,\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I}, 0\right)=\lim _{t_{\mathrm{j}} \rightarrow 0} N\left(r t_{\mathrm{j}}, u_{x_{0}}^{*}, 0\right) \\
& =N\left(0+, u_{x_{0}}^{*}, 0\right)=\kappa\left(x_{0}\right) .
\end{aligned}
$$

Thus, Almgren's frequency of $\left(u_{x_{0}}^{*}\right)_{0}^{I}$ is constant $\kappa\left(x_{0}\right)$ on $0<r<1$ which is possible only if $\left(u_{x_{0}}^{*}\right)_{0}^{I}$ is a $\kappa\left(x_{0}\right)$-homogeneous solution of the Signorini problem in $B_{1}$, see Theorem 9.4 in [48]. Finally, by using the homogeneity, we readily extend $\left(u_{x_{0}}^{*}\right)_{0}^{I}$ to a solution of the Signorini problem in all of $\mathbb{R}^{n}$. This completes the proof for $\left(u_{x_{0}}^{*}\right)_{0}^{I}$.

The corresponding result for $\left(U_{x_{0}}^{*}\right)_{x_{0}, t_{\mathrm{j}}}^{A}$ follows now by a change of variables.
With Proposition 2.8.1 at hand, we can repeat the argument in the proof of Lemma 1.6.1 with $u_{x_{0}}^{*}$ to obtain the following, which is possible since $u_{x_{0}}^{*}$ satisfies the complementarity condition and an Almgren-type monotonicity formula with a blowup as a nonzero solution of the Signorini problem.

Lemma 2.8.1 (Minimal frequency). Let $U$ be an $A$-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$. If $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(U)$, then

$$
\kappa\left(x_{0}\right) \geq \frac{3}{2} .
$$

Consequently, we also have

$$
\widehat{N}_{\kappa_{0}}^{A}\left(t, U_{x_{0}}^{*}, x_{0}\right)=\widehat{N}_{\kappa_{0}}\left(t, u_{x_{0}}^{*}, 0\right) \geq 3 / 2 \quad \text { for } 0<t<t_{0} .
$$

Lemma 2.8.1 readily gives the following. (see Corollary 2.8.2)
Corollary 2.8.2. Let $U$ be an $A$-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$ and $x_{0}$ a free boundary point. Then

$$
W_{3 / 2}^{A}\left(t, U_{x_{0}}^{*}, x_{0}\right)=\operatorname{det} \mathfrak{a}_{x_{0}} W_{3 / 2}\left(t, u_{x_{0}}^{*}, 0\right) \geq 0, \quad \text { for } 0<t<t_{0}
$$

### 2.9 Growth estimates

The first result in this section (Lemma 2.9.1) provides growth estimates for the quasisymmetric almost minimizers near free boundary points $x_{0}$ with $\kappa\left(x_{0}\right) \geq \kappa$. Such estimates were obtained in Lemma 1.7.1 in the case $A \equiv I$ as a consequence of Weiss-type monotonicity formulas. However, they contain an unwanted logarithmic term that creates difficulties in the blowup analysis of the problem.

The next two results (Lemmas 2.9.2 and 2.9.3) remove the logarithmic term from these estimates for $\kappa=3 / 2$, by establishing first a growth rate for $W_{3 / 2}$. (Recall that $\kappa\left(x_{0}\right) \geq 3 / 2$ at every free boundary point $x_{0}$, by Lemma 2.8.1.) These are analogous to Lemmas 1.7.3, 1.7.4 in the case $A \equiv I$ and follow from the so-called epiperimetric inequality for $\kappa=3 / 2$ (see e.g. Themrem 1.7.2). Later, in Section 2.12, we remove the logarithmic term also in the case $\kappa=2 m<\kappa_{0}, m \in \mathbb{N}$, see Lemma 2.12.1.

The results in this section are stated in terms of both $u_{x_{0}}^{*}$ and $U_{x_{0}}^{*}$, as we need both forms in the subsequent arguments. We note that the estimates for $u_{x_{0}}^{*}$ follow directly from Lemmas 1.7.1, 1.7.3, 1.7.4 and the ones for $U_{x_{0}}^{*}$ are obtained by using the deskewing procedure and therefore we skip all proofs in this section.

In the estimates below, as well in the rest of the chapter, we use the notation

$$
R_{0}:=(1 / 2) \Lambda^{-1 / 2}
$$

which is the radius of the largest ball $B_{R_{0}}$, where $u_{x_{0}}^{*}$ is guaranteed to exists for any $x_{0} \in B_{1 / 2}$ for an almost minimizer $U$ in $B_{1}$.

Lemma 2.9.1 (Weak growth estimate). Let $U$ be an A-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$ and $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(U)$. If

$$
\kappa\left(x_{0}\right) \geq \kappa
$$

for some $\kappa \leq \kappa_{0}$, then

$$
\begin{aligned}
\int_{\partial B_{t}}\left(u_{x_{0}}^{*}\right)^{2} & \leq C\left\|u_{x_{0}}^{*}\right\|_{W^{1,2}\left(B_{R_{0}}\right)}^{2}\left(\log \frac{1}{t}\right) t^{n+2 \kappa-1}, \\
\int_{B_{t}}\left|\nabla u_{x_{0}}^{*}\right|^{2} & \leq C\left\|u_{x_{0}}^{*}\right\|_{W^{1,2}\left(B_{R_{0}}\right)}^{2}\left(\log \frac{1}{t}\right) t^{n+2 \kappa-2}, \\
\int_{\partial E_{t}\left(x_{0}\right)}\left(U_{x_{0}}^{*}\right)^{2} & \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}^{2}\left(\log \frac{1}{t}\right) t^{n+2 \kappa-1}, \\
\int_{E_{t}\left(x_{0}\right)}\left|\nabla U_{x_{0}}^{*}\right|^{2} & \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}^{2}\left(\log \frac{1}{t}\right) t^{n+2 \kappa-2},
\end{aligned}
$$

for $0<t<t_{0}=t_{0}\left(n, \alpha, M, \kappa_{0}\right)$ and $C=C\left(n, \alpha, M, \kappa_{0}\right)$.
Lemma 2.9.2. Let $U$ and $x_{0}$ be as above. Then, there exists $\delta=\delta(n, \alpha)>0$ such that

$$
\begin{gathered}
\left.0 \leq W_{3 / 2}\left(t, u_{x_{0}}^{*}, 0\right) \leq C\left\|u_{x_{0}}^{*}\right\|_{W^{1,2}\left(B_{R_{0}}\right.}^{2}\right)^{\delta} \\
0 \leq W_{3 / 2}^{A}\left(t, U_{x_{0}}^{*}, x_{0}\right) \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} t^{\delta}
\end{gathered}
$$

for $0<t<t_{0}=t_{0}(n, \alpha, M)$ and $C=C(n, \alpha, M)$.
Lemma 2.9.3 (Optimal growth estimate). Let $U$ and $x_{0}$ be as above. Then,

$$
\begin{aligned}
& \int_{\partial B_{t}}\left(u_{x_{0}}^{*}\right)^{2} \leq C\left\|u_{x_{0}}^{*}\right\|_{W^{1,2}\left(B_{R_{0}}\right)}^{2} t^{n+2} \\
& \int_{B_{t}}\left|\nabla u_{x_{0}}^{*}\right|^{2} \leq C\left\|u_{x_{0}}^{*}\right\|_{W^{1,2}\left(B_{R_{0}}\right)}^{2} t^{n+1} \\
& \int_{\partial E_{t}\left(x_{0}\right)}\left(U_{x_{0}}^{*}\right)^{2} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} t^{n+2} \\
& \int_{E_{t}\left(x_{0}\right)}\left|\nabla U_{x_{0}}^{*}\right|^{2} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}^{2} t^{n+1}
\end{aligned}
$$

for $0<t<t_{0}=t_{0}(n, \alpha, M)$ and $C=C(n, \alpha, M)$.

### 2.10 3/2-almost homogeneous rescalings and blowups

In this section we study another kind of rescalings and blowups that will play a fundamental role in the analysis of regular free boundary points where $\kappa\left(x_{0}\right)=3 / 2$ (see the next section), namely $3 / 2$-almost homogeneous blowups. The main result that we prove in this section is the uniqueness and Hölder continuous dependence of such blowups at a free boundary point $x_{0}$ (Lemma 2.10.3).

For a function $v$ in $B_{1}$ and $x_{0} \in B_{1 / 2}^{\prime}$, we define the 3/2-almost homogeneous rescalings of $v$ at $x_{0}$ by

$$
v_{x_{0}, t}^{\phi}(x)=\frac{v\left(t x+x_{0}\right)}{\phi(t)}, \quad \phi(t)=\mathrm{e}^{-\left(\frac{3 b}{2 \alpha}\right) t^{\alpha}} t^{3 / 2}
$$

with $b$ as in the Weiss-type monotonicity formulas $W_{3 / 2}^{A}$ and $W_{3 / 2}$. When $x_{0}=0$, we simply write $v_{0, t}^{\phi}=v_{t}^{\phi}$.

The name is explained by the fact that

$$
\lim _{t \rightarrow 0} \frac{\phi(t)}{t^{3 / 2}}=1
$$

and the reason to look at such rescalings instead of $3 / 2$-homogeneous rescalings (that would correspond to $\phi(t)=t^{3 / 2}$ ) is how they play well with the Weiss-type monotonicity formulas $W_{3 / 2}^{A}$ and $W_{3 / 2}$.

Now, if $U$ is an $A$-quasisymmetric almost minimizer and $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(U)$, then for any fixed $R>1$, if $t=t_{\mathrm{j}}>0$ is small, then by Lemma 2.9.3,

$$
\begin{gather*}
\int_{B_{R}}\left|\nabla\left(u_{x_{0}}^{*}\right)_{t}^{\phi}\right|^{2}=\frac{\mathrm{e}^{\frac{3 b}{\alpha} t^{\alpha}}}{t^{n+1}} \int_{B_{R t}}\left|\nabla u_{x_{0}}^{*}\right|^{2} \leq C\left\|u_{x_{0}}^{*}\right\|_{W^{1,2}\left(B_{R_{0}}\right)}^{2} R^{n+1},  \tag{2.10.1}\\
\int_{\partial B_{R}}\left(\left(u_{x_{0}}^{*}\right)_{t}^{\phi}\right)^{2}=\frac{\mathrm{e}^{\frac{3 b}{\alpha} t^{\alpha}}}{t^{n+2}} \int_{\partial B_{R t}}\left(u_{x_{0}}^{*}\right)^{2} \leq C\left\|u_{x_{0}}^{*}\right\|_{W^{1,2}\left(B_{R_{0}}\right)}^{2} R^{n+2}, \tag{2.10.2}
\end{gather*}
$$

with $C=C(n, \alpha, M), R_{0}=(1 / 2) \Lambda^{-1 / 2}$. Hence, $\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{\phi}$ is a bounded sequence in $W^{1,2}\left(B_{R}\right)$. Next, arguing as in the proof of Proposition 2.8.1, we will have that

$$
\begin{equation*}
\left\|\widehat{\nabla\left(u_{x_{0}}^{*}\right)_{t}^{\phi}}\right\|_{C^{0, \beta}(K)} \leq C\left\|\left(u_{x_{0}}^{*}\right)_{t}^{\phi}\right\|_{W^{1,2}\left(B_{R}\right)}, \tag{2.10.3}
\end{equation*}
$$

with $C=C(n, \alpha, M, R, K)$ for $K \Subset B_{R}$. Thus, by letting $R \rightarrow \infty$ and using Cantor's diagonal argument, we can conclude that over a subsequence $t=t_{\mathrm{j}} \rightarrow 0+$,

$$
\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{\phi} \rightarrow\left(u_{x_{0}}^{*}\right)_{0}^{\phi} \quad \text { in } \quad C_{\mathrm{loc}}^{1}\left(\mathbb{R}_{ \pm}^{n} \cup \mathbb{R}^{n-1}\right)
$$

We call such $\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ a $3 / 2$-homogeneous blowup of $u_{x_{0}}^{*}$ at 0 . (We may skip the "almost" modifier here as the limit is the same as for $3 / 2$-homogeneous rescalings.) Furthermore, from the relation

$$
\left(u_{x_{0}}^{*}\right)_{t}^{\phi}(y)=\left(U_{x_{0}}^{*}\right)_{x_{0}, t}^{\phi}\left(\overline{\mathfrak{a}}_{x_{0}} y\right),
$$

we also conclude that for any sequence $t_{\mathrm{j}} \rightarrow 0+$, there is a subsequence, still denoted by $t_{\mathrm{j}}$, such that

$$
\left(U_{x_{0}}^{*}\right)_{x_{0}, t_{\mathrm{j}}}^{\phi} \rightarrow\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi} \quad \text { in } \quad C_{\mathrm{loc}}^{1}\left(\mathbb{R}_{ \pm}^{n} \cup \mathbb{R}^{n-1}\right)
$$

Apriori, the blowups $\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ and $\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}$ may depend on the sequence $t_{\mathrm{j}} \rightarrow 0+$. However, this does not happen in the case of $3 / 2$-homogeneous blowups. We start with what we call a rotation estimate for rescalings.

Lemma 2.10.1 (Rotation estimate). Let $U$ be an $A$-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}, x_{0} \in B_{1 / 2}^{\prime}$ a free boundary point, and $\delta$ as in Lemma 2.9.2. Then,

$$
\begin{aligned}
\int_{\partial B_{1}}\left|\left(u_{x_{0}}^{*}\right)_{t}^{\phi}-\left(u_{x_{0}}^{*}\right)_{s}^{\phi}\right| & \leq C\left\|u_{x_{0}}^{*}\right\|_{W^{1,2}\left(B_{R_{0}}\right)} t^{\delta / 2} \\
\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, t}^{\phi}-\left(U_{x_{0}}^{*}\right)_{x_{0}, s}^{\phi}\right| & \leq C\|U\|_{W^{1,2}\left(B_{1}\right)} t^{\delta / 2}
\end{aligned}
$$

for $s<t<t_{0}=t_{0}(n, \alpha, M)$ and $C=C(n, \alpha, M)$.

Proof. This is an analogue of Lemma 1.8.2, which follows from the computation done in the proof of Lemma 1.7.1, the growth estimate for $W_{3 / 2}$ in Lemma 1.7.3 and a dyadic argument. The analogues of those results in our case are stated in Lemma 2.9.1 and 2.9.2. This proves the lemma for $u_{x_{0}}^{*}$. The estimate for $\left(U_{x_{0}}^{*}\right)_{x_{0}, t}^{\phi}$ then follows from the equality

$$
\left(u_{x_{0}}^{*}\right)_{t}^{\phi}(y)=\left(U_{x_{0}}^{*}\right)_{x_{0}, t}^{\phi}\left(\overline{\mathfrak{a}}_{x_{0}} y\right), \quad y \in B_{R_{0} / t} .
$$

The uniqueness of $3 / 2$-homogeneous blowup now follows.
Lemma 2.10.2. Let $\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}$ and $\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ be blowups of $\left(U_{x_{0}}^{*}\right)_{x_{0}, t}^{\phi}$ and $\left(u_{x_{0}}^{*}\right)_{t}^{\phi}$, respectively, at a free boundary point $x_{0} \in B_{1 / 2}^{\prime}$. Then,

$$
\begin{aligned}
\int_{\partial B_{1}}\left|\left(u_{x_{0}}^{*}\right)_{t}^{\phi}-\left(u_{x_{0}}^{*}\right)_{0}^{\phi}\right| \leq C\left\|u_{x_{0}}^{*}\right\|_{W^{1,2}\left(B_{R_{0}}\right.} t^{\delta / 2} \\
\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, t}^{\phi}-\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}\right| \leq C\|U\|_{W^{1,2}\left(B_{1}\right)} t^{\delta / 2}
\end{aligned}
$$

for $0<t<t_{0}(n, \alpha, M)$ and $C=C(n, \alpha, M)$, where $\delta=\delta(n, \alpha)>0$ is as in Lemma 2.10.1. In particular, the blowups $\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ and $\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}$ are unique.

Proof. If $\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ is the limit of $\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{\phi}$ for $t_{\mathrm{j}} \rightarrow 0$, then the first part of the lemma follows immediately from Lemma 2.10.1, by taking $s=t_{\mathrm{j}} \rightarrow 0$ and passing to the limit.

To see the uniqueness of blowups, we observe that $\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ is a solution of the Signorini problem in $B_{1}$, by arguing as in the proof of Proposition 2.8.1 for Almgren blowups. Now, if $v_{0}$ is another blowup, over a possibly different sequence $t_{\mathrm{j}} \rightarrow 0$, then passing to the limit in the first part of the lemma we will have

$$
\int_{\partial B_{1}}\left|v_{0}-\left(u_{x_{0}}^{*}\right)_{0}^{\phi}\right|^{2}=0
$$

implying that both $v_{0}$ and $\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ are solutions of the Signorini problem in $B_{1}$ with the same boundary values on $\partial B_{1}$. By the uniqueness of such solutions, we have $v_{0}=\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ in $B_{1}$.

The equality propagates to all of $\mathbb{R}^{n}$ by the unique continuation of harmonic functions in $\mathbb{R}_{ \pm}^{n}$. This completes the proof for $u_{x_{0}}^{*}$. An analogous argument holds for $U_{x_{0}}^{*}$ using the equalities

$$
\begin{array}{ll}
\left(u_{x_{0}}^{*}\right)_{t}^{\phi}(y)=\left(U_{x_{0}}^{*}\right)_{x_{0}, t}^{\phi}\left(\overline{\mathfrak{a}}_{x_{0}} y\right), & y \in B_{R_{0} / t}, \\
\left(u_{x_{0}}^{*}\right)_{0}^{\phi}(y)=\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}\left(\overline{\mathfrak{a}}_{x_{0}} y\right), & y \in \mathbb{R}^{n} .
\end{array}
$$

The rotation estimate for rescalings implies not only the uniqueness of blowups and the convergence rate to blowups, but also the continuous dependence of blowups on a free boundary point.

Lemma 2.10.3 (Continuous dependence of blowups). There exists $\rho=\rho(n, \alpha, M)>0$ such that if $x_{0}, y_{0} \in B_{\rho}$ are free boundary points of $U$, then

$$
\begin{array}{r}
\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right| \leq C \mid x_{0}-y_{0}^{\gamma}, \\
\int_{\partial B_{1}}\left|\left(u_{x_{0}}^{*}\right)_{0}^{\phi}-\left(u_{y_{0}}^{*}\right)_{0}^{\phi}\right| \leq C\left|x_{0}-y_{0}\right|^{\gamma}, \\
\int_{\partial B_{1}}\left|\left(u_{x_{0}}^{*}\right)_{0}^{\phi}-\left(u_{y_{0}}^{*}\right)_{0}^{\phi}\right| \leq C\left|x_{0}-y_{0}\right|^{\gamma}, \tag{2.10.6}
\end{array}
$$

with $C=C\left(n, \alpha, M,\|U\|_{W^{1,2}\left(B_{1}\right)}\right), \gamma=\gamma(n, \alpha, M)>0$.

Proof. Step 1. Let $d=\left|x_{0}-y_{0}\right|$ and $d^{\tau} \leq r \leq 2 d^{\tau}$ with $\tau=\tau(\alpha) \in(0,1)$ to be determined later.

Next note that we can incorporate the weight $\mu_{x_{0}} / \operatorname{det} \mathfrak{a}_{x_{0}}$ with $\mu_{x_{0}}$ as in (2.2.4) in the integral on the left hand side of (2.10.4) because of the bounds

$$
\left(\frac{\lambda}{\Lambda}\right)^{1 / 2} \leq \frac{\mu_{x_{0}}}{\operatorname{det} \mathfrak{a}_{x_{0}}} \leq\left(\frac{\Lambda}{\lambda}\right)^{1 / 2}
$$

Then, by using Lemma 2.10.2, we have

$$
\begin{align*}
& \int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right| \frac{\mu_{x_{0}}}{\operatorname{det} \mathfrak{a}_{x_{0}}} \\
& \leq \int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left(\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}-\left(U_{x_{0}}^{*}\right)_{x_{0}, r}^{\phi}\right|+\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, r}^{\phi}-\left(U_{x_{0}}^{*}\right)_{y_{0}, r}^{\phi}\right|\right. \\
&\left.\quad+\left|\left(U_{x_{0}}^{*}\right)_{y_{0}, r}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}\right|+\left|\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right|\right) \frac{\mu_{x_{0}}}{\operatorname{det} \mathfrak{a}_{x_{0}}} \\
& \quad+\int_{\mathfrak{a}_{y_{0}} \partial B_{1}}\left|\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right| \frac{\mu_{y_{0}}}{\operatorname{det}}  \tag{2.10.7}\\
& \quad \quad-\int_{\mathfrak{a}_{y_{0}} \partial B_{1}}\left|\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right| \frac{\mu_{y_{0}}}{\operatorname{det} \mathfrak{a}_{y_{0}}} \\
& \leq 2 C r^{\delta / 2}+I_{r}+I I_{r}+I I I_{r} \\
& \leq C d^{\tau \delta / 2}+I_{r}+I I_{r}+I I I_{r},
\end{align*}
$$

where

$$
\begin{aligned}
I_{r} & =\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, r}^{\phi}-\left(U_{x_{0}}^{*}\right)_{y_{0}, r}^{\phi}\right| \frac{\mu_{x_{0}}}{\operatorname{det} \mathfrak{a}_{x_{0}}}, \\
I I_{r} & =\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{y_{0}, r}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}\right| \frac{\mu_{x_{0}}}{\operatorname{det} \mathfrak{a}_{x_{0}}}, \\
I I I_{r} & =\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right| \frac{\mu_{x_{0}}}{\operatorname{det} \mathfrak{a}_{x_{0}}}-\int_{\mathfrak{a}_{y_{0}} \partial B_{1}}\left|\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right| \frac{\mu_{y_{0}}}{\operatorname{det} \mathfrak{a}_{y_{0}}} .
\end{aligned}
$$

Step 2. By the definition of the almost homogeneous rescalings, we have

$$
I_{r} \leq \frac{C}{d^{\tau(n+1 / 2)}} \int_{\mathfrak{a}_{x_{0}} \partial B_{r}}\left|U_{x_{0}}^{*}\left(z+x_{0}\right)-U_{x_{0}}^{*}\left(z+y_{0}\right)\right| d S_{z}
$$

This gives

$$
\begin{aligned}
\frac{1}{d^{\tau}} \int_{d^{\tau}}^{2 d^{\tau}} I_{r} d r & \leq \frac{C}{d^{\tau(n+3 / 2)}} \int_{d^{\tau}}^{2 d^{\tau}} \int_{\mathfrak{a}_{x_{0}} \partial B_{r}}\left|U_{x_{0}}^{*}\left(z+x_{0}\right)-U_{x_{0}}^{*}\left(z+y_{0}\right)\right| d S_{z} d r \\
& \leq \frac{C}{d^{\tau(n+3 / 2)}} \int_{\mathfrak{a}_{x_{0}}\left(B_{\left.2 d^{\tau} \backslash B_{d^{\tau}}\right)}\left|U_{x_{0}}^{*}\left(z+x_{0}\right)-U_{x_{0}}^{*}\left(z+y_{0}\right)\right| d z\right.} \quad=\frac{C}{d^{\tau(n+3 / 2)}} \int_{\mathfrak{a}_{x_{0}}\left(B_{\left.2 d^{\tau} \backslash B_{d^{\tau}}\right)}\left|\int_{0}^{1} \frac{d}{d s}\left[U_{x_{0}}^{*}\left(z+x_{0}(1-s)+y_{0} s\right)\right] d s\right| d z\right.} \\
& \leq \frac{C}{d^{\tau(n+3 / 2)}}\left|x_{0}-y_{0}\right| \int_{0}^{1} \int_{\mathfrak{a}_{x_{0}}\left(B_{\left.2 d^{\tau} \backslash B_{d} \tau\right)}\left|\nabla U_{x_{0}}^{*}\left(z+x_{0}(1-s)+y_{0} s\right)\right| d z d s\right.} \\
& \leq \frac{C}{d^{\tau(n+3 / 2)-1}} \int_{0}^{1} \int_{\mathfrak{a}_{x_{0}} B_{2 d^{\tau}}+\left[x_{0}(1-s)+y_{0} s\right]}\left|\nabla U_{x_{0}}^{*}\right| d z d s .
\end{aligned}
$$

Notice that the last integral is taken over

$$
\begin{aligned}
\mathfrak{a}_{x_{0}} B_{2 d^{\tau}}+\left[x_{0}(1-s)+y_{0} s\right] & =\mathfrak{a}_{x_{0}}\left[B_{2 d^{\tau}}+s \mathfrak{a}_{x_{0}}^{-1}\left(y_{0}-x_{0}\right)\right]+x_{0} \\
& \subset \mathfrak{a}_{x_{0}} B_{2 d^{\tau}+\lambda^{-1 / 2} d}+x_{0} \subset E_{3 d^{\tau}}\left(x_{0}\right),
\end{aligned}
$$

if $\rho=\rho(n, \alpha, M)$ is small so that $(2 \rho)^{1-\tau} \leq \lambda^{1 / 2}$ which readily implies $d^{1-\tau} \leq \lambda^{1 / 2}$. Thus,

$$
\begin{aligned}
\frac{1}{d^{\tau}} \int_{d^{\tau}}^{2 d^{\tau}} I_{r} d r & \leq \frac{C}{d^{\tau(n+3 / 2)-1}} \int_{0}^{1} \int_{E_{3 d^{\tau}\left(x_{0}\right)}}\left|\nabla U_{x_{0}}^{*}\right| d z d s \\
& \leq \frac{C}{d^{\tau(n / 2+3 / 2)-1}}\left(\int_{E_{3 d^{\tau}\left(x_{0}\right)}}\left|\nabla U_{x_{0}}^{*}\right|^{2}\right)^{1 / 2} \\
& \leq C\|U\|_{W^{1,2}\left(B_{1}\right)} d^{1-\tau}
\end{aligned}
$$

where the third inequality follows from Lemma 2.9.3.
Step 3. By the definition of rescalings and symmetrizations, we have

$$
\begin{aligned}
I I_{r} & \leq \frac{C}{d^{\tau(n+1 / 2)}} \int_{\mathfrak{a}_{x_{0}} \partial B_{r}+y_{0}}\left|U_{x_{0}}^{*}(z)-U_{y_{0}}^{*}(z)\right| d S_{z} \\
& \leq \frac{C}{d^{\tau(n+1 / 2)}} \int_{\mathfrak{a}_{x_{0}} \partial B_{r}+y_{0}}\left|U\left(P_{x_{0}} z\right)-U\left(P_{y_{0}} z\right)\right| d S_{z} .
\end{aligned}
$$

This gives

$$
\begin{aligned}
\frac{1}{d^{\tau}} \int_{d^{\tau}}^{2 d^{\tau}} I I_{r} d r & \leq \frac{C}{d^{\tau(n+3 / 2)}} \int_{\mathfrak{a}_{x_{0}}\left(B_{2 d} \tau \backslash B_{d} \tau\right)+y_{0}}\left|U\left(P_{x_{0}} z\right)-U\left(P_{y_{0}} z\right)\right| d z \\
& \leq \frac{C}{d^{\tau(n+3 / 2)}} \int_{\mathfrak{a}_{x_{0}}\left(B_{2 d} \tau \backslash B_{d^{\tau}}\right)+y_{0}} \int_{0}^{1}\left|\frac{d}{d s}\left[U\left(\left[(1-s) P_{x_{0}}+s P_{y_{0}}\right] z\right)\right]\right| d s d z \\
& \leq \frac{C\left|P_{x_{0}}-P_{y_{0}}\right|}{d^{\tau(n+3 / 2)}} \int_{0}^{1} \int_{\mathfrak{a}_{x_{0}}\left(B_{2 d} \tau \backslash B_{d} \tau\right)+y_{0}}\left|\nabla U\left(\left[(1-s) P_{x_{0}}+s P_{y_{0}}\right] z\right)\right| d z d s .
\end{aligned}
$$

Now we do the change of variables

$$
y=\left[(1-s) P_{x_{0}}+s P_{y_{0}}\right] z .
$$

Since $P_{x_{0}}$ and $P_{y_{0}}$ are upper-triangular matrices with diagonal entries $1,1, \ldots, 1,-1$, so is $(1-s) P_{x_{0}}+s P_{y_{0}}$. Thus

$$
\left|\operatorname{det}\left[(1-s) P_{x_{0}}+s P_{y_{0}}\right]\right|=1
$$

Moreover, $y \in\left[(1-s) P_{x_{0}}+s P_{y_{0}}\right]\left(\mathfrak{a}_{x_{0}} B_{2 d^{\tau}}+y_{0}\right)$. Since

$$
\mathfrak{a}_{x_{0}} B_{2 d^{\tau}}+y_{0} \subset \mathfrak{a}_{y_{0}} B_{2(\Lambda / \lambda)^{1 / 2} d^{\tau}}+y_{0}=E_{2(\Lambda / \lambda)^{1 / 2} d^{\tau}}\left(y_{0}\right),
$$

we have

$$
P_{y_{0}}\left(\mathfrak{a}_{x_{0}} B_{2 d^{\tau}}+y_{0}\right) \subset P_{y_{0}} E_{2(\Lambda / \lambda)^{1 / 2} d^{\tau}}\left(y_{0}\right)=E_{2(\Lambda / \lambda)^{1 / 2} d^{\tau}}\left(y_{0}\right) .
$$

Similarly, since

$$
\begin{aligned}
\mathfrak{a}_{x_{0}} B_{2 d^{\tau}}+y_{0} & =E_{2 d^{\tau}}\left(x_{0}\right)+\left(y_{0}-x_{0}\right) \subset B_{2 \Lambda^{1 / 2} d^{\tau}}\left(x_{0}\right)+\left(y_{0}-x_{0}\right) \\
& \subset B_{4 \Lambda^{1 / 2} d^{\tau}}\left(x_{0}\right) \subset E_{4(\Lambda / \lambda)^{1 / 2} d^{\tau}}\left(x_{0}\right),
\end{aligned}
$$

we have

$$
P_{x_{0}}\left(\mathfrak{a}_{x_{0}} B_{2 d^{\tau}}+y_{0}\right) \subset E_{4(\Lambda / \lambda)^{1 / 2} d^{\tau}}\left(x_{0}\right) .
$$

Thus

$$
\begin{aligned}
y \in & (1-s) P_{x_{0}}\left(\mathfrak{a}_{x_{0}} B_{2 d^{\tau}}+y_{0}\right)+s P_{y_{0}}\left(\mathfrak{a}_{x_{0}} B_{2 d^{\tau}}+y_{0}\right) \\
& \subset(1-s) E_{4(\Lambda / \lambda)^{1 / 2} d^{\tau}}\left(x_{0}\right)+s E_{2(\Lambda / \lambda)^{1 / 2} d^{\tau}}\left(y_{0}\right) \\
& \subset B_{6\left(\Lambda / \lambda^{1 / 2}\right) d^{\tau}}+x_{0}+s\left(y_{0}-x_{0}\right) \\
& \subset B_{7\left(\Lambda / \lambda^{1 / 2}\right) d^{\tau}}+x_{0} \subset E_{7(\Lambda / \lambda) d^{\tau}}\left(x_{0}\right) .
\end{aligned}
$$

Therefore,

$$
\begin{aligned}
\frac{1}{d^{\tau}} \int_{d^{\tau}}^{2 d^{\tau}} I I_{r} d r & \leq \frac{C}{d^{\tau(n+3 / 2)-\alpha}} \int_{0}^{1} \int_{E_{7(\Lambda / \lambda) d^{\tau}\left(x_{0}\right)}}|\nabla U| d z d s \\
& \leq \frac{C}{d^{\tau(n / 2+3 / 2)-\alpha}}\left(\int_{E_{7(\Lambda / \lambda) d^{\tau}\left(x_{0}\right)}}|\nabla U|^{2}\right)^{1 / 2}
\end{aligned}
$$

$$
\begin{aligned}
& \leq \frac{C}{d^{\tau(n / 2+3 / 2)-\alpha}}\left(\int_{E_{7(\Lambda / \lambda) d^{\tau}\left(x_{0}\right)}}\left|\nabla U_{x_{0}}^{*}\right|^{2}\right)^{1 / 2} \\
& \leq C\|U\|_{W^{1,2}\left(B_{1}\right)} d^{\alpha-\tau}
\end{aligned}
$$

for small $\rho$, where the third inequality follows from the quasisymmetry property and the last inequality from Lemma 2.9.3.

Step 4. By the change of variables, we have

$$
\begin{aligned}
I I I_{r} & =\int_{\partial B_{1}}\left|\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}\left(\mathfrak{a}_{x_{0}} z\right)-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\left(\mathfrak{a}_{x_{0}} z\right)\right|-\int_{\partial B_{1}}\left|\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}\left(\mathfrak{a}_{y_{0}} z\right)-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\left(\mathfrak{a}_{y_{0}} z\right)\right| \\
& \leq \int_{\partial B_{1}}\left|\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}\left(\mathfrak{a}_{x_{0}} z\right)-\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}\left(\mathfrak{a}_{y_{0}} z\right)\right|+\int_{\partial B_{1}}\left|\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\left(\mathfrak{a}_{x_{0}} z\right)-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\left(\mathfrak{a}_{y_{0}} z\right)\right| \\
& \leq C\left(\left\|\nabla\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}\right\|_{L^{\infty}\left(B_{\Lambda^{1 / 2}}\right)}+\left\|\nabla\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right\|_{L^{\infty}\left(B_{\Lambda^{1 / 2}}\right)}\right)\left|\mathfrak{a}_{x_{0}}-\mathfrak{a}_{y_{0}}\right|,
\end{aligned}
$$

where we have used the fact that both $\mathfrak{a}_{x_{0}} z$ and $\mathfrak{a}_{y_{0}} z$ are contained in $\overline{B_{\Lambda^{1 / 2}}}$ for $z \in \partial B_{1}$. To estimate the gradients of rescalings we first observe that by the inclusion $B_{r \Lambda^{1 / 2}}\left(y_{0}\right) \subset$ $E_{r(\Lambda / \lambda)^{1 / 2}}\left(y_{0}\right) \subset B_{r \Lambda / \lambda^{1 / 2}}\left(y_{0}\right)$, we have

$$
\left\|\nabla\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}\right\|_{L^{\infty}\left(B_{\Lambda^{1 / 2}}\right)} \leq \frac{C}{r^{1 / 2}}\left\|\nabla U_{y_{0}}^{*}\right\|_{L^{\infty}\left(B_{r \Lambda^{1 / 2}}\left(y_{0}\right)\right)} \leq \frac{C}{r^{1 / 2}}\|\nabla U\|_{L^{\infty}\left(B_{r \Lambda / \lambda^{1 / 2}}\left(y_{0}\right)\right)} .
$$

Let $U_{y_{0}, r}(x):=U\left(r\left(x-y_{0}\right)+y_{0}\right)$. Then, arguing as in the proof of Proposition 2.8.1, we have

$$
\left\|\nabla U_{y_{0}, r}\right\|_{L^{\infty}\left(B_{\Lambda / \lambda^{1 / 2}}\left(y_{0}\right)\right)} \leq C(n, \alpha, M)\left\|U_{y_{0}, r}\right\|_{W^{1,2}\left(B_{2 \Lambda / \lambda^{1 / 2}}\left(y_{0}\right)\right)}
$$

Thus

$$
\begin{aligned}
\|\nabla U\|_{L^{\infty}\left(B_{r \Lambda / \lambda^{1 / 2}}\left(y_{0}\right)\right)} & =\frac{1}{r}\left\|\nabla U_{y_{0}, r}\right\|_{L^{\infty}\left(B_{\Lambda / \lambda^{1 / 2}}\left(y_{0}\right)\right)} \\
& \leq \frac{C}{r}\left\|U_{y_{0}, r}\right\|_{W^{1,2}\left(B_{2 \Lambda / \lambda^{1 / 2}}\left(y_{0}\right)\right)} \\
& \leq \frac{C}{r^{n / 2+1}}\|U\|_{L^{2}\left(B_{2 r \Lambda / \lambda^{1 / 2}}\left(y_{0}\right)\right)}+\frac{C}{r^{n / 2}}\|\nabla U\|_{L^{2}\left(B_{2 r \Lambda / \lambda^{1 / 2}}\left(y_{0}\right)\right)} \\
& \leq \frac{C}{r^{n / 2+1}}\left\|U_{y_{0}}^{*}\right\|_{L^{2}\left(E_{2 r \Lambda / \lambda}\left(y_{0}\right)\right)}+\frac{C}{r^{n / 2}}\left\|\nabla U_{y_{0}}^{*}\right\|_{L^{2}\left(E_{2 r \Lambda / \lambda}\left(y_{0}\right)\right)} \\
& \leq C r^{1 / 2}\|U\|_{W^{1,2}\left(B_{1}\right)},
\end{aligned}
$$

where we have used the inclusion $B_{2 r \Lambda / \lambda^{1 / 2}}\left(y_{0}\right) \subset E_{2 r \Lambda / \lambda}\left(y_{0}\right)$ and the quasisymmetry property in the third inequality and Lemma 2.9.3 in the forth. Therefore,

$$
\left\|\nabla\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}\right\|_{L^{\infty}\left(B_{\Lambda^{1 / 2}}\right)} \leq \frac{C}{r^{1 / 2}}\|\nabla U\|_{L^{\infty}\left(B_{r \Lambda / \lambda^{1 / 2}}\left(y_{0}\right)\right)} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}
$$

Moreover, by $C_{\text {loc }}^{1}$ convergence of $\left(U_{y_{0}}^{*}\right)_{y_{0}, r}^{\phi}$ to $\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}$, we also have

$$
\begin{equation*}
\left\|\nabla\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right\|_{L^{\infty}\left(B_{\Lambda^{1 / 2}}\right)}=\lim _{r_{\mathrm{j}} \rightarrow 0+}\left\|\nabla\left(U_{y_{0}}^{*}\right)_{y_{0}, r_{j}}^{\phi}\right\|_{L^{\infty}\left(B_{\Lambda^{1 / 2}}\right)} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)} . \tag{2.10.8}
\end{equation*}
$$

Therefore,

$$
I I I_{r} \leq C\left|\mathfrak{a}_{x_{0}}-\mathfrak{a}_{y_{0}}\right|\|U\|_{W^{1,2}\left(B_{1}\right)} \leq C\|U\|_{W^{1,2}\left(B_{1}\right)} d^{\alpha}
$$

Step 5. Now we are ready to prove (2.10.4). Using the estimates in Steps 2-4 and taking the average over $d^{\tau} \leq r \leq 2 d^{\tau}$, we have

$$
\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right| \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}\left(d^{\tau \delta / 2}+d^{1-\tau}+d^{\alpha-\tau}+d^{\alpha}\right)
$$

If we simply take $\tau=\alpha / 2$, then we conclude

$$
\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right| \leq C\left|x_{0}-y_{0}\right|^{\gamma}
$$

with $\gamma=\alpha \delta / 4$ and $C=C\left(n, \alpha, M,\|U\|_{W^{1,2}\left(B_{1}\right)}\right)$.
Step 6. To prove (2.10.5), we first observe that from (2.10.4),

$$
\begin{aligned}
\int_{\partial B_{1}}\left|\left(u_{x_{0}}^{*}\right)_{0}^{\phi}(z)-\left(u_{y_{0}}^{*}\right)_{0}^{\phi}\left(\overline{\mathfrak{a}}_{y_{0}}^{-1} \overline{\mathfrak{a}}_{x_{0}} z\right)\right| & =\int_{\partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}\left(\overline{\mathfrak{a}}_{x_{0}} z\right)-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\left(\overline{\mathfrak{a}}_{x_{0}} z\right)\right| \\
& =\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}-\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right| \frac{\mu_{x_{0}}}{\operatorname{det} \mathfrak{a}_{x_{0}}} \\
& \leq C\left|x_{0}-y_{0}\right|^{\gamma} .
\end{aligned}
$$

On the other hand,

$$
\begin{aligned}
\int_{\partial B_{1}}\left|\left(u_{y_{0}}^{*}\right)_{0}^{\phi}(z)-\left(u_{y_{0}}^{*}\right)_{0}^{\phi}\left(\overline{\mathfrak{a}}_{y_{0}}^{-1} \overline{\mathfrak{a}}_{x_{0}} z\right)\right| & =\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(u_{y_{0}}^{*}\right)_{0}^{\phi}\left(\overline{\mathfrak{a}}_{x_{0}}^{-1} z\right)-\left(u_{y_{0}}^{*}\right)_{0}^{\phi}\left(\overline{\mathfrak{a}}_{y_{0}}^{-1} z\right)\right| \frac{\mu_{x_{0}}}{\operatorname{det} \mathfrak{a}_{x_{0}}} \\
& \leq C\left\|\nabla\left(u_{y_{0}}^{*}\right)_{0}^{\phi}\right\|_{L^{\infty}\left(B_{(\Lambda / \lambda)^{1 / 2}}\right)}\left|\overline{\mathfrak{a}}_{x_{0}}^{-1}-\overline{\mathfrak{a}}_{y_{0}}^{-1}\right| \\
& \leq C\left\|\nabla\left(U_{y_{0}}^{*}\right)_{y_{0}, 0}^{\phi}\right\|_{L^{\infty}\left(B_{\Lambda / \lambda^{1 / 2}}\right)}\left|x_{0}-y_{0}\right|^{\alpha} \\
& \leq C\|U\|_{W^{1,2}\left(B_{1}\right)}\left|x_{0}-y_{0}\right|^{\alpha},
\end{aligned}
$$

where the last inequality follows from (2.10.8). (It is easy to see that we can enlarge the domain in (2.10.8).) Therefore, combining the preceding two estimates, we conclude that

$$
\int_{\partial B_{1}}\left|\left(u_{x_{0}}^{*}\right)_{0}^{\phi}-\left(u_{y_{0}}^{*}\right)_{0}^{\phi}\right| \leq C\left|x_{0}-y_{0}\right|^{\gamma} .
$$

Step 7. Finally, (2.10.5) implies (2.10.6), by arguing precisely as in Proposition 7.4 in [20].

### 2.11 Regularity of the regular set

In this section we combine the uniqueness and Hölder continuous dependence of 3/2homogeneous blowups of the symmetrized almost minimizers $\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}$ (Lemma 2.10.3) with a classification of such blowups at so-called regular points (Proposition 2.11.1) to prove one of the main results of this chapter, the $C^{1, \gamma}$ regularity of the regular set (Theorem 2.11.5). While some arguments follow directly from those in the case $A \equiv I$ by a coordinate transformation $\bar{T}_{x_{0}}$, the dependence of these transformations on $x_{0}$ creates an additional difficulty.

We start by defining the regular set.

Definition 2.11.1 (Regular points). For an A-quasisymmetric almost minimizer $U$ for the A-Signorini problem in $B_{1}$, we say that a free boundary point $x_{0}$ of $U$ is regular if

$$
\kappa\left(x_{0}\right)=3 / 2 .
$$

We denote the set of all regular points of $U$ by $\mathcal{R}(U)$ and call it the regular set.

We explicitly observe here that $3 / 2<2 \leq \kappa_{0}$, so the fact $x_{0} \in \mathcal{R}(U)$ is independent of the choice of $\kappa_{0} \geq 2$, see Remark 2.7.3.

The proofs of the following two results (Lemma 2.11.1 and Proposition 2.11.1) are established precisely as in Lemma 1.9.1 and Proposition 1.9.1 for the transformed functions $u_{x_{0}}^{*}$. The equivalent statements for $U_{x_{0}}^{*}$ are obtained by changing back to the original variables.

Lemma 2.11.1 (Nondegeneracy at regular points). Let $x_{0} \in B_{1 / 2}^{\prime} \cap \mathcal{R}(U)$ for an $A$ quasisymmetric almost minimizer $U$ for the $A$-Signorini problem in $B_{1}$. Then, for $\kappa=3 / 2$,

$$
\liminf _{t \rightarrow 0} \int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left(\left(U_{x_{0}}^{*}\right)_{x_{0}, t}^{\phi}\right)^{2} \mu_{x_{0}}=\operatorname{det} \mathfrak{a}_{x_{0}} \liminf _{t \rightarrow 0} \int_{\partial B_{1}}\left(\left(u_{x_{0}}^{*}\right)_{t}^{\phi}\right)^{2}>0 .
$$

Proposition 2.11.1. If $\kappa\left(x_{0}\right)<2$, then necessarily $\kappa\left(x_{0}\right)=3 / 2$ and

$$
\begin{aligned}
\left(u_{x_{0}}^{*}\right)_{0}^{\phi}(z) & =a_{x_{0}} \operatorname{Re}\left(z^{\prime} \cdot \nu_{x_{0}}+i\left|z_{n}\right|\right)^{3 / 2} \\
\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}(x) & =a_{x_{0}} \operatorname{Re}\left(\left(\overline{\mathfrak{a}}_{x_{0}}^{-1} x\right)^{\prime} \cdot \nu_{x_{0}}+i\left|\left(\overline{\mathfrak{a}}_{x_{0}}^{-1} x\right)_{n}\right|\right)^{3 / 2}
\end{aligned}
$$

for some $a_{x_{0}}>0, \nu_{x_{0}} \in \partial B_{1}^{\prime}$.

The next two corollaries are obtained by repeating the same arguments as in Corollaries 1.9.2 and 1.9.3.

Corollary 2.11.2 (Almgren's frequency gap). Let $U$ and $x_{0}$ be as in Lemma 2.11.1. Then either

$$
\kappa\left(x_{0}\right)=3 / 2 \quad \text { or } \quad \kappa\left(x_{0}\right) \geq 2
$$

Corollary 2.11.3. The regular set $\mathcal{R}(U)$ is a relatively open subset of the free boundary.
The combination of Proposition 2.11.1 and Lemma 2.10.3 implies the following lemma.
Lemma 2.11.4. Let $U$ and $x_{0}$ be as in Lemma 2.11.1. Then there exists $\rho>0$, depending on $x_{0}$ such that $B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(U) \subset \mathcal{R}(U)$ and if

$$
\left(u_{\bar{x}}^{*}\right)_{0}^{\phi}(z)=a_{\bar{x}} \operatorname{Re}\left(z^{\prime} \cdot \nu_{\bar{x}}+i\left|z_{n}\right|\right)^{3 / 2}
$$

is the unique 3/2-homogeneous blowup of $u_{\bar{x}}^{*}$ at $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u)$, then

$$
\begin{aligned}
& \left|a_{\bar{x}}-a_{\bar{y}}\right| \leq C_{0}|\bar{x}-\bar{y}|^{\gamma}, \\
& \left|\nu_{\bar{x}}-\nu_{\bar{y}}\right| \leq C_{0}|\bar{x}-\bar{y}|^{\gamma},
\end{aligned}
$$

for any $\bar{x}, \bar{y} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(u)$ with a constant $C_{0}$ depending on $x_{0}$.
Proof. The proof follows by repeating the argument in Lemma 7.5 in [20] with $\left(u_{\bar{x}}^{*}\right)_{0}^{\phi},\left(u_{\bar{y}}^{*}\right)_{0}^{\phi}$.

Now we are ready to prove the main result on the regularity of the regular set.
Theorem 2.11.5 ( $C^{1, \gamma}$ regularity of the regular set). Let $U$ be an $A$-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$. Then, if $x_{0} \in B_{1 / 2}^{\prime} \cap \mathcal{R}(U)$, there exists $\rho>0$, depending on $x_{0}$ such that, after a possible rotation of coordinate axes in $\mathbb{R}^{n-1}$, one has $B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(U) \subset \mathcal{R}(U)$, and

$$
B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(U)=B_{\rho}^{\prime}\left(x_{0}\right) \cap\left\{x_{n-1}=g\left(x_{1}, \ldots, x_{n-2}\right)\right\},
$$

for $g \in C^{1, \gamma}\left(\mathbb{R}^{n-2}\right)$ with an exponent $\gamma=\gamma(n, \alpha, M) \in(0,1)$.
Proof. The proof of the theorem is similar to those of in Theorem 1.2 in [20] and Theorem 1.9.5. However, we provide full details since there are technical differences.

Step 1. By relative openness of $\mathcal{R}(U)$ in $\Gamma(U)$, for small $\rho>0$ we have $B_{2 \rho}^{\prime}\left(x_{0}\right) \cap \Gamma(U) \subset$ $\mathcal{R}(U)$. We then claim that for any $\varepsilon>0$, there is $r_{\varepsilon}>0$ such that for $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(U)$, $r<r_{\varepsilon}$, we have that

$$
\left\|\left(u_{\bar{x}}^{*}\right)_{r}^{\phi}-\left(u_{\bar{x}}^{*}\right)_{0}^{\phi}\right\|_{C^{1}\left(\overline{B_{1}^{ \pm}}\right)}<\varepsilon .
$$

Assuming the contrary, there is a sequence of points $\bar{x}_{\mathrm{j}} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(U)$ and radii $r_{\mathrm{j}} \rightarrow 0$ such that

$$
\left\|\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{r_{\mathrm{j}}}^{\phi}-\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{0}^{\phi}\right\|_{C^{1}\left(\overline{B_{1}^{ \pm}}\right)} \geq \varepsilon_{0}
$$

for some $\varepsilon_{0}>0$. Taking a subsequence if necessary, we may assume $\bar{x}_{\mathrm{j}} \rightarrow \bar{x}_{0} \in \overline{B_{\rho}^{\prime}\left(x_{0}\right)} \cap$ $\Gamma(U)$. Using estimates (2.10.1)-(2.10.3), we can see that $\nabla\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{r_{j}}^{\phi}$ are uniformly bounded
in $C^{0, \beta}\left(B_{2}^{ \pm} \cup B_{2}^{\prime}\right)$. Since $\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{r_{\mathrm{j}}}^{\phi}(0)=0$, we also have that $\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{r_{\mathrm{j}}}^{\phi}$ is uniformly bounded in $C^{1, \beta}\left(B_{2}^{ \pm} \cup B_{2}^{\prime}\right)$. Thus, we may assume that for some $w$

$$
\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{r_{\mathrm{j}}}^{\phi} \rightarrow w \quad \text { in } C^{1}\left(\overline{B_{1}^{ \pm}}\right) .
$$

By arguing as in the proof of Proposition 2.8.1, we see that the limit $w$ is a solution of the Signorini problem in $B_{1}$. Further, by Lemma 2.10.2, we have

$$
\left\|\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{r_{\mathrm{j}}}^{\phi}-\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{0}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)} \rightarrow 0 .
$$

On the other hand, by Lemma 2.11.4, we have

$$
\left(u_{\bar{x}_{\mathrm{x}}}^{*}\right)_{0}^{\phi} \rightarrow\left(u_{\bar{x}_{0}}^{*}\right)_{0}^{\phi} \quad \text { in } C^{1}\left(\overline{B_{1}^{ \pm}}\right),
$$

and thus

$$
w=\left(u_{\bar{x}_{0}}^{*}\right)_{0}^{\phi} \quad \text { on } \partial B_{1} .
$$

Since both $w$ and $\left(u_{\bar{x}_{0}}^{*}\right)_{0}^{\phi}$ are solutions of the Signorini problem, they must coincide also in $B_{1}$. Therefore

$$
\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{r_{\mathrm{j}}}^{\phi} \rightarrow\left(u_{\bar{x}_{0}}^{*}\right)_{0}^{\phi} \quad \text { in } C^{1}\left(\overline{B_{1}^{ \pm}}\right),
$$

implying also that

$$
\left\|\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{r_{\mathrm{j}}}^{\phi}-\left(u_{\bar{x}_{\mathrm{j}}}^{*}\right)_{0}^{\phi}\right\|_{C^{1}\left(\overline{B_{1}^{ \pm}}\right)} \rightarrow 0,
$$

which contradicts our assumption.
Step 2. For a given $\varepsilon>0$ and a unit vector $\nu \in \mathbb{R}^{n-1}$ define the cone

$$
\mathcal{C}_{\varepsilon}(\nu)=\left\{x^{\prime} \in \mathbb{R}^{n-1}: x^{\prime} \cdot \nu>\varepsilon\left|x^{\prime}\right|\right\} .
$$

By Lemma 2.11.4, we may assume $a_{\bar{x}} \geq \frac{a_{x_{0}}}{2}$ for $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(U)$ by taking $\rho$ small. For such $\rho$, we then claim that for any $\varepsilon>0$, there is $r_{\varepsilon}>0$ such that for any $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(U)$, we have

$$
\mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime} \subset\left\{u_{\bar{x}}^{*}(\cdot, 0)>0\right\} .
$$

Indeed, denoting $\mathcal{K}_{\varepsilon}(\nu)=\mathcal{C}_{\varepsilon} \cap \partial B_{1 / 2}^{\prime}$, we have for some universal $C_{\varepsilon}>0$

$$
\mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right) \Subset\left\{\left(u_{\bar{x}}^{*}\right)_{0}^{\phi}(\cdot, 0)>0\right\} \cap B_{1}^{\prime} \quad \text { and } \quad\left(u_{\bar{x}}^{*} \phi_{0}^{\phi}(\cdot, 0) \geq a_{\bar{x}} C_{\varepsilon} \geq \frac{a_{x_{0}}}{2} C_{\varepsilon} \quad \text { on } \mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right) .\right.
$$

Since $\frac{a_{x_{0}}}{2} C_{\varepsilon}$ is independent of $\bar{x}$, by Step 1 we can find $r_{\varepsilon}>0$ such that for $r<2 r_{\varepsilon}$,

$$
\left(u_{\bar{x}}^{*} \phi_{r}^{\phi}(\cdot, 0)>0 \quad \text { on } \mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right) .\right.
$$

This implies that for $r<2 r_{\varepsilon}$,

$$
u_{\bar{x}}^{*}(\cdot, 0)>0 \quad \text { on } \quad r \mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right)=\mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap \partial B_{r / 2}^{\prime} .
$$

Taking the union over all $r<2 r_{\varepsilon}$, we obtain

$$
u_{\bar{x}}^{*}(\cdot, 0)>0 \quad \text { on } \mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime} .
$$

Step 3. We claim that for given $\varepsilon>0$, there exists $r_{\varepsilon}>0$ such that for any $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(U)$, we have $-\left(\mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) \subset\left\{u_{\bar{x}}^{*}(\cdot, 0)=0\right\}$.

Indeed, we first note that

$$
-\partial_{x_{n}}^{+}\left(u_{\bar{x}}^{*}\right)_{0}^{\phi} \geq a_{\bar{x}} C_{\varepsilon}>\left(\frac{a_{x_{0}}}{2}\right) C_{\varepsilon} \quad \text { on }-\mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right),
$$

for a universal constant $C_{\varepsilon}>0$. From Step 1, there exists $r_{\varepsilon}>0$ such that for $r<2 r_{\varepsilon}$,

$$
-\partial_{x_{n}}^{+}\left(u_{\bar{x}}^{*}\right)_{r}^{\phi}(\cdot, 0)>0 \quad \text { on }-\mathcal{K}_{\varepsilon}\left(\nu_{\bar{x}}\right) .
$$

By arguing as in Step 2, we obtain

$$
-\partial_{x_{n}}^{+} u_{\bar{x}}^{*}(\cdot, 0)>0 \quad \text { on }-\left(\mathcal{C}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) .
$$

By the complementarity condition in Lemma 2.6.6, we therefore conclude that

$$
-\left(\mathcal{C}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) \subset\left\{-\partial_{x_{n}}^{+} u_{\bar{x}}^{*}(\cdot, 0)>0\right\} \subset\left\{u_{\bar{x}}^{*}(\cdot, 0)=0\right\}
$$

Step 4. By direct computation, we have

$$
\mathcal{C}_{\Lambda^{1 / 2} \lambda^{-1 / 2} \varepsilon}\left(\nu_{\bar{x}}^{A}\right) \cap B_{\lambda^{1 / 2} r_{\varepsilon}}^{\prime} \subset \overline{\mathfrak{a}}_{\bar{x}}\left(\mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime}\right),
$$

where

$$
\nu_{\bar{x}}^{A}:=\frac{\left(\overline{\mathfrak{a}}_{x}^{-1}\right)^{\operatorname{tr}} \nu_{\bar{x}}}{\left|\left(\overline{\mathfrak{a}}_{x}^{-1}\right)^{\operatorname{tr}} \nu_{\bar{x}}\right|} .
$$

(Here $(\cdot)^{\text {tr }}$ stands for the transpose of the matrix.) Indeed, if $y^{\prime} \in \mathcal{C}_{\Lambda^{1 / 2} \lambda^{-1 / 2} \varepsilon}\left(\nu_{\bar{x}}^{A}\right) \cap B_{\lambda^{1 / 2} r_{\varepsilon}}^{\prime}$, then

$$
y^{\prime} \in B_{\lambda^{1 / 2} r_{\varepsilon}}^{\prime}=\overline{\mathfrak{a}}_{\bar{x}}\left(\overline{\mathfrak{a}}_{\bar{x}}^{-1} B_{\lambda^{1 / 2} r_{\varepsilon}}^{\prime}\right) \subset \overline{\mathfrak{a}}_{\bar{x}} B_{r_{\varepsilon}}^{\prime}
$$

and

$$
\begin{aligned}
\left\langle\overline{\mathfrak{a}}_{x}^{-1} y^{\prime}, \nu_{\bar{x}}\right\rangle & =\left\langle y^{\prime},\left(\overline{\mathfrak{a}}_{x}^{-1}\right)^{\operatorname{tr}} \nu_{\bar{x}}\right\rangle=\left\langle y^{\prime}, \nu_{\bar{x}}^{A}\right\rangle\left|\left(\overline{\mathfrak{a}}_{x}^{-1}\right)^{\operatorname{tr}} \nu_{\bar{x}}\right| \\
& \geq\left(\Lambda^{1 / 2} \lambda^{-1 / 2} \varepsilon\left|y^{\prime}\right|\right)\left(\Lambda^{-1 / 2}\right) \\
& =\lambda^{-1 / 2} \varepsilon\left|y^{\prime}\right| \geq \varepsilon\left|\overline{\mathfrak{a}}_{\bar{x}}^{-1} y^{\prime}\right| .
\end{aligned}
$$

Combining this with Step 2 and Step 3, for $\bar{x} \in B_{\rho}^{\prime}\left(x_{0}\right) \cap \Gamma(U)$,

$$
\begin{aligned}
\bar{x}+\left(\mathcal{C}_{\Lambda^{1 / 2} \lambda^{-1 / 2}}\left(\nu_{\bar{x}}^{A}\right) \cap B_{\lambda^{1 / 2} r_{\varepsilon}}^{\prime}\right) & \subset \bar{x}+\overline{\mathfrak{a}}_{\bar{x}}\left(\mathcal{C}_{\varepsilon}\left(\nu_{\bar{x}}\right) \cap B_{r_{\varepsilon}}^{\prime}\right) \\
& \subset\left\{U_{\bar{x}}^{*}(\cdot, 0)>0\right\}, \\
\bar{x}-\left(\mathcal{C}_{\Lambda^{1 / 2} \lambda^{-1 / 2}}\left(\nu_{\bar{x}}^{A}\right) \cap B_{\lambda^{1 / 2} r_{\varepsilon}}^{\prime}\right) & \subset\left\{U_{\bar{x}}^{*}(\cdot, 0)=0\right\} .
\end{aligned}
$$

Step 5. By rotation in $\mathbb{R}^{n-1}$ we may assume $\nu_{x_{0}}^{A}=\mathrm{e}_{n-1}$. For any $\varepsilon>0$, by Lemma 2.11.4 and the Hölder continuity of $A$, we can take $\rho_{\varepsilon}=\rho\left(x_{0}, \varepsilon, M\right)$, possibly smaller than $\rho$ in the previous steps, such that

$$
\mathcal{C}_{2 \Lambda^{1 / 2} \lambda^{-1 / 2} \varepsilon}\left(\mathrm{e}_{n-1}\right) \cap B_{\lambda^{1 / 2} r_{\varepsilon}}^{\prime} \subset \mathcal{C}_{\Lambda^{1 / 2} \lambda^{-1 / 2} \varepsilon}\left(\nu_{\bar{x}}^{A}\right) \cap B_{\lambda^{1 / 2 r_{\varepsilon}}}^{\prime},
$$

for $\bar{x} \in B_{\rho_{\varepsilon}}^{\prime}\left(x_{0}\right) \cap \Gamma(U)$. By Step 4, we also have

$$
\begin{aligned}
& \bar{x}+\left(\mathcal{C}_{2 \Lambda^{1 / 2} \lambda^{-1 / 2} \varepsilon}\left(\mathrm{e}_{n-1}\right) \cap B_{\lambda^{1 / 2} r_{\varepsilon}}^{\prime}\right) \subset\{U(\cdot, 0)>0\}, \\
& \bar{x}-\left(\mathcal{C}_{2 \Lambda^{1 / 2} \lambda^{-1 / 2} \varepsilon}\left(\mathrm{e}_{n-1}\right) \cap B_{\lambda^{1 / 2} r_{\varepsilon}}^{\prime}\right) \subset\{U(\cdot, 0)=0\} .
\end{aligned}
$$

Now, fixing $\varepsilon=\varepsilon_{0}$, by the standard arguments, we conclude that there exists a Lipschitz function $g: \mathbb{R}^{n-2} \rightarrow \mathbb{R}$ with $|\nabla g| \leq C_{n, M} / \varepsilon_{0}$ such that

$$
\begin{aligned}
& B_{\rho_{\varepsilon_{0}}}^{\prime}\left(x_{0}\right) \cap\{U(\cdot, 0)=0\}=B_{\rho_{\varepsilon_{0}}}^{\prime}\left(x_{0}\right) \cap\left\{x_{n-1} \leq g\left(x^{\prime \prime}\right)\right\}, \\
& B_{\rho_{\varepsilon_{0}}}^{\prime}\left(x_{0}\right) \cap\{U(\cdot, 0)>0\}=B_{\rho_{\varepsilon_{0}}}^{\prime}\left(x_{0}\right) \cap\left\{x_{n-1}>g\left(x^{\prime \prime}\right)\right\}
\end{aligned}
$$

Step 6. Taking $\varepsilon \rightarrow 0$ in Step $5, \Gamma(U)$ is differentiable at $x_{0}$ with normal $\nu_{x_{0}}^{A}$. Recentering at any $\bar{x} \in B_{\rho_{e_{0}}}^{\prime}\left(x_{0}\right) \cap \Gamma(U)$, we see that $\Gamma(U)$ has a normal $\nu_{\bar{x}}^{A}$ at $\bar{x}$. By noticing that $\bar{x} \mapsto \nu_{\bar{x}}^{A}$ is $C^{0, \gamma}$, we conclude that the function $g$ in Step 5 is $C^{1, \gamma}$. This completes the proof.

### 2.12 Singular points

In this section we study another type of free boundary points for almost minimizers, the so-called singular set $\Sigma(U)$. Because of the machinery developed in the earlier sections, we are able to prove a stratification type result for $\Sigma(U)$ (Theorem 2.12.4), following a similar approach for the minimizers and almost minimizers with $A=I$.

Definition 2.12.1 (Singular points). Let $U$ be an $A$-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$. We say that a free boundary point $x_{0}$ is singular if the coincidence set $\Lambda(U)=\{U(\cdot, 0)=0\} \subset B_{1}^{\prime}$ has zero $H^{n-1}$-density at $x_{0}$, i.e.,

$$
\lim _{r \rightarrow 0+} \frac{H^{n-1}\left(\Lambda(U) \cap B_{r}^{\prime}\left(x_{0}\right)\right)}{H^{n-1}\left(B_{r}^{\prime}\right)}=0
$$

We denote the set of all singular points by $\Sigma(U)$ and call it the singular set.

Denote by $\overline{\mathfrak{a}}_{x_{0}}^{\prime}$ the $(n-1) \times(n-1)$ submatrix of $\overline{\mathfrak{a}}_{x_{0}}$ formed by the first $(n-1)$ rows and columns. We then claim that there are constants $C, c>0$ depending only on $n$, $\lambda$, and $\Lambda$ such that

$$
\begin{equation*}
c \leq\left|\operatorname{det} \overline{\mathfrak{a}}_{x_{0}}^{\prime}\right| \leq C \tag{2.12.1}
\end{equation*}
$$

Indeed, this follows from the ellipticity of $\mathfrak{a}_{x_{0}}$ and the invariance of both $\mathbb{R}^{n-1} \times\{0\}$ and $\{0\} \times \mathbb{R}$ under $\overline{\mathfrak{a}}_{x_{0}}$, since we have

$$
\left|\operatorname{det} \overline{\mathfrak{a}}_{x_{0}}^{\prime}\left(\overline{\mathfrak{a}}_{x_{0}}\right)_{n n}\right|=\left|\operatorname{det} \overline{\mathfrak{a}}_{x_{0}}\right|=\left|\operatorname{det} \mathfrak{a}_{x_{0}}\right|
$$

and

$$
\left|\left(\overline{\mathfrak{a}}_{x_{0}}\right)_{n n}\right|=\left|\left\langle\overline{\mathfrak{a}}_{x_{0}} \mathrm{e}_{n}, \mathrm{e}_{n}\right\rangle\right|=\left|\overline{\mathfrak{a}}_{x_{0}} \mathrm{e}_{n}\right| \in\left[\lambda^{1 / 2}, \Lambda^{1 / 2}\right] .
$$

Recall now that for $x_{0} \in \Gamma(u), u_{x_{0}}(y)=U\left(\overline{\mathfrak{a}}_{x_{0}} y+x_{0}\right)$ and note that $\overline{\mathfrak{a}}_{x_{0}}^{\prime} B_{r}^{\prime}+x_{0}=E_{r}^{\prime}\left(x_{0}\right)$. Thus,

$$
\begin{equation*}
H^{n-1}\left(\Lambda(U) \cap E_{r}^{\prime}\left(x_{0}\right)\right)=\left|\operatorname{det} \overline{\mathfrak{a}}_{x_{0}}^{\prime}\right| H^{n-1}\left(\Lambda\left(u_{x_{0}}^{*}\right) \cap B_{r}^{\prime}\right) \tag{2.12.2}
\end{equation*}
$$

Now, by (2.12.2) and (2.12.1), together with $B_{\lambda^{1 / 2} r}\left(x_{0}\right) \subset E_{r}\left(x_{0}\right) \subset B_{\Lambda^{1 / 2_{r}}}\left(x_{0}\right)$, we have

$$
\begin{aligned}
\lim _{r \rightarrow 0+} \frac{H^{n-1}\left(\Lambda(U) \cap B_{r}^{\prime}\left(x_{0}\right)\right)}{H^{n-1}\left(B_{r}^{\prime}\right)}=0 & \Longleftrightarrow \lim _{r \rightarrow 0+} \frac{H^{n-1}\left(\Lambda(U) \cap E_{r}^{\prime}\left(x_{0}\right)\right)}{H^{n-1}\left(E_{r}^{\prime}\left(x_{0}\right)\right)}=0 \\
& \Longleftrightarrow \lim _{r \rightarrow 0+} \frac{H^{n-1}\left(\Lambda\left(u_{x_{0}}^{*}\right) \cap B_{r}^{\prime}\right)}{H^{n-1}\left(B_{r}^{\prime}\right)}=0
\end{aligned}
$$

In terms of Almgren rescalings $\left(u_{x_{0}}^{*}\right)_{r}^{I}$, we can rewrite the condition above as

$$
\lim _{r \rightarrow 0+} H^{n-1}\left(\Lambda\left(\left(u_{x_{0}}^{*}\right)_{r}^{I}\right) \cap B_{1}^{\prime}\right)=0
$$

We then have the following characterization of singular points.

Proposition 2.12.1 (Characterization of singular points). Let $U$ be an A-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$, and $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(U)$ be such that $\kappa\left(x_{0}\right)=\kappa<\kappa_{0}$. Then the following statements are equivalent.
(i) $x_{0} \in \Sigma(U)$.
(ii) any Almgren blowup $\left(u_{x_{0}}^{*}\right)_{0}^{I}$ of $u_{x_{0}}^{*}$ at 0 is a nonzero polynomial from the class

$$
\begin{aligned}
& \mathcal{Q}_{\kappa}=\{q: q \text { is homogeneous polynomial of degree } \kappa \text { such that } \\
& \qquad \begin{array}{l}
\left.\Delta q=0, q\left(y^{\prime}, 0\right) \geq 0, q\left(y^{\prime}, y_{n}\right)=q\left(y^{\prime},-y_{n}\right)\right\}
\end{array}
\end{aligned}
$$

(iii) any Almgren blowup $\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{A}$ of $U_{x_{0}}^{*}$ at $x_{0}$ is a nonzero polynomial from the class

$$
\begin{aligned}
& \mathcal{Q}_{\kappa}^{A, x_{0}}=\{p: p \text { is homogeneous polynomial of degree } \kappa \text { such that } \\
& \left.\qquad \operatorname{div}\left(A\left(x_{0}\right) \nabla p\right)=0, p\left(x^{\prime}, 0\right) \geq 0, p(x)=p\left(P_{x_{0}} x\right)\right\}
\end{aligned}
$$

(iv) $\kappa\left(x_{0}\right)=2 m$ for some $m \in \mathbb{N}$.

Proof. This is the analogue of Proposition 1.10.1 in the case $A \equiv I$.
Clearly, (ii) and (iii) are equivalent. By Proposition 2.8.1, any Almgren blowup $\left(u_{x_{0}}^{*}\right)_{0}^{I}$ of $u_{x_{0}}^{*}$ at 0 is a nonzero global solution of the Signorini problem, homogeneous of degree $\kappa$. Moreover, $\left(u_{x_{0}}^{*}\right)_{0}^{I}$ is a $C_{\text {loc }}^{1}$ limit of Almgren rescalings $\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{I}$ in $\mathbb{R}_{ \pm}^{n} \cup \mathbb{R}^{n-1}$. Since $u_{x_{0}}^{*}$ also satisfies the complementarity condition in Lemma 2.6.6, the equivalence among (i), (ii) and (iv) follows by repeating the arguments in Proposition 2.8.1.

In order to proceed with the blowup analysis at singular points, we need to remove the logarithmic term from the growth estimates in Lemma 2.9.1. This was achieved in

Lemma 1.10.6 in the case $A \equiv I$ by using a bootstrapping argument Lemmas 1.10.2-1.10.4, Corollary 1.10.5, based on the log-epiperimetric inequality of [27]. All the arguments above work directly for $u_{x_{0}}^{*}$ (and then for $U_{x_{0}}^{*}$, by deskewing) and we obtain the following optimal growth estimate.

Lemma 2.12.1 (Optimal growth estimate at singular points). Let $U$ be an A-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$. If $x_{0} \in B_{1 / 2}^{\prime} \cap \Gamma(U)$ and $\kappa\left(x_{0}\right)=\kappa<\kappa_{0}, \kappa=2 m, m \in \mathbb{N}$, then there are $t_{0}$ and $C$, depending on $n, \alpha, M, \kappa, \kappa_{0}$, $\|U\|_{W^{1,2}\left(B_{1}\right)}$, such that for $0<t<t_{0}$,

$$
\begin{aligned}
\int_{\partial B_{t}}\left(u_{x_{0}}^{*}\right)^{2} & \leq C t^{n+2 \kappa-1}, \quad \int_{B_{t}}\left|\nabla u_{x_{0}}^{*}\right|^{2} \leq C t^{n+2 \kappa-2}, \\
\int_{\partial E_{t}\left(x_{0}\right)}\left(U_{x_{0}}^{*}\right)^{2} & \leq C t^{n+2 \kappa-1}, \quad \int_{E_{t}\left(x_{0}\right)}\left|\nabla U_{x_{0}}^{*}\right|^{2} \leq C t^{n+2 \kappa-2}
\end{aligned}
$$

With this growth estimate at hand, we now proceed as in the beginning of Section 2.10 but with $\kappa=2 m<\kappa_{0}$ in place of $\kappa=3 / 2$. Namely, for such $\kappa$, let

$$
\phi(r)=\phi_{\kappa}(r):=\mathrm{e}^{-\left(\frac{\kappa b}{\alpha}\right) r^{\alpha}} r^{\kappa}, \quad 0<r<t_{0},
$$

where $b=\frac{M\left(n+2 \kappa_{0}\right)}{\alpha}$ is as in Weiss-type monotonicity formula. Then, define the $\kappa$-almost homogeneous rescalings of a function $v$ at $x_{0}$ by

$$
v_{x_{0}, r}^{\phi}(x):=\frac{v\left(r x+x_{0}\right)}{\phi(r)} .
$$

Again, when $x_{0}=0$, we simply write $v_{0, r}^{\phi}=v_{r}^{\phi}$.
The growth estimates in Lemma 2.12.1 enable us to consider $\kappa$-homogeneous blowups

$$
\begin{aligned}
\left(u_{x_{0}}^{*}\right)_{t_{\mathrm{j}}}^{\phi} \rightarrow\left(u_{x_{0}}^{*}\right)_{0}^{\phi} & \text { in } C_{\mathrm{loc}}^{1}\left(\mathbb{R}_{ \pm}^{n} \cup \mathbb{R}^{n-1}\right), \\
\left(U_{x_{0}}^{*}\right)_{x_{0}, t_{\mathrm{j}}}^{\phi} \rightarrow\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi} & \text { in } C_{\mathrm{loc}}^{1}\left(\mathbb{R}_{ \pm}^{n} \cup \mathbb{R}^{n-1}\right),
\end{aligned}
$$

for $t=t_{\mathrm{j}} \rightarrow 0+$, similar to 3/2-homogeneous blowups in Section 2.10.

Furthermore, the arguments in Proposition 1.10.2 also go through for $u_{x_{0}}^{*}$ (and then for $U_{x_{0}}^{*}$, by deskewing), and we obtain the following rotation estimate for almost homogeneous rescalings.

Proposition 2.12.2 (Rotation estimate). For $U$ and $x_{0}$ as in Lemma 2.12.1, there exist $C>0$ and $t_{0}>0$ such that

$$
\begin{array}{r}
\int_{\partial B_{1}}\left|\left(u_{x_{0}}^{*}\right)_{t}^{\phi}-\left(u_{x_{0}}^{*}\right)_{s}^{\phi}\right| \leq C\left(\log \frac{1}{t}\right)^{-\frac{1}{n-2}}, \\
\int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left|\left(U_{x_{0}}^{*}\right)_{x_{0}, t}^{\phi}-\left(U_{x_{0}}^{*}\right)_{x_{0}, s}^{\phi}\right| \leq C\left(\log \frac{1}{t}\right)^{-\frac{1}{n-2}},
\end{array}
$$

for $0<s<t<t_{0}$. In particular, the blowups $\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ and $\left(U_{x_{0}}^{*}\right)_{x_{0}, 0}^{\phi}$ are unique.
We next show that the rotation estimate as above holds uniformly for $u_{x_{0}}^{*}$ replaced with its Almgren rescalings $\left(u_{x_{0}}^{*}\right)_{r}^{I}, 0<r<1$. (Note that the objects $\left[\left(u_{x_{0}}^{*}\right)_{r}^{I}\right]_{t}^{\phi}$ in the proposition below are $\kappa$-almost homogeneous rescalings of Almgren rescalings.)

Proposition 2.12.3. For $U$ and $x_{0}$ as in Lemma 2.12.1 and $0<r<1$, there are $C>0$ and $t_{0}>0$, independent of $r$ such that

$$
\int_{\partial B_{1}}\left|\left[\left(u_{x_{0}}^{*}\right)_{r}^{I}\right]_{t}^{\phi}-\left[\left(u_{x_{0}}^{*}\right)_{r}^{I}\right]_{s}^{\phi}\right| \leq C\left(\log \frac{1}{t}\right)^{-\frac{1}{n-2}}
$$

for $0<s<t<t_{0}$. In particular, the $\kappa$-homogeneous blowup $\left[\left(u_{x_{0}}^{*}\right)_{r}^{I}\right]_{0}^{\phi}$ is unique.
Proof. We first observe that since $u_{x_{0}}^{*}$ has the almost Signorini property at $0,\left(u_{x_{0}}^{*}\right)_{r}^{I}$ also has the almost Signorini property at 0 . This implies that $W_{\kappa}\left(\rho,\left(u_{x_{0}}^{*}\right)_{r}^{I}, 0\right)$ and $\widehat{N}_{\kappa_{0}}\left(\rho,\left(u_{x_{0}}^{*}\right)_{r}^{I}, 0\right)$ are monotone nondecreasing on $\rho$. Thus

$$
\widehat{N}_{\kappa_{0}}\left(0+,\left(u_{x_{0}}^{*}\right)_{r}^{I}, 0\right)=\lim _{\rho \rightarrow 0} \widehat{N}_{\kappa_{0}}\left(\rho,\left(u_{x_{0}}^{*}\right)_{r}^{I}, 0\right)=\lim _{\rho \rightarrow 0} \widehat{N}_{\kappa_{0}}\left(\rho r, u_{x_{0}}^{*}, 0\right)=\kappa\left(x_{0}\right)=\kappa .
$$

Fix $R>1$. If $t$ is small, then we can argue as in the proof of Proposition 2.8.1 to obtain that for any $K \Subset B_{R}^{ \pm} \cup B_{R}^{\prime}$,

$$
\left\|\left[\left(u_{x_{0}}^{*}\right)_{r}^{I}\right]_{t}^{\phi}\right\|_{C^{1, \beta}(K)} \leq C(n, \alpha, M, R, K)\left\|\left[\left(u_{x_{0}}^{*}\right)_{r}^{I}\right]_{t}^{\phi}\right\|_{W^{1,2}\left(B_{R}\right)}
$$

Those are all we need to proceed all the arguments with $\left(u_{x_{0}}^{*}\right)_{r}^{I}$ as in Lemmas 1.10.2-1.10.4, Corollary 1.10.5, Lemma 1.10.6, and Proposition 1.10.2. This completes the proof.

Once we have Proposition 2.12.3, we can argue as in Lemma 1.10.8 to obtain the nondegeneracy for $u_{x_{0}}^{*}$, and also for $U_{x_{0}}^{*}$.

Lemma 2.12.2 (Nondegeneracy at singular points). Let $U$ and $x_{0}$ be as in Lemma 2.12.1. Then

$$
\begin{array}{r}
\liminf _{t \rightarrow 0} \int_{\partial B_{1}}\left(\left(u_{x_{0}}^{*}\right)_{t}^{\phi}\right)^{2}=\liminf _{t \rightarrow 0} \frac{1}{t^{n+2 \kappa-1}} \int_{\partial B_{t}}\left(u_{x_{0}}^{*}\right)^{2}>0, \\
\liminf _{t \rightarrow 0} \int_{\mathfrak{a}_{x_{0}} \partial B_{1}}\left(\left(U_{x_{0}}^{*}\right)_{x_{0}, t}^{\phi}\right)^{2}=\liminf _{t \rightarrow 0} \frac{1}{t^{n+2 \kappa-1}} \int_{\partial E_{t}\left(x_{0}\right)}\left(U_{x_{0}}^{*}\right)^{2}>0 .
\end{array}
$$

To state our main result on the singular set, we need to introduce certain subsets of $\Sigma(U)$. For $\kappa=2 m<\kappa_{0}, m \in \mathbb{N}$, let

$$
\Sigma_{\kappa}(U):=\left\{x_{0} \in \Sigma(U): \kappa\left(x_{0}\right)=\kappa\right\}=\Gamma_{\kappa}(U) .
$$

Note that the last equality follows from the implication (iv) $\Rightarrow$ (i) in Proposition 2.12.1.

Lemma 2.12.3. The set $\Sigma_{\kappa}(U)$ is of topological type $F_{\sigma}$; i.e., it is a countable union of closed sets.

Proof. For $\mathrm{j} \in \mathbb{N}, \mathrm{j} \geq 2$, let

$$
F_{\mathrm{j}}:=\left\{x_{0} \in \Sigma_{\kappa}(U) \cap \overline{B_{1-1 / \mathrm{j}}}: \frac{1}{\mathrm{j}} \leq \frac{1}{\rho^{n+2 \kappa-1}} \int_{\partial E_{\rho}\left(x_{0}\right)}\left(U_{x_{0}}^{*}\right)^{2} \leq \mathrm{j} \text { for } 0<\rho<\frac{1}{2 \mathrm{j}}\right\} .
$$

Note that if $x_{\mathrm{j}} \rightarrow x_{0}$, then by the local uniform continuity of $U$ and $A$,

$$
\int_{\partial E_{\rho}\left(x_{\mathrm{i}}\right)}\left(U_{x_{\mathrm{i}}}^{*}\right)^{2} \rightarrow \int_{\partial E_{\rho}\left(x_{0}\right)}\left(U_{x_{0}}^{*}\right)^{2} .
$$

Using this, together with Lemma 2.12.1, Lemma 2.12.2 and Lemma 2.9.1, we can argue as in Lemma 1.10 .9 to prove that $\Sigma_{\kappa}(U)=\cup_{\mathrm{j}=2}^{\infty} F_{\mathrm{j}}$ and each $F_{\mathrm{j}}$ is closed.

Next, for $\kappa=2 m<\kappa_{0}, m \in \mathbb{N}$ and $x_{0} \in \Sigma_{\kappa}(U)$, we define

$$
d_{x_{0}}^{(\kappa)}:=\operatorname{dim}\left\{\xi \in \mathbb{R}^{n-1}: \xi \cdot \nabla_{y^{\prime}}\left(u_{x_{0}}^{*}\right)_{0}^{\phi}\left(y^{\prime}, 0\right) \equiv 0 \text { on } \mathbb{R}^{n-1}\right\},
$$

which has the meaning of the dimension of $\Sigma_{\kappa}\left(u_{x_{0}}^{*}\right)$ at 0 , and where $\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ is the unique $\kappa$ homogeneous blowup of $u_{x_{0}}^{*}$ at 0 . We note here that $d_{x_{0}}^{(\kappa)}$ can only take the values $0,1, \ldots, n-$ 2. Indeed, otherwise $\left(u_{x_{0}}^{*}\right)_{0}^{\phi}$ would vanish identically on $\Pi$ and consequently on $\mathbb{R}^{n}$, since it is a solution of the Signorini problem, even symmetric with respect to $\Pi$ (see [14]). However, that would contradict the nondegeneracy Lemma 2.12.2. Then, for $d=0,1, \ldots, n-2$, let

$$
\Sigma_{\kappa}^{d}(U):=\left\{x_{0} \in \Sigma_{\kappa}(U): d_{x_{0}}^{(\kappa)}=d\right\} .
$$

Theorem 2.12.4 (Structure of the singular set). Let $U$ be an A-quasisymmetric almost minimizer for the $A$-Signorini problem in $B_{1}$. Then for every $\kappa=2 m<\kappa_{0}, m \in \mathbb{N}$, and $d=0,1, \ldots, n-2$, the set $\Sigma_{\kappa}^{d}(U)$ is contained in the union of countably many submanifolds of dimension $d$ and class $C^{1, \log }$.

Proof. We follow the idea in Theorem 1.10.10. For $x_{0} \in \Sigma_{\kappa}(U) \cap B_{1 / 2}^{\prime}$, let $q_{x_{0}} \in \mathcal{Q}_{\kappa}$ denote the unique $\kappa$-homogeneous blowup of $u_{x_{0}}^{*}$ at 0 . By the optimal growth (Lemma 2.12.1) and the nondegeneracy (Lemma 2.12.2), we can write

$$
q_{x_{0}}=\eta_{x_{0}} q_{x_{0}}^{I}, \quad \eta_{x_{0}}>0, \quad\left\|q_{x_{0}}^{I}\right\|_{L^{2}\left(\partial B_{1}\right)}=1
$$

where $q_{x_{0}}^{I} \in \mathcal{Q}_{\kappa}$ is the corresponding Almgren blowup. If $x_{1}, x_{2} \in \Sigma_{\kappa}(U) \cap B_{1 / 2}^{\prime}$, for $t>0$, to be chosen below, we can write

$$
\begin{align*}
\left\|q_{x_{1}}-q_{x_{2}}\right\|_{L^{1}\left(\partial B_{1}\right)} & \leq\left\|q_{x_{1}}-\left(u_{x_{1}}^{*}\right)_{t}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)}+\left\|\left(u_{x_{1}}^{*}\right)_{t}^{\phi}-\left(u_{x_{2}}^{*}\right)_{t}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)}+\left\|q_{x_{2}}-\left(u_{x_{2}}^{*}\right)_{t}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)} \\
& \leq C\left(\log \frac{1}{t}\right)^{-\frac{1}{n-2}}+\left\|\left(u_{x_{1}}^{*}\right)_{t}^{\phi}-\left(u_{x_{2}}^{*}\right)_{t}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)} \tag{2.12.3}
\end{align*}
$$

where we have used Proposition 2.12.2 in the second inequality. Moreover, we have

$$
\begin{align*}
\left\|\left(u_{x_{1}}^{*}\right)_{t}^{\phi}-\left(u_{x_{2}}^{*}\right)_{t}^{\phi}\right\|_{L^{1}\left(\partial B_{1}\right)}= & \left.\frac{1}{2 \phi(t)} \int_{\partial B_{1}} \right\rvert\, U\left(t \overline{\mathfrak{a}}_{x_{1}} y+x_{1}\right)+U\left(P_{x_{1}}\left(t \overline{\mathfrak{a}}_{x_{1}} y+x_{1}\right)\right) \\
& \quad-U\left(t \overline{\mathfrak{a}}_{x_{2}} y+x_{2}\right)-U\left(P_{x_{2}}\left(t \overline{\mathfrak{a}}_{x_{2}} y+x_{2}\right)\right) \mid d S_{y} \\
\leq & \frac{C}{t^{\kappa}} \int_{\partial B_{1}}\left(\left|U\left(t \overline{\mathfrak{a}}_{x_{1}} y+x_{1}\right)-U\left(t \overline{\mathfrak{a}}_{x_{2}} y+x_{2}\right)\right|\right. \\
& \quad+\left|U\left(P_{x_{1}}\left(t \overline{\mathfrak{a}}_{x_{1}} y+x_{1}\right)\right)-U\left(P_{x_{1}}\left(t \overline{\mathfrak{a}}_{x_{2}} y+x_{2}\right)\right)\right| \\
& \left.\quad+\left|U\left(P_{x_{1}}\left(t \overline{\mathfrak{a}}_{x_{2}} y+x_{2}\right)\right)-U\left(P_{x_{2}}\left(t \overline{\mathfrak{a}}_{x_{2}} y+x_{2}\right)\right)\right|\right) d S_{y} \\
\leq & \frac{C}{t^{\kappa}}\|\nabla U\|_{L^{\infty}\left(B_{1}\right)}\left(\left|\overline{\mathfrak{a}}_{x_{1}}-\overline{\mathfrak{a}}_{x_{2}}\right|+\left|x_{1}-x_{2}\right|+\left|P_{x_{1}}-P_{x_{2}}\right|\right) \\
\leq & C \frac{\left|x_{1}-x_{2}\right|^{\alpha}}{t^{\kappa}}=C\left|x_{1}-x_{2}\right|^{\alpha / 2}, \tag{2.12.4}
\end{align*}
$$

if we choose $t=\left|x_{1}-x_{2}\right|^{\frac{\alpha}{2 \kappa}}$ and have $\left|x_{1}-x_{2}\right|<\left(1 / 4 \Lambda^{-1} \lambda^{1 / 2}\right)^{\frac{2 \kappa}{\alpha}}$. Combining (2.12.3) and (2.12.4), we obtain

$$
\left\|q_{x_{1}}-q_{x_{2}}\right\|_{L^{1}\left(\partial B_{1}\right)} \leq C\left(\log \frac{1}{\left|x_{1}-x_{2}\right|}\right)^{-\frac{1}{n-2}}
$$

After this, we can repeat the argument in the proof of Theorem 1.10.10 to obtain the estimates that for $x_{0} \in \Sigma_{\kappa}(U) \cap B_{1 / 2}^{\prime}$, there is $\delta=\delta\left(x_{0}\right)>0$ such that

$$
\begin{gathered}
\left|\eta_{x_{1}}-\eta_{x_{2}}\right| \leq C\left(\log \frac{1}{\left|x_{1}-x_{2}\right|}\right)^{-\frac{1}{2(n-2)}} \\
\left\|q_{x_{1}}^{I}-q_{x_{2}}^{I}\right\|_{L^{\infty}\left(B_{1}\right)} \leq C\left(\log \frac{1}{\left|x_{1}-x_{2}\right|}\right)^{-\frac{1}{2(n-2)}}, \quad x_{1}, x_{2} \in \Sigma_{\kappa}(U) \cap B_{\delta}^{\prime}\left(x_{0}\right) .
\end{gathered}
$$

Now, we also have the similar result for $U_{x_{0}}^{*}$. For $x_{0} \in \Sigma_{\kappa}(U) \cap B_{1 / 2}$, where $\kappa=2 m, m \in \mathbb{N}$, let $p_{x_{0}} \in \mathcal{Q}_{\kappa}^{A, x_{0}}$ be the unique $\kappa$-homogeneous blowup of $U_{x_{0}}^{*}$ at $x_{0}$. Then we can write

$$
p_{x_{0}}=\eta_{x_{0}}^{A} p_{x_{0}}^{A}, \quad \eta_{x_{0}}^{A}>0, \quad\left\|p_{x_{0}}^{A}\right\|_{L^{2}\left(\partial B_{1}\right)}=1
$$

where $p_{x_{0}}^{A} \in \mathcal{Q}_{\kappa}^{A, x_{0}}$ is the corresponding Almgren blowup of $U_{x_{0}}^{*}$. Using that

$$
q_{x_{0}}^{I}(z)=\left(\operatorname{det} \mathfrak{a}_{x_{0}}\right)^{1 / 2} p_{x_{0}}^{A}\left(\mathfrak{a}_{x_{0}} z\right), \quad q_{x_{0}}(z)=p_{x_{0}}\left(\mathfrak{a}_{x_{0}} z\right),
$$

together with the ellipticity and Hölder continuity of $\mathfrak{a}_{x_{0}}$ and the homogeneity of blowups, we easily conclude that for $x_{0} \in \Sigma_{\kappa}(U) \cap B_{1 / 2}^{\prime}$, there is $\delta=\delta\left(x_{0}\right)>0$ such that

$$
\begin{aligned}
\left|\eta_{x_{1}}^{A}-\eta_{x_{2}}^{A}\right| \leq C\left(\log \frac{1}{\left|x_{1}-x_{2}\right|}\right)^{-\frac{1}{2(n-2)}} \\
\left\|p_{x_{1}}^{A}-p_{x_{2}}^{A}\right\|_{L^{\infty}\left(B_{1}\right)} \leq C\left(\log \frac{1}{\left|x_{1}-x_{2}\right|}\right)^{-\frac{1}{2(n-2)}}, \quad x_{1}, x_{2} \in \Sigma_{\kappa}(U) \cap B_{\delta}^{\prime}\left(x_{0}\right) .
\end{aligned}
$$

Once we have these estimates, as well as Lemma 2.12.3, we can apply the Whitney Extension Theorem of Fefferman [56], to complete the proof, similar to that of Theorem 1.7 in [27].

## 2.A Example of almost minimizers

Example 2.A.1. Let $U$ be a solution of the $A$-Signorini problem in $B_{1}$ with velocity field $b \in L^{p}\left(B_{1}\right), p>n:$

$$
\begin{aligned}
& -\operatorname{div}(A \nabla U)+\langle b(x), \nabla U\rangle=0 \quad \text { in } B_{1}^{ \pm}, \\
& U \geq 0, \quad\left\langle A \nabla U, \nu^{+}\right\rangle+\left\langle A \nabla U, \nu^{-}\right\rangle \geq 0 \\
& U\left(\left\langle A \nabla U, \nu^{+}\right\rangle+\left\langle A \nabla U, \nu^{-}\right\rangle\right)=0 \quad \text { on } B_{1}^{\prime},
\end{aligned}
$$

where $\nu^{ \pm}=\mp \mathrm{e}_{n}$ and $\left\langle A \nabla U, \nu^{ \pm}\right\rangle$on $B_{1}^{\prime}$ are understood as the limits from inside $B_{1}^{ \pm}$. We interpret this in the weak sense that $U$ satisfies the variational inequality

$$
\int_{B_{1}}\langle A \nabla U, \nabla(W-U)\rangle+\langle b, \nabla U\rangle(W-U) \geq 0
$$

for any competitor $W \in \mathfrak{K}_{0, U}\left(B_{1}, \Pi\right)$. Then $U$ is an almost minimizer of the $A$-Signorini problem in $B_{1}$ with thin obstacle $\psi=0$ on $\Pi=\mathbb{R}^{n-1} \times\{0\}$ and a gauge function $\omega(r)=$ $C r^{1-n / p}, C=C(n, p, \lambda, \Lambda)\|b\|_{L^{p}\left(B_{1}\right)}^{2}$.

Proof. For any $E_{r}\left(x_{0}\right) \Subset B_{1}$ and $W \in \mathfrak{K}_{0, U}\left(E_{r}\left(x_{0}\right), \Pi\right)$, we extend $W$ as equal to $U$ in $B_{1} \backslash E_{r}\left(x_{0}\right)$ to obtain

$$
\begin{equation*}
\int_{E_{r}\left(x_{0}\right)}\langle A \nabla U, \nabla(W-U)\rangle+\langle b, \nabla U\rangle(W-U) \geq 0 . \tag{2.A.1}
\end{equation*}
$$

Let $V$ be the minimizer of the energy functional

$$
\int_{E_{r}\left(x_{0}\right)}\langle A \nabla V, \nabla V\rangle \quad \text { on } \mathfrak{K}_{0, U}\left(E_{r}\left(x_{0}\right), \Pi\right)
$$

Then it follows from a standard variation argument that $V$ satisfies the variational inequality

$$
\begin{equation*}
\int_{E_{r}\left(x_{0}\right)}\langle A \nabla V, \nabla(W-V)\rangle \geq 0 \quad \text { for any } W \in \mathfrak{K}_{0, U}\left(E_{r}\left(x_{0}\right), \Pi\right) \tag{2.A.2}
\end{equation*}
$$

Taking $W=U \pm(U-V)^{+}$in (2.A.1) and $W=V+(U-V)^{+}$in (2.A.2), we obtain

$$
\int_{E_{r}\left(x_{0}\right)}\left\langle A \nabla(U-V)^{+}, \nabla(U-V)^{+}\right\rangle \leq-\int_{E_{r}\left(x_{0}\right)}\langle b, \nabla U\rangle(U-V)^{+} .
$$

Similarly, taking $W=U+(V-U)^{+}$in (2.A.1) and $W=V \pm(V-U)^{+}$in (2.A.2), we get

$$
\int_{E_{r}\left(x_{0}\right)}\left\langle A \nabla(V-U)^{+}, \nabla(V-U)^{+}\right\rangle \leq \int_{E_{r}\left(x_{0}\right)}\langle b, \nabla U\rangle(V-U)^{+} .
$$

These two inequalities give

$$
\int_{E_{r}\left(x_{0}\right)}\langle A \nabla(U-V), \nabla(U-V)\rangle \leq \int_{E_{r}\left(x_{0}\right)}|b||\nabla U \| U-V| .
$$

Applying Hölder's inequality,

$$
\begin{aligned}
\int_{E_{r}\left(x_{0}\right)}|\nabla(U-V)|^{2} & \leq \lambda^{-1} \int_{E_{r}\left(x_{0}\right)}\langle A \nabla(U-V), \nabla(U-V)\rangle \\
& \leq \lambda^{-1}\|b\|_{L^{p}\left(E_{r}\left(x_{0}\right)\right)}\|\nabla U\|_{L^{2}\left(E_{r}\left(x_{0}\right)\right)}\|U-V\|_{L^{p^{*}}\left(E_{r}\left(x_{0}\right)\right)}
\end{aligned}
$$

with $p^{*}=2 p /(p-2)$. Since $U-V \in W_{0}^{1,2}\left(E_{r}\left(x_{0}\right)\right)$ and $\operatorname{diam}\left(E_{r}\left(x_{0}\right)\right) \leq 2 \Lambda^{1 / 2} r$, from the Sobolev's inequality,

$$
\|U-V\|_{L^{p^{*}}\left(E_{r}\left(x_{0}\right)\right)} \leq C(n, p, \lambda, \Lambda) r^{1-n / p}\|\nabla(U-V)\|_{L^{2}\left(E_{r}\left(x_{0}\right)\right)} .
$$

Now we have

$$
\begin{equation*}
\int_{E_{r}\left(x_{0}\right)}|\nabla(U-V)|^{2} \leq C r^{2(1-n / p)} \int_{E_{r}\left(x_{0}\right)}|\nabla U|^{2} \tag{2.A.3}
\end{equation*}
$$

with $C=C(n, p, \lambda, \Lambda)\|b\|_{L^{p}\left(B_{1}\right)}^{2}$. Thus,

$$
\begin{aligned}
\int_{E_{r}\left(x_{0}\right)}\langle A \nabla U, \nabla U\rangle-\int_{E_{r}\left(x_{0}\right)}\langle & A \nabla V, \nabla V\rangle=\int_{E_{r}\left(x_{0}\right)}\langle A \nabla(U+V), \nabla(U-V)\rangle \\
& \leq C \int_{E_{r}\left(x_{0}\right)}|\nabla(U+V)||\nabla(U-V)| \\
& \leq C r^{\gamma} \int_{E_{r}\left(x_{0}\right)}\left(|\nabla U|^{2}+|\nabla V|^{2}\right)+C r^{-\gamma} \int_{E_{r}\left(x_{0}\right)}|\nabla(U-V)|^{2} \\
& \leq C r^{\gamma} \int_{E_{r}\left(x_{0}\right)}\langle A \nabla U, \nabla U\rangle+C r^{\gamma} \int_{E_{r}\left(x_{0}\right)}\langle A \nabla V, \nabla V\rangle \\
& +C r^{2(1-n / p)-\gamma} \int_{E_{r}\left(x_{0}\right)}\langle A \nabla U, \nabla U\rangle
\end{aligned}
$$

where we applied Young's inequality and used (2.A.3) at the end. We choose $\gamma=1-n / p$ to complete the proof.

## 3. ALMOST MINIMIZERS FOR CERTAIN FRACTIONAL VARIATIONAL PROBLEMS

### 3.1 Introduction and Main Results

### 3.1.1 Fractional harmonic functions

Given $0<s<1$, we say that a function $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right):=L^{1}\left(\mathbb{R}^{n},\left(1+|x|^{n+2 s}\right)^{-1}\right)$ is $s$-fractional harmonic in an open set $\Omega \subset \mathbb{R}^{n}$ if

$$
\begin{equation*}
\left(-\Delta_{x}\right)^{s} u(x):=C_{n, s} \text { p.v. } \int_{\mathbb{R}^{n}} \frac{u(x)-u(x+z)}{|z|^{n+2 s}}=0 \quad \text { in } \Omega, \tag{3.1.1}
\end{equation*}
$$

where p.v. stands for Cauchy's principal value and $C_{n, s}$ is a normalization constant. The formula above is just one of many equivalent definitions of the fractional Laplacian $\left(-\Delta_{x}\right)^{s}$, another one being a pseudo-differential operator with Fourier symbol $|\xi|^{2 s}$. We refer to a recent review of Garofalo [60] for basic properties of $\left(-\Delta_{x}\right)^{s}$, as well as many historical remarks concerning that operator.

In recent years, there has been a surge of interest in nonlocal problems involving the fractional Laplacian, when it was discovered that the problems can be localized by the use of the so-called Caffarelli-Silvestre extension procedure [10]. Namely, for $a=1-2 s \in(-1,1)$, let

$$
P(x, y):=C_{n, a} \frac{|y|^{1-a}}{\left(|x|^{2}+|y|^{2}\right)^{\frac{n+1-a}{2}}}, \quad(x, y) \in \mathbb{R}^{n} \times \mathbb{R}_{+}=\mathbb{R}_{+}^{n+1}
$$

(to be called the Poisson kernel for the extension operator $L_{a}$ ) and consider the convolution, still denoted by $u$,

$$
u(x, y):=u * P(\cdot, y)=\int_{\mathbb{R}^{n}} u(z) P(x-z, y) d z, \quad(x, y) \in \mathbb{R}_{+}^{n+1}
$$

Note that $u(x, y)$ solves the Cauchy problem

$$
\begin{aligned}
L_{a} u:=\operatorname{div}\left(|y|^{a} \nabla u\right)=0 & \text { in } \mathbb{R}_{+}^{n+1}, \\
u(x, 0)=u(x) & \text { on } \mathbb{R}^{n},
\end{aligned}
$$

where $\nabla=\nabla_{x, y}$ is the full gradient in $x$ and $y$ variables. $L_{a}$ is known as the Caffarelli-Silvestre extension operator. Then, one can recover $\left(-\Delta_{x}\right)^{s} u$ as the fractional normal derivative on $\mathbb{R}^{n}$

$$
\left(-\Delta_{x}\right)^{s} u(x)=-C_{n, a} \lim _{y \rightarrow 0+} y^{a} \partial_{y} u(x, y), \quad x \in \mathbb{R}^{n}
$$

to be understood in the appropriate sense of traces. Now, going back to the definition (3.1.1), if we consider the even reflection of $u$ in $y$-variable to all of $\mathbb{R}^{n+1}$, i.e.,

$$
u(x, y)=u(x,-y), \quad x \in \mathbb{R}^{n}, y<0
$$

then the following fact holds: $u(x)$ is $s$-fractional harmonic in $\Omega$ if and only if $u(x, y)$ satisfies

$$
\begin{equation*}
L_{a} u=0 \quad \text { in } \widetilde{\Omega}:=\mathbb{R}_{-}^{n+1} \cup(\Omega \times\{0\}) \cup \mathbb{R}_{+}^{n+1} \tag{3.1.2}
\end{equation*}
$$

(We will refer to solutions of $L_{a} u=0$ as $L_{a}$-harmonic functions.) This is essentially Lemma 4.1 in [10]. Since $L_{a} u=0$ in $\mathbb{R}_{ \pm}^{n}$ by definition, the condition (3.1.2) is equivalent to asking

$$
L_{a} u=0 \quad \text { in } \mathbb{B}_{r}\left(x_{0}\right),
$$

for any ball $\mathbb{B}_{r}\left(x_{0}\right)$ centered at $x_{0} \in \Omega$ such that $\mathbb{B}_{r}\left(x_{0}\right) \Subset \widetilde{\Omega}$, or equivalently $B_{r}\left(x_{0}\right) \Subset \Omega$. Now, observing that the solutions of the above equation are minimizers of the weighted Dirichlet energy $\int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla v|^{2}|y|^{a}$, we obtain the following fact.

Proposition 3.1.1. A function $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$ is s-fractional harmonic in $\Omega$ if and only if its reflected Caffarelli-Silvestre extension $u(x, y)$ is in $W_{\text {loc }}^{1,2}\left(\widetilde{\Omega},|y|^{a}\right)$ and for any ball $\mathbb{B}_{r}\left(x_{0}\right)$ with $x_{0} \in \Omega$ such that $B_{r}\left(x_{0}\right) \Subset \Omega$, we have

$$
\int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla u|^{2}|y|^{a} \leq \int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla v|^{2}|y|^{a}
$$

for any $v \in u+W_{0}^{1,2}\left(\mathbb{B}_{r}\left(x_{0}\right),|y|^{a}\right)$.

We take this proposition as the starting point for the definition of almost $s$-fractional harmonic functions, in the spirit of Anzellotti [31].

Definition 3.1.1 (Almost $s$-fractional harmonic functions). Let $r_{0}>0$ and $\omega:\left(0, r_{0}\right) \rightarrow$ $[0, \infty)$ be a modulus of continuity ${ }^{1}$. We say that a function $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$ is almost $s$-fractional harmonic in an open set $\Omega \subset \mathbb{R}^{n}$, with a gauge function $\omega$, if its reflected Caffarelli-Silvestre extension $u(x, y)$ is in $W_{\mathrm{loc}}^{1,2}\left(\widetilde{\Omega},|y|^{a}\right)$ and for any ball $\mathbb{B}_{r}\left(x_{0}\right)$ with $x_{0} \in \Omega$ and $0<r<r_{0}$ such that $B_{r}\left(x_{0}\right) \Subset \Omega$, we have

$$
\begin{equation*}
\int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla u|^{2}|y|^{a} \leq(1+\omega(r)) \int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla v|^{2}|y|^{a}, \tag{3.1.3}
\end{equation*}
$$

for any $v \in u+W_{0}^{1,2}\left(\mathbb{B}_{r}\left(x_{0}\right),|y|^{a}\right)$.

### 3.1.2 Fractional obstacle problem

A function $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$ is said to solve the $s$-fractional obstacle problem with obstacle $\psi$ in an open set $\Omega \subset \mathbb{R}^{n}$, if

$$
\begin{equation*}
\min \left\{\left(-\Delta_{x}\right)^{s} u, u-\psi\right\}=0 \quad \text { in } \Omega \tag{3.1.4}
\end{equation*}
$$

We refer to [11], [13], [61] for general introduction and basic results on this problem. With the help of the reflected Caffarelli-Silvestre extension, we can rewrite the problem as a Signorini-type problem for the operator $L_{a}$ :

$$
\begin{aligned}
L_{a} u=0 & \text { in } \mathbb{R}_{ \pm}^{n+1} \\
\min \left\{-\partial_{y}^{a} u, u-\psi\right\}=0 & \text { in } \Omega
\end{aligned}
$$

where

$$
\partial_{y}^{a} u(x, 0):=\lim _{y \rightarrow 0+} y^{a} \partial_{y} u(x, y) .
$$

This, in turn, can be written in the following variational form, see [13].
$\overline{{ }^{1} \text { i. } . \text {., a nondecreasing function with } \omega(0+)}=0$

Proposition 3.1.2. A function $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$ solves (3.1.4) if and only if its reflected CaffarelliSilvestre extension $u(x, y)$ is in $W_{\mathrm{loc}}^{1,2}(\widetilde{\Omega})$ and for any ball $\mathbb{B}_{r}\left(x_{0}\right)$ with $x_{0} \in \Omega$ such that $B_{r}\left(x_{0}\right) \Subset \Omega$, we have

$$
\int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla u|^{2}|y|^{a} \leq \int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla v|^{2}|y|^{a}
$$

for any $v \in \mathfrak{K}_{\psi, u}\left(\mathbb{B}_{r}\left(x_{0}\right),|y|^{a}\right):=\left\{v \in u+W_{0}^{1,2}\left(\mathbb{B}_{r}\left(x_{0}\right),|y|^{a}\right): v \geq \psi\right.$ on $\left.B_{r}\left(x_{0}\right)\right\}$.
Definition 3.1.2 (Almost minimizers for $s$-fractional obstacle problem). Let $r_{0}>0$ and $\omega:\left(0, r_{0}\right) \rightarrow[0, \infty)$ be a modulus of continuity. We say that a function $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$ is an almost minimizer for the $s$-fractional obstacle problem in an open set $\Omega \subset \mathbb{R}^{n}$, with a gauge function $\omega$, if its reflected Caffarelli-Silvestre extension $u(x, y)$ is in $W_{\operatorname{loc}}^{1,2}\left(\widetilde{\Omega},|y|^{a}\right)$ and for any ball $\mathbb{B}_{r}\left(x_{0}\right)$ with $x_{0} \in \Omega$ and $0<r<r_{0}$ such that $B_{r}\left(x_{0}\right) \Subset \Omega$, we have

$$
\begin{equation*}
\int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla u|^{2}|y|^{a} \leq(1+\omega(r)) \int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla v|^{2}|y|^{a}, \tag{3.1.5}
\end{equation*}
$$

for any $v \in \mathfrak{K}_{\psi, u}\left(\mathbb{B}_{r}\left(x_{0}\right),|y|^{a}\right)$.
The notion of almost minimizers above is related to the one for the thin obstacle problem $(s=1 / 2)$ studied in Chapter 1 , but there are certain important differences. In Definition 3.1.2, we ask the almost minimizing property (3.1.5) to hold only for balls centered on the "thin space" $\mathbb{R}^{n}$, while in Chapter 1 , we ask that property for balls centered at any point in an open set in the "thick space" $\mathbb{R}^{n+1}$. In a sense, this means that here we think of the perturbation from minimizers as living on the thin space, while in Chapter 1 they live in the thick space.

### 3.1.3 Main results and structure

In this chapter, our main concern is the regularity of almost minimizers in their original variables.

We start with examples of almost minimizers in Section 3.2. We then proceed to prove the following results, echoing those in [31] and Chapter 1.

Theorem J. Let $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$ be almost $s$-fractional harmonic in $\Omega$. Then
(i) $u$ is almost Lipschitz in $\Omega$, i.e, $u \in C^{0, \sigma}(\Omega)$ for any $0<\sigma<1$.
(ii) If $\omega(r)=r^{\alpha}$, then $u \in C^{1, \beta}(\Omega)$ for some $\beta=\beta_{n, a, \alpha}>0$.
(iii) If $0<s<1 / 2$ or $s=1 / 2$ and $\omega(r)=r^{\alpha}$ for some $\alpha>0$, then $u$ is actually $s$-fractional harmonic in $\Omega$.

In the case of the $s$-fractional obstacle problem, our results are obtained under the assumption that $1 / 2 \leq s<1$. Also, because of the technical nature of the problem, we restrict ourselves to the case $\psi=0$.

Theorem K. Let $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$ be an almost minimizer for the s-fractional obstacle problem with obstacle $\psi=0$ in $\Omega$.
(i) If $1 / 2 \leq s<1$, then $u \in C^{0, \sigma}(\Omega)$ for any $0<\sigma<1$.
(ii) If $1 / 2 \leq s<1$ and $\omega(r)=r^{\alpha}$ for some $\alpha>0$, then $u \in C^{1, \beta}(\Omega)$ for some $\beta=\beta_{n, a, \alpha}>$ 0.

The proofs follow the general approach in [31] and Chapter 1 by first obtaining growth estimates for minimizers (see Section 3.3) and then deriving their perturbed versions for almost minimizers (Section 3.4 for $s$-fractional harmonic functions and Section 3.5 for the $s$-fractional obstacle problem). The regularity then follows by an embedding theorem of a Morrey-Campanato-type space into the Hölder space, which we included in Appendix 3.A. Finally, Appendix 3.B contains the proof of orthogonal polynomial expansion of $L_{a}$-harmonic functions, that we rely on in deriving the growth estimates in Section 3.3. The polynomial expansion has other interesting corollaries such as the (known) real-analyticity of $s$-fractional harmonic functions, which are of independent interest.

### 3.1.4 Notation

Throughout this chapter we use the following notation. $\mathbb{R}^{n}$ is the $n$-dimensional Euclidean space. The points of $\mathbb{R}^{n+1}$ are denoted by $X=(x, y)$, where $x=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}, y \in \mathbb{R}$. We routinely identify $x \in \mathbb{R}^{n}$ with $(x, 0) \in \mathbb{R}^{n} \times\{0\} . \mathbb{R}_{ \pm}^{n+1}$ stands for open halfspaces $\left\{X=(x, y) \in \mathbb{R}^{n+1}: \pm y>0\right\}$.

We use the following notations for balls of radius $r$ in $\mathbb{R}^{n}$ and $\mathbb{R}^{n+1}$

$$
\begin{aligned}
\mathbb{B}_{r}(X) & =\left\{Z \in \mathbb{R}^{n+1}:|X-Z|<r\right\}, & & \text { (Euclidean) ball in } \mathbb{R}^{n+1} \\
\mathbb{B}_{r}^{ \pm}(x) & =\mathbb{B}_{r}(x, 0) \cap\{ \pm y>0\}, & & \text { half-ball in } \mathbb{R}^{n+1}, \\
B_{r}(x) & =\mathbb{B}_{r}(x, 0) \cap\{y=0\}, & & \text { ball in } \mathbb{R}^{n} .
\end{aligned}
$$

We typically drop the center from the notation if it is the origin. Thus, $\mathbb{B}_{r}=\mathbb{B}_{r}(0), B_{r}=$ $B_{r}(0)$, etc.

Next, $\nabla u=\nabla_{X} u=\left(\partial_{x_{1}} u, \ldots, \partial_{x_{n}} u, \partial_{y} u\right)$ stands for the full gradient, while $\nabla_{x} u=$ $\left(\partial_{x_{1}} u, \ldots, \partial_{x_{n}} u\right)$. We also use the standard notations for partial derivatives, such as $\partial_{x_{\mathrm{i}}} u$, $u_{x_{\mathrm{i}}}, u_{y}$ etc.

In integrals, we often drop the variable and the measure of integration if it with respect to the Lebesgue measure or the surface measure. Thus,

$$
\int_{\mathbb{B}_{r}} u|y|^{a}=\int_{\mathbb{B}_{r}} u(X)|y|^{a} d X, \quad \int_{\partial \mathbb{B}_{r}} u|y|^{a}=\int_{\partial \mathbb{B}_{r}} u(X)|y|^{a} d S_{X},
$$

where $S_{X}$ stands for the surface measure.
By $L^{2}\left(\mathbb{B}_{R},|y|^{a}\right)$ and $L^{2}\left(\partial \mathbb{B}_{R},|y|^{a}\right)$ we indicate the weighted Lebesgue spaces of functions with the norms

$$
\begin{aligned}
\|u\|_{L^{2}\left(\mathbb{B}_{R},|y|^{a}\right)}^{2} & =\int_{\mathbb{B}_{R}} u^{2}|y|^{a} \\
\|u\|_{L^{2}\left(\partial \mathbb{B}_{R},|y|^{a}\right)}^{2} & =\int_{\partial \mathbb{B}_{R}} u^{2}|y|^{a} .
\end{aligned}
$$

$W^{1,2}\left(\mathbb{B}_{R},|y|^{a}\right)$ is the corresponding weighted Sobolev space of functions with the norm

$$
\|u\|_{W^{1,2}\left(\mathbb{B}_{R},|y|^{a}\right)}^{2}=\|u\|_{L^{2}\left(\mathbb{B}_{R},|y|^{a}\right)}^{2}+\|\nabla u\|_{L^{2}\left(\mathbb{B}_{R},|y|^{a}\right)}^{2} .
$$

We also use other typical notations for Sobolev spaces. Thus, $W_{0}^{1,2}\left(\mathbb{B}_{R},|y|^{a}\right)$ stands for the closure of $C_{0}^{\infty}\left(\mathbb{B}_{R}\right)$ in $W^{1,2}\left(\mathbb{B}_{R},|y|^{a}\right)$.

For $x \in \mathbb{R}^{n}$ and $r>0$, we indicate by $\langle u\rangle_{x, r}$ the $|y|^{a}$-weighted integral mean value of a function $u$ over $\mathbb{B}_{r}(x)$. That is,

$$
\langle u\rangle_{x, r}=f_{\mathbb{B}_{r}(x)} u|y|^{a}:=\frac{1}{\omega_{n+1+a} r^{n+1+a}} \int_{\mathbb{B}_{r}(x)} u|y|^{a},
$$

where $\omega_{n+1+a}=\int_{\mathbb{B}_{1}}|y|^{a}$ is the $|y|^{a}$-weighted volume of the unit ball $\mathbb{B}_{1}$ in $\mathbb{R}^{n+1}$. (Note that here and throughout the thesis, the sign $f$ denotes the integral mean value with respect to the weighted measure $|y|^{a} d X$.) Finally, similarly to the other notations, we drop the origin if it is 0 and write $\langle u\rangle_{r}$ for $\langle u\rangle_{0, r}$.

### 3.2 Examples of almost minimizers

Before we proceed with the proofs of the main results, we would like to give some examples of almost minimizers.

Example 3.2.1. Let $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$ be a solution of

$$
\left(-\Delta_{x}\right)^{s} u+b(x) \cdot \nabla_{x} u=0 \quad \text { in } \Omega
$$

where $b=\left(b^{1}, b^{2}, \ldots, b^{n}\right) \in W^{1, \infty}(\Omega)$ and $1 / 2<s<1$ (or $-1<a<0$ ). Then $u$ is almost $s$-fractional harmonic with a gauge function $\omega(r)=C r^{-a}$ (note that $-a>0$ ).

Proof. Consider a ball $\mathbb{B}_{r}\left(x_{0}\right)$ centered at $x_{0} \in \Omega$ such that $B_{r}\left(x_{0}\right) \Subset \Omega$. Without loss of generality assume that $x_{0}=0$. Let $v$ be the minimizer of

$$
\int_{\mathbb{B}_{r}}|\nabla v|^{2}|y|^{a}
$$

on $u+W_{0}^{1,2}\left(\mathbb{B}_{r},|y|^{a}\right)$. Then

$$
\int_{\mathbb{B}_{r}} \nabla v \nabla(u-v)|y|^{a}=0
$$

and as a consequence,

$$
\int_{\mathbb{B}_{r}}\left(|\nabla u|^{2}-|\nabla v|^{2}\right)|y|^{a}=\int_{\mathbb{B}_{r}}|\nabla(u-v)|^{2}|y|^{a} .
$$

Then, we have

$$
\begin{aligned}
\int_{\mathbb{B}_{r}}\left(|\nabla u|^{2}-|\nabla v|^{2}\right)|y|^{a} & =2 \int_{\mathbb{B}_{r}^{+}}|\nabla(u-v)|^{2}|y|^{a} \\
& =2 \int_{\mathbb{B}_{r}^{+}}|\nabla(u-v)|^{2}|y|^{a}+\operatorname{div}\left(|y|^{a} \nabla(u-v)\right)(u-v) \\
& =2 \int_{\mathbb{B}_{r}^{+}} \operatorname{div}\left(|y|^{a} \nabla\left(\frac{(u-v)^{2}}{2}\right)\right) \\
& =2 \int_{\left(\partial \mathbb{B}_{r}\right)^{+}}|y|^{a}(u-v)\left(u_{\nu}-v_{\nu}\right)-2 \int_{\mathbb{B}_{r}}(u-v)\left(\partial_{y}^{a} u-\partial_{y}^{a} v\right) \\
& =C \int_{\mathbb{B}_{r}}(u-v)\left(-\Delta_{x}\right)^{s} u \\
& =-C \int_{\mathbb{B}_{r}}(u-v) b^{i} u_{x_{i}}
\end{aligned}
$$

with $C=C_{n, a}$. Next, extending $b^{i}$ to $\mathbb{R}^{n+1}$ by $b^{i}(x, y):=b^{i}(x)$, we have

$$
\begin{aligned}
\int_{\mathbb{B}_{r}}\left(|\nabla u|^{2}-|\nabla v|^{2}\right)|y|^{a}= & -C \int_{\mathbb{B}_{r}^{\prime}}(u-v) b^{i} u_{x_{i}} \\
= & C \int_{\mathbb{B}_{r}^{+}} \partial_{y}\left((u-v) b^{i} u_{x_{i}}\right) \\
= & C \int_{\mathbb{B}_{r}^{+}}\left(u_{y}-v y_{y}\right) b^{i} u_{x_{i}}+(u-v) b^{i} u_{x_{i} y} \\
\leq & C\|b\|_{W^{1, \infty}(\Omega)} \int_{\mathbb{B}_{r}^{+}}|\nabla u|^{2}+|\nabla v|^{2} \\
& +C \int_{\partial\left(\mathbb{B}_{r}^{+}\right)}(u-v) b^{i} u_{y} \nu_{x_{i}}-C \int_{\mathbb{B}_{r}^{+}} \partial_{x_{i}}\left((u-v) b^{i}\right) u_{y} \\
= & C\|b\|_{W^{1, \infty}(\Omega)} \int_{\mathbb{B}_{r}^{+}}|\nabla u|^{2}+|\nabla v|^{2} \\
& \quad-C \int_{\mathbb{B}_{r}^{+}}\left(\left(u_{x_{i}}-v_{x_{i}}\right) b^{i}+(u-v) b_{x_{i}}^{i}\right) u_{y} \\
\leq & C\|b\|_{W^{1, \infty}(\Omega)} \int_{\mathbb{B}_{r}^{+}}|\nabla u|^{2}+|\nabla v|^{2}+|u-v|^{2} .
\end{aligned}
$$

Using Poincare's inequality, it follows that

$$
\begin{aligned}
\int_{\mathbb{B}_{r}}\left(|\nabla u|^{2}-|\nabla v|^{2}\right)|y|^{a} & \leq C \int_{\mathbb{B}_{r}}|\nabla u|^{2}+|\nabla v|^{2} \leq C r^{-a} \int_{\mathbb{B}_{r}}\left(|\nabla u|^{2}+|\nabla v|^{2}\right)|y|^{a} \\
& \leq C r^{-a} \int_{\mathbb{B}_{r}}|\nabla u|^{2}|y|^{a} .
\end{aligned}
$$

Hence,

$$
\int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla u|^{2}|y|^{a} \leq\left(1+C r^{-a}\right) \int_{\mathbb{B}_{r}\left(x_{0}\right)}|\nabla v|^{2}|y|^{a},
$$

for $0<r<r_{0}$, with $C$ and $r_{0}$ depending on $n, a$, and $\|b\|_{W^{1, \infty}(\Omega)}$.
Example 3.2.2. Let $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$ be a solution of the obstacle problem for fractional Laplacian with drift

$$
\min \left\{\left(-\Delta_{x}\right)^{s} u+b(x) \cdot \nabla_{x} u, u\right\}=0 \quad \text { in } \Omega,
$$

where $b=\left(b^{1}, b^{2}, \ldots, b^{n}\right) \in W^{1, \infty}(\Omega)$ and $1 / 2<s<1$ (or $-1<a<0$ ). Then $u$ is an almost minimizer for $s$-fractional obstacle problem in $\Omega$ with an obstacle $\psi=0$ and a gauge function $\omega(r)=C r^{-a}$.

The obstacle problem above has been studied earlier in [17] and [57].

Proof. We argue similarly to Example 3.2.1. Let $\mathbb{B}_{r}\left(x_{0}\right)$ centered at $x_{0} \in \Omega$ such that $B_{r}\left(x_{0}\right) \Subset \Omega$. Without loss of generality assume that $x_{0}=0$. Let $v$ be the minimizer of

$$
\int_{\mathbb{B}_{r}}|\nabla v|^{2}|y|^{a}
$$

on $\mathfrak{K}_{0, u}\left(\mathbb{B}_{r},|y|^{a}\right)=\left\{v \in u+W_{0}^{1,2}\left(\mathbb{B}_{r},|y|^{a}\right): v \geq 0\right.$ on $\left.B_{r}\right\}$. Next, we write

$$
\begin{aligned}
\int_{\mathbb{B}_{r}}\left(|\nabla u|^{2}-|\nabla v|^{2}\right)|y|^{a} & =2 \int_{\mathbb{B}_{r}} \nabla u \nabla(u-v)|y|^{a}-\int_{\mathbb{B}_{r}}|\nabla(u-v)|^{2}|y|^{a} \\
& \leq 2 \int_{\mathbb{B}_{r}} \nabla u \nabla(u-v)|y|^{a} \\
& =4 \int_{\mathbb{B}_{r}^{+}} \nabla u \nabla(u-v)|y|^{a}+\operatorname{div}\left(|y|^{a} \nabla u\right)(u-v) \\
& =-4 \int_{B_{r}}(u-v) \partial_{y}^{a} u \\
& =C \int_{B_{r}}(u-v)\left(-\Delta_{x}\right)^{s} u \\
& =C\left[-\int_{B_{r} \cap\{u>0\}}(u-v) b^{\mathrm{i}} u_{x_{i}}+\int_{B_{r} \cap\{u=0\}}(-v)\left(-\Delta_{x}\right)^{s} u\right] \\
& \leq C\left[-\int_{B_{r} \cap\{u>0\}}(u-v) b^{i} u_{x_{i}}-\int_{B_{r} \cap\{u=0\}}(-v) b^{i} u_{x_{i}}\right]
\end{aligned}
$$

$$
=-C \int_{B_{r}}(u-v) b^{i} u_{x_{i}},
$$

where we used that $(-\Delta)^{s} u+b^{i} u_{x_{i}} \geq 0$ and $-v \leq 0$ on $B_{r} \cap\{u=0\}$ in the last inequality.
Then we complete the proof as in Example 3.2.1.

### 3.3 Growth estimates for minimizers

In this section we prove growth estimates for $L_{a}$-harmonic functions and solutions of the Signorini problem for $L_{a}$, i.e., minimizers $v$ of the weighted Dirichlet integral

$$
\int_{\mathbb{B}_{r}}|\nabla v|^{2}|y|^{a}
$$

on $v+W_{0}^{1,2}\left(\mathbb{B}_{r},|y|^{a}\right)$ or on the thin obstacle constraint set $\mathfrak{K}_{0, v}\left(\mathbb{B}_{r},|y|^{a}\right)$.
The idea is that these estimates will extend to almost minimizers and will ultimately imply their regularity with the help of Morrey-Campanato-type space embedding.

The proofs in this section are akin to those in Chapter 1 for almost minimizers of the thin obstacle problem. Yet, one has to be careful with different growth rates for tangential and normal derivatives.

### 3.3.1 Growth estimates for $L_{a}$-harmonic functions

Lemma 3.3.1. Let $v \in W^{1,2}\left(\mathbb{B}_{R},|y|^{a}\right)$ be a solution of $L_{a} v=0$ in $\mathbb{B}_{R}$. If $v$ is even in $y$, then for $0<\rho<R$

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v\right|^{2}|y|^{a} & \leq\left(\frac{\rho}{R}\right)^{n+1+a} \int_{\mathbb{B}_{R}}\left|\nabla_{x} v\right|^{2}|y|^{a} \\
\int_{\mathbb{B}_{\rho}}\left|v_{y}\right|^{2}|y|^{a} & \leq\left(\frac{\rho}{R}\right)^{n+3+a} \int_{\mathbb{B}_{R}}\left|v_{y}\right|^{2}|y|^{a} .
\end{aligned}
$$

Proof. Note that we can write

$$
v(x, y)=\sum_{k=0}^{\infty} p_{k}(x, y)
$$

where $p_{k}$ 's are $L_{a}$-harmonic homogeneous polynomials of degree $k$ (see Appendix 3.B). Then $\left\{\partial_{x_{\mathrm{i}}} p_{k}\right\}_{k=1}^{\infty}$ are $L_{a}$-harmonic homogeneous polynomials of degree $k-1$, and thus orthogonal in $L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)$. Thus,

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v\right|^{2}|y|^{a} & =\sum_{k=1}^{\infty} \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} p_{k}\right|^{2}|y|^{a}=\sum_{k=1}^{\infty}\left(\frac{\rho}{R}\right)^{n+1+a+2(k-1)} \int_{\mathbb{B}_{R}}\left|\nabla_{x} p_{k}\right|^{2}|y|^{a} \\
& \leq\left(\frac{\rho}{R}\right)^{n+1+a} \sum_{k=1}^{\infty} \int_{\mathbb{B}_{R}}\left|\nabla_{x} p_{k}\right|^{2}|y|^{a}=\left(\frac{\rho}{R}\right)^{n+1+a} \int_{\mathbb{B}_{R}}\left|\nabla_{x} v\right|^{2}|y|^{a} .
\end{aligned}
$$

Similarly, $\left\{|y|^{a} \partial_{y} p_{k}\right\}_{k=1}^{\infty}$ are $L_{-a}$-harmonic homogeneous functions of degree $k-1+a$, and thus orthogonal in $L^{2}\left(\partial \mathbb{B}_{1},|y|^{-a}\right)$. Notice that since $p_{1}(x, y)=p_{1}(x)$ is independent of $y$ variable by the even symmetry, we have $|y|^{a} \partial_{y} p_{1}=0$. Thus,

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|v_{y}\right|^{2}|y|^{a} & =\left.\left.\int_{\mathbb{B}_{\rho}}| | y\right|^{a} v_{y}\right|^{2}|y|^{-a}=\left.\left.\sum_{k=2}^{\infty} \int_{\mathbb{B}_{\rho}}| | y\right|^{a} \partial_{y} p_{k}\right|^{2}|y|^{-a} \\
& =\left.\sum_{k=2}^{\infty}\left(\frac{\rho}{R}\right)^{n+1-a+2(k-1+a)} \int_{\mathbb{B}_{R}}|y|^{a} \partial_{y} p_{k}\right|^{2}|y|^{-a} \leq\left(\frac{\rho}{R}\right)^{n+3+a} \int_{\mathbb{B}_{R}}\left|v_{y}\right|^{2}|y|^{a} .
\end{aligned}
$$

Lemma 3.3.2. Let $v$ be a solution of $L_{a} v=0$ in $\mathbb{B}_{R}$, even in $y$. Then, for $0<\rho<R$,

$$
\begin{equation*}
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \leq\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{R}\right|^{2}|y|^{a} . \tag{3.3.1}
\end{equation*}
$$

Proof. First note that since $L_{a}\left(\nabla_{x} v\right)=0$ in $\mathbb{B}_{R},\left\langle\nabla_{x} v\right\rangle=\nabla_{x} v(0)$ by the mean value theorem for $L_{a}$-harmonic functions, see Lemma 2.9 in [13]. If we use the expansion $v=\sum_{k=0}^{\infty} p_{k}(x, y)$ in $\mathbb{B}_{R}$ as in the proof of Lemma 3.3.1, then $\nabla_{x} v-\nabla_{x} v(0)=\sum_{k=2}^{\infty} \nabla_{x} p_{k}$ and consequently

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\nabla_{x} v(0)\right|^{2}|y|^{a} & =\sum_{k=2}^{\infty} \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} p_{k}\right|^{2}|y|^{a}=\sum_{k=2}^{\infty}\left(\frac{\rho}{R}\right)^{n+a+2 k-1} \int_{\mathbb{B}_{R}}\left|\nabla_{x} p_{k}\right|^{2}|y|^{a} \\
& \leq\left(\frac{\rho}{R}\right)^{n+a+3} \sum_{k=2}^{\infty} \int_{\mathbb{B}_{R}}\left|\nabla_{x} p_{k}\right|^{2}|y|^{a} \\
& =\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} v-\nabla_{x} v(0)\right|^{2}|y|^{a} .
\end{aligned}
$$

### 3.3.2 Growth estimates for the solutions of the Signorini problem for $L_{a}$

Our estimates for the solutions of the Signorini problem will require an assumption that $1 / 2 \leq s<1$, or $a \leq 0$. Also, unless stated otherwise, the obstacle $\psi$ is assumed to be zero.

The first estimate is the analogue of Lemma 3.3.1, but with less information of the growth of $v_{y}$.

Lemma 3.3.3. Let $v$ be a solution of the Signorini problem for $L_{a}$ in $\mathbb{B}_{R}$, even in $y$, with $a \leq 0$. Then, for $0<\rho<R$

$$
\begin{equation*}
\int_{\mathbb{B}_{\rho}}|\nabla v|^{2}|y|^{a} \leq\left(\frac{\rho}{R}\right)^{n+1+a} \int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a} . \tag{3.3.2}
\end{equation*}
$$

Proof. We use the following property: if $v$ is as in the statement of the lemma, then $v_{x_{i}}$, $i=1, \ldots, n$, and $y|y|^{a-1} v_{y}$ are Hölder continuous in $\mathbb{B}_{R}$, see [13]. Moreover, we have that

$$
L_{a}\left(v_{x_{i}}^{ \pm}\right) \geq 0, \quad L_{-a}\left(\left(y|y|^{a-1} v_{y}\right)^{ \pm}\right) \geq 0 \quad \text { in } \mathbb{B}_{R} .
$$

This follows from the fact that $L_{a} v_{x_{i}}=0$ in $\left\{ \pm v_{x_{i}}>0\right\}$ and $L_{-a}\left(y|y|^{a-1} v_{y}\right)=0$ in $\left\{ \pm y|y|^{a-1} v_{y}>0\right\}$, by the complementarity condition $v_{y} v=0$ on $B_{R}$, as well as an argument in Exercise 2.6 or Exercise 9.5 in [48]. As a consequence, we have

$$
L_{a}\left(\left|\nabla_{x} v\right|^{2}\right) \geq 0, \quad L_{-a}\left(\left.\left.| | y\right|^{a} v_{y}\right|^{2}\right) \geq 0 \quad \text { in } \mathbb{B}_{R}
$$

We next use the following $|y|^{a}$-weighted sub-mean value property for $L_{a}$-subharmonic functions: If $L_{a} w \geq 0$ weakly in $\mathbb{B}_{R},-1<a<1$, then

$$
\rho \mapsto \frac{1}{\rho^{n+1+a}} \int_{B_{\rho}} w|y|^{a}
$$

is nondecreasing. This follows by integration from the spherical sub-mean value property, see Lemma 2.9 in [13]. Thus, we have that

$$
\rho \mapsto \frac{1}{\rho^{n+1+a}} \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v\right|^{2}|y|^{a}
$$

$$
\rho \mapsto \frac{1}{\rho^{n+1-a}} \int_{\mathbb{B}_{\rho}}|y|^{a} u_{y}^{2}
$$

are monotone nondecreasing for $0<\rho<R$. This implies

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v\right|^{2}|y|^{a} & \leq\left(\frac{\rho}{R}\right)^{n+1+a} \int_{\mathbb{B}_{R}}\left|\nabla_{x} v\right|^{2}|y|^{a} \\
\int_{\mathbb{B}_{\rho}} v_{y}^{2}|y|^{a} & \leq\left(\frac{\rho}{R}\right)^{n+1-a} \int_{\mathbb{B}_{R}} v_{y}^{2}|y|^{a} .
\end{aligned}
$$

In the case $a \leq 0$, we therefore conclude that the bound (3.3.2) holds.
Lemma 3.3.4. Let $v$ be a solution of the Signorini problem for $L_{a}$ in $\mathbb{B}_{R}$, even in $y$, with $a \leq 0$. If $v(0)=0$, then there exists $C=C_{n, \alpha}$ such that for $0<\rho<r<(3 / 4) R$,

$$
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \leq\left(\frac{\rho}{r}\right)^{n+a+3} \int_{\mathbb{B}_{r}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{r}\right|^{2}|y|^{a}+C \frac{\rho^{n+2}}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right)}^{2}
$$

Proof. Define

$$
\varphi(r):=\frac{1}{r^{n+a+3}} \int_{\mathbb{B}_{r}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{r}\right|^{2}|y|^{a} .
$$

Then,

$$
\begin{aligned}
\varphi(r) & =\frac{1}{r^{n+a+3}}\left[\int_{\mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a}-2\left\langle\nabla_{x} v\right\rangle_{r} \int_{\mathbb{B}_{r}} \nabla_{x} v|y|^{a}+\left\langle\nabla_{x} v\right\rangle_{r}^{2} \int_{\mathbb{B}_{r}}|y|^{a}\right] \\
& =\frac{1}{r^{n+a+3}}\left[\int_{\mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a}-\frac{1}{\omega_{n+1+a} r^{n+1+a}}\left(\int_{\mathbb{B}_{r}} \nabla_{x} v|y|^{a}\right)^{2}\right] .
\end{aligned}
$$

Thus, using the Cauchy-Schwarz and Young's inequality, we obtain

$$
\begin{aligned}
\varphi^{\prime}(r)= & \frac{1}{r^{n+a+3}}\left[-\frac{n+a+3}{r} \int_{\mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a}+\int_{\partial \mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a}\right. \\
& +\frac{n+a+3}{\omega_{n+1+a} r^{n+2+a}}\left(\int_{\mathbb{B}_{r}} \nabla_{x} v|y|^{a}\right)^{2}+\frac{n+1+a}{\omega_{n+1+a} r^{n+2+a}}\left(\int_{\mathbb{B}_{r}} \nabla_{x} v|y|^{a}\right)^{2} \\
& \left.-\frac{2}{\omega_{n+1+a} r^{n+1+a}}\left(\int_{\mathbb{B}_{r}} \nabla_{x} v|y|^{a}\right)\left(\int_{\partial \mathbb{B}_{r}} \nabla_{x} v|y|^{a}\right)\right] \\
\geq- & \frac{C}{r^{n+a+3}}\left[\frac{1}{r} \int_{\mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a}+\left(\frac{1}{r} \int_{\mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a}\right)^{1 / 2}\left(\int_{\partial \mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a}\right)^{1 / 2}\right] \\
\geq- & \frac{C}{r^{n+a+3}}\left[\frac{1}{r} \int_{\mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a}+\int_{\partial \mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a}\right] .
\end{aligned}
$$

Next, we note that

$$
\left[\nabla_{x} v\right]_{C^{0, s}\left(\mathbb{B}_{3 / 4 R}\right)} \leq \frac{C_{n, s}}{R^{1+s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right)} .
$$

Indeed, this follows from the known interior regularity for solutions of the Signorini problem for $L_{a}$ in $\mathbb{B}_{1}$ in the case $R=1$, see e.g. [13], and a simple scaling argument for all $R>0$. Noting also that $\nabla_{x} v(0)=0$, since $v$ attains its minimum on $B_{r}$ at 0 , we have that for $X \in \overline{\mathbb{B}_{r}}$ with $r<(3 / 4) R$

$$
\left|\nabla_{x} v(X)\right|=\left|\nabla_{x} v(X)-\nabla_{x} v(0)\right| \leq \frac{C}{R^{1+s}} r^{s}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right)}
$$

and so

$$
\frac{1}{r} \int_{\mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a}+\int_{\partial \mathbb{B}_{r}}\left|\nabla_{x} v\right|^{2}|y|^{a} \leq C \frac{r^{n+1}}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right)}^{2} .
$$

This gives

$$
\varphi^{\prime}(r) \geq-\frac{C}{r^{a+2}} \frac{1}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right)}^{2} .
$$

Thus, for $0<\rho<r<(3 / 4) R$,

$$
\varphi(r)-\varphi(\rho)=\int_{\rho}^{r} \varphi^{\prime}(t) d t \geq-C \frac{\rho^{-1-a}-r^{-1-a}}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right)}^{2}
$$

Therefore,

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} & =\rho^{n+a+3} \varphi(\rho) \\
& \leq \rho^{n+a+3}\left(\varphi(r)+C \frac{\rho^{-1-a}-r^{-1-a}}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right)}^{2}\right) \\
& \leq\left(\frac{\rho}{r}\right)^{n+a+3} \int_{\mathbb{B}_{r}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{r}\right|^{2}|y|^{a}+C \frac{\rho^{n+2}}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right)}^{2} .
\end{aligned}
$$

Lemma 3.3.5. Let $v$ be a solution of the Signorini problem for $L_{a}$ in $\mathbb{B}_{R}$, even in $y$. Then there are $C_{1}=C_{n, a}, C_{2}=C_{n, a}$ such that for all $0<\rho<S<(3 / 8) R$,

$$
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \leq C_{1}\left(\frac{\rho}{S}\right)^{n+a+3} \int_{\mathbb{B}_{S}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{S}\right|^{2}|y|^{a}+C_{2} \frac{S^{n+2}}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right)}^{2} .
$$

Proof. If $\rho \geq S / 8$, then we immediately have

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} & \leq C\left(\frac{8 \rho}{S}\right)^{n+a+3} \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \\
& \leq C\left(\frac{\rho}{S}\right)^{n+a+3} \int_{\mathbb{B}_{S}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{S}\right|^{2}|y|^{a} .
\end{aligned}
$$

Thus we may assume $\rho<S / 8$. Due to Lemma 3.3.4, we may assume $v(0)>0$. Let $d:=\operatorname{dist}(0,\{v(\cdot, 0)=0\})>0$. Then $L_{a} v=0$ in $\mathbb{B}_{d}$. Thus, if $d \geq S$, we may use Lemma 3.3.2 to obtain

$$
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \leq\left(\frac{\rho}{S}\right)^{n+a+3} \int_{\mathbb{B}_{S}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{S}\right|^{2}|y|^{a} .
$$

Thus we may also assume $d<S$.
Case 1. $S / 4 \leq d(<S)$.
Case 1.1. Suppose $0<\rho<d(<S)$. Then using $L_{a}\left(\nabla_{x} v\right)=0$ in $\mathbb{B}_{d}$ again,

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} & \leq\left(\frac{\rho}{d}\right)^{n+a+3} \int_{\mathbb{B}_{d}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{d}\right|^{2}|y|^{a} \\
& \leq C\left(\frac{\rho}{S}\right)^{n+a+3} \int_{\mathbb{B}_{S}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{S}\right|^{2}|y|^{a} .
\end{aligned}
$$

Case 1.2. Suppose $\rho \geq d(\geq S / 4)$. Then

$$
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \leq\left(\frac{4 \rho}{S}\right)^{n+a+3} \int_{\mathbb{B}_{S}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{S}\right|^{2}|y|^{a}
$$

Case 2. $0<d<S / 4$.
Case 2.1. Suppose $\rho<d / 2$. Take $x_{1} \in \partial\left(B_{d}\right)$ such that $v\left(x_{1}\right)=0$. Then using inclusions $\mathbb{B}_{\rho} \subset \mathbb{B}_{d / 2} \subset \mathbb{B}_{(3 / 2) d}\left(x_{1}\right) \subset \mathbb{B}_{S / 2}\left(x_{1}\right) \subset \mathbb{B}_{R / 2}\left(x_{1}\right), L_{a} v=0$ in $\mathbb{B}_{d}$ and the preceding Lemma 3.3.4, we obtain

$$
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a}
$$

$$
\begin{aligned}
& \leq\left(\frac{2 \rho}{d}\right)^{n+a+3} \int_{\mathbb{B}_{d / 2}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{d / 2}\right|^{2}|y|^{a} \\
& \leq\left(\frac{2 \rho}{d}\right)^{n+a+3} \int_{\mathbb{B}_{(3 / 2) d}\left(x_{1}\right)}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{x_{1},(3 / 2) d}\right|^{2}|y|^{a} \\
& \leq\left(\frac{2 \rho}{d}\right)^{n+a+3}\left[\left(\frac{3 d}{S}\right)^{n+a+3} \int_{\mathbb{B}_{S / 2}\left(x_{1}\right)}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{x_{1}, S / 2}\right|^{s}|y|^{a}+C \frac{S^{n+2}}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R / 2}\left(x_{1}\right)\right)}^{2}\right] \\
& \leq C\left(\frac{\rho}{S}\right)^{n+a+3} \int_{\mathbb{B}_{S}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{S}\right|^{2}|y|^{a}+C \frac{S^{n+2}}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right) .}^{2} .
\end{aligned}
$$

Case 2.2. Suppose $d / 2 \leq \rho$. Then we see that $\mathbb{B}_{\rho} \subset \mathbb{B}_{3 \rho}\left(x_{1}\right) \subset \mathbb{B}_{S / 2}\left(x_{1}\right) \subset \mathbb{B}_{S}$. As we did in Case 2.1, we have

$$
\begin{aligned}
& \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \\
& \leq \int_{\mathbb{B}_{3 \rho}\left(x_{1}\right)}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{x_{1}, 3 \rho}\right|^{2}|y|^{a} \\
& \leq C\left(\frac{\rho}{S}\right)^{n+a+3} \int_{\mathbb{B}_{S / 2}\left(x_{1}\right)}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{x_{1}, S / 2}\right|^{2}|y|^{a}+C \frac{S^{n+2}}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R / 2}\left(x_{1}\right)\right)}^{2} \\
& \leq C\left(\frac{\rho}{S}\right)^{n+a+3} \int_{\mathbb{B}_{S}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{S}\right|^{2}|y|^{a}+C \frac{S^{n+2}}{R^{2+2 s}}\|v\|_{L^{\infty}\left(\mathbb{B}_{R}\right)}^{2} .
\end{aligned}
$$

Corollary 3.3.6. Let $v$ be a solution of the Signorini problem for $L_{a}$ in $\mathbb{B}_{R}$, even in $y$. Then there are $C_{1}=C_{n, a}, C_{2}=C_{n, a}$ such that for all $0<\rho<S<(3 / 16) R$,

$$
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \leq C_{1}\left(\frac{\rho}{S}\right)^{n+a+3} \int_{\mathbb{B}_{S}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{S}\right|^{2}|y|^{2}+C_{2} \frac{S^{n+2}}{R^{2+2 s}}\left\langle v^{2}\right\rangle_{R} .
$$

Proof. Since $v^{ \pm}=\max ( \pm v, 0) \geq 0$ and $L_{a}\left(v^{ \pm}\right)=0$ in $\left\{v^{ \pm}>0\right\}$, we have $L_{a}\left(v^{ \pm}\right) \geq 0$ in $\mathbb{B}_{R}$. (For this, one may follow the argument in Exercise 2.6 or Exercise 9.5 in [48].) Thus, we have by Theorem 2.3.1 in [62]

$$
\sup _{\mathbb{B}_{R / 2}} v^{ \pm} \leq C\left(\frac{1}{\omega_{n+1+a} R^{n+1+a}} \int_{\mathbb{B}_{R}}\left(v^{ \pm}\right)^{2}|y|^{a}\right)^{1 / 2} .
$$

Hence,

$$
\|v\|_{L^{\infty}\left(\mathbb{B}_{R / 2}\right)}^{2} \leq C\left\langle v^{2}\right\rangle_{R}
$$

which completes the proof.

### 3.4 Almost $s$-fractional harmonic functions

In this section we prove Theorem J, by deducing growth estimates for almost $s$-fractional harmonic functions from that of $s$-fractional harmonic functions and then applying the Morrey-Campanato space embedding to deduce the regularity of almost $s$-fractional harmonic functions.

Theorem 3.4.1 (Almost Lipschitz regularity). If $u$ is an almost $s$-fractional harmonic function in $B_{1}, 0<s<1$, then $u \in C^{0, \sigma}\left(B_{1}\right)$ for any $0<\sigma<1$.

Proof. Let $K$ be a compact subset of $B_{1}$ containing 0 . Take $\delta=\delta_{n, \omega, \sigma, K}>0$ such that $\delta<$ $\operatorname{dist}\left(K, \partial B_{1}\right)$ and $\omega(\delta) \leq \varepsilon$, where $\varepsilon=\varepsilon_{2, n+1+a, n-1+a+2 \sigma}$ is as Lemma 1.2.2. For $0<R<\delta$, let $v$ be a minimizer of

$$
\int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a}
$$

on $u+W_{0}^{1,2}\left(\mathbb{B}_{R},|y|^{a}\right)$. Then $L_{a} v=0$ in $\mathbb{B}_{R}$. In particular,

$$
\int_{\mathbb{B}_{R}} \nabla v \cdot \nabla(u-v)|y|^{a}=0
$$

and hence

$$
\begin{aligned}
\int_{\mathbb{B}_{R}}|\nabla(u-v)|^{2}|y|^{a} & =\int_{\mathbb{B}_{R}}|\nabla u|^{2}|y|^{a}-\int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a}-2 \int_{\mathbb{B}_{R}} \nabla v \cdot \nabla(u-v)|y|^{a} \\
& \leq \omega(R) \int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a} .
\end{aligned}
$$

Moreover, by Lemma 3.3.1, for $0<\rho<R$ we have

$$
\int_{\mathbb{B}_{\rho}}|\nabla v|^{2}|y|^{a} \leq\left(\frac{\rho}{R}\right)^{n+1+a} \int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a} .
$$

Thus

$$
\int_{\mathbb{B}_{\rho}}|\nabla u|^{2}|y|^{a} \leq 2 \int_{\mathbb{B}_{\rho}}|\nabla v|^{2}|y|^{a}+2 \int_{\mathbb{B}_{\rho}}|\nabla(u-v)|^{2}|y|^{a}
$$

$$
\begin{aligned}
& \leq 2\left(\frac{\rho}{R}\right)^{n+1+a} \int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a}+2 \int_{\mathbb{B}_{\rho}}|\nabla(u-v)|^{2}|y|^{a} \\
& \leq 2\left(\frac{\rho}{R}\right)^{n+1+a} \int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a}+2 \omega(R) \int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a} \\
& \leq 2\left[\left(\frac{\rho}{R}\right)^{n+1+a}+\varepsilon\right] \int_{\mathbb{B}_{R}}|\nabla u|^{2}|y|^{a} .
\end{aligned}
$$

By Lemma 1.2.2,

$$
\int_{\mathbb{B}_{\rho}}|\nabla u|^{2}|y|^{a} \leq C_{n, a, \sigma}\left(\frac{\rho}{R}\right)^{n-1+a+2 \sigma} \int_{\mathbb{B}_{R}}|\nabla u|^{2}|y|^{a},
$$

for any $0<\sigma<1$. Taking $R \nearrow \delta$ we have

$$
\begin{equation*}
\int_{\mathbb{B}_{\rho}}|\nabla u|^{2}|y|^{a} \leq C_{n, a, \sigma, \delta}\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \rho^{n-1+a+2 \sigma} . \tag{3.4.1}
\end{equation*}
$$

By weighted Poincaré inequality (Theorem 1.5 in [62])

$$
\int_{\mathbb{B}_{\rho}}\left|u-\langle u\rangle_{\rho}\right|^{2}|y|^{a} \leq C_{n, a, \sigma, \delta}\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \rho^{n+1+a+2 \sigma} .
$$

Now, a similar estimates holds at all point $x_{0} \in K$, which implies the Hölder continuity of $u$ (see Theorem 3.A.1) with

$$
\|u\|_{C^{0, \sigma}(K)} \leq C_{n, a, \omega, \sigma, K}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)} .
$$

Theorem 3.4.2 ( $C^{1, \beta}$ regularity). If $u$ is an almost s-fractional harmonic function in $B_{1}$, $0<s<1$, with gauge function $\omega(r)=r^{\alpha}$, $\alpha>0$, then $\nabla_{x} u \in C^{0, \beta}\left(B_{1}\right)$ for some $\beta=$ $\beta(n, s, \alpha)$.

Proof. Let $K \Subset B_{1}$ be a ball and take $0<\delta<\operatorname{dist}\left(K, \partial B_{1}\right)$. Let $B_{R}\left(x_{0}\right) \Subset B_{1}$ with $0<R<\delta$, for $x_{0} \in K$. For simplicity write $x_{0}=0$, and let $v$ be the $L_{a}$-harmonic function in $\mathbb{B}_{R}$ with $v=u$ on $\partial \mathbb{B}_{R}$. Then, by Jensen's inequality we have

$$
\int_{\mathbb{B}_{\rho}}\left|\left\langle\nabla_{x} u\right\rangle_{\rho}-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \leq \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a},
$$

and hence

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{\rho}\right|^{2}|y|^{a} \leq & 3 \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a}+3 \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} \\
& +3 \int_{\mathbb{B}_{\rho}}\left|\left\langle\nabla_{x} u\right\rangle_{\rho}-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \\
\leq & 3 \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a}+6 \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} .
\end{aligned}
$$

Similarly,

$$
\int_{\mathbb{B}_{R}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{R}\right|^{2}|y|^{a} \leq 3 \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{R}\right|^{2}|y|^{a}+6 \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} .
$$

Next let $\beta \in(0, \alpha / 2)$. Then using the estimate (3.4.1) in the proof of Theorem 3.4.1 with $\sigma=1+\beta-\frac{\alpha}{2}$, we have

$$
\begin{aligned}
\int_{\mathbb{B}_{R}}|\nabla u-\nabla v|^{2}|y|^{a} & =\int_{\mathbb{B}_{R}}|\nabla u|^{2}|y|^{a}-\int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a} \leq R^{\alpha} \int_{\mathbb{B}_{R}}|\nabla u|^{2}|y|^{a} \\
& \leq C\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} R^{n+1+a+2 \beta} .
\end{aligned}
$$

Then, with the help of Lemma 3.3.2, we have that for $\rho<R$

$$
\begin{aligned}
& \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{\rho}\right|^{2}|y|^{a} \\
& \leq C \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a}+C \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} \\
& \leq C\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{R}\right|^{2}|y|^{a}+C \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} \\
& \leq C\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{R}\right|^{2}|y|^{a}+C \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} \\
& \leq C\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{R}\right|^{2}|y|^{a}+C\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} R^{n+1+a+2 \beta} .
\end{aligned}
$$

Hence, by Lemma 1.2.2, we obtain that for $\rho<R$

$$
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{\rho}\right|^{2}|y|^{a}
$$

$$
\leq C\left[\left(\frac{\rho}{R}\right)^{n+1+a+2 \beta} \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{R}\right|^{2}|y|^{a}+\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \rho^{n+1+a+2 \beta}\right] .
$$

Taking $R \nearrow \delta$, we have

$$
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{\rho}\right|^{2}|y|^{a} \leq C_{n, a, \alpha, \beta, K}\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \rho^{n+1+a+2 \beta} .
$$

Now, a similar estimate holds for any $x_{0} \in K$. Fixing $\beta$ and applying Theorem 3.A.1, we have

$$
\left\|\nabla_{x} u\right\|_{C^{0, \beta}(K)} \leq C_{n, a, \alpha, K}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)} .
$$

Remark 3.4.3. From the assumption for almost minimizers that the Caffarelli-Silvestre extension $u \in W_{\text {loc }}^{1,2}$ we know only that $\nabla_{x} u \in L_{\text {loc }}^{2}$, which is not sufficient to deduce the existence of the trace of $\nabla_{x} u$ on $B_{1}$. However, in the proof of Theorem 3.4.2 we showed that $\nabla_{x} u$ is in a Morrey-Campanato space, which implies the existence of the trace as the limit of averages

$$
T\left(\nabla_{x} u\right)\left(x_{0}\right)=\lim _{r \rightarrow 0+}\left\langle\nabla_{x} u\right\rangle_{x_{0}, r} .
$$

It is not hard to see that $T\left(\nabla_{x} u\right)$ is the distributional derivative $\nabla_{x} u$ on $B_{1}$. Indeed, if $\eta \in C_{0}^{\infty}\left(B_{1}\right)$, then extending it to $\mathbb{R}^{n+1}$ by $\eta(x, y)=\eta(x)$, we have

$$
\begin{aligned}
\int_{B_{1}} T\left(\partial_{x_{\mathrm{i}}} u\right) \eta & =\lim _{r \rightarrow 0+} \int_{B_{1}}\left\langle\partial_{x_{\mathrm{i}}} u\right\rangle_{x, r} \eta=\lim _{r \rightarrow 0+} \int_{B_{1}} \partial_{x_{\mathrm{i}}} u\langle\eta\rangle_{x, r} \\
& =\lim _{r \rightarrow 0+}-\int_{B_{1}} u\left\langle\partial_{x_{\mathrm{i}}} \eta\right\rangle_{x, r}=-\int_{B_{1}} u \partial_{x_{\mathrm{i}}} \eta .
\end{aligned}
$$

Theorem 3.4.4. Let $u$ be an almost s-fractional harmonic function in $B_{1}$ for $0<s<1 / 2$ or $s=1 / 2$ and a gauge function $\omega(r)=r^{\alpha}$ for some $\alpha>0$. Then $u$ is actually $s$-fractional harmonic in $B_{1}$.

Proof. We argue as in the proof Theorem 3.4.1. Let $K, \delta, R, v$ be as in the proof of that theorem. Then, by Lemma 3.3.1, for $0<\rho<R$

$$
\int_{\mathbb{B}_{\rho}}\left|v_{y}\right|^{2}|y|^{a} \leq\left(\frac{\rho}{R}\right)^{n+3+a} \int_{\mathbb{B}_{R}}\left|v_{y}\right|^{2}|y|^{a} .
$$

Thus, for any $0<\sigma<1$, we have

$$
\begin{aligned}
\left.\left.\int_{\mathbb{B}_{\rho}}| | y\right|^{a} u_{y}\right|^{2}|y|^{-a} & \leq 2 \int_{\mathbb{B}_{\rho}}\left|v_{y}\right|^{2}|y|^{a}+2 \int_{\mathbb{B}_{\rho}}\left|u_{y}-v_{y}\right|^{2}|y|^{a} \\
& \leq 2\left(\frac{\rho}{R}\right)^{n+3+a} \int_{\mathbb{B}_{R}}\left|v_{y}\right|^{2}|y|^{a}+2 \int_{\mathbb{B}_{\rho}}\left|u_{y}-v_{y}\right|^{2}|y|^{a} \\
& \leq 4\left(\frac{\rho}{R}\right)^{n+3+a} \int_{\mathbb{B}_{R}}\left|u_{y}\right|^{2}|y|^{a}+6 \int_{\mathbb{B}_{R}}\left|u_{y}-v_{y}\right|^{2}|y|^{a} \\
& \leq\left.\left. 4\left(\frac{\rho}{R}\right)^{n+3+a} \int_{\mathbb{B}_{R}}| | y\right|^{a} u_{y}\right|^{2}|y|^{-a}+6 \omega(R) \int_{\mathbb{B}_{R}}|\nabla u|^{2}|y|^{a} \\
& \leq\left.\left. 4\left(\frac{\rho}{R}\right)^{n+3+a} \int_{\mathbb{B}_{R}}| | y\right|^{a} u_{y}\right|^{2}|y|^{-a}+C_{n, a, \sigma, \delta} \omega(R)\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)} R^{n-1+a+2 \sigma},
\end{aligned}
$$

where we used (3.4.1) in the last inequality.
Consider now the two cases in statement of the theorem.
Case 1. $0<s<1 / 2$ (or $a>0$ ). In this case by Lemma 1.2.2,

$$
\begin{aligned}
& \left.\left.\int_{\mathbb{B}_{\rho}}| | y\right|^{a} u_{y}\right|^{2}|y|^{-a} \\
& \quad \leq C\left[\left.\left.\left(\frac{\rho}{R}\right)^{n-1+a+2 \sigma} \int_{\mathbb{B}_{R}}| | y\right|^{a} u_{y}\right|^{2}|y|^{-a}+\omega(\delta)\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \rho^{n-1+a+2 \sigma}\right] \\
& \quad \leq C\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \rho^{n+1-a+(-2+2 a+2 \sigma)} .
\end{aligned}
$$

Now we take $\sigma=1-a / 2 \in(0,1)$ to have $-2+2 a+2 \sigma=a>0$. Varying the center, we have a similar bound at every $x \in K$. Then, by Theorem 3.A.1, we obtain that the limit of the averages $T\left(y|y|^{a-1} u_{y}\right)=0$ on $B_{1}$. This implies that $\left(-\Delta_{x}\right)^{s} u=0$ on $B_{1}$. Indeed, arguing as in Remark 3.4.3, by considering the mollifications $u_{\varepsilon}$ in $x$-variable, we note that

$$
\left.\left.\int_{\mathbb{B}_{\rho}}| | y\right|^{a}\left(u_{\varepsilon}\right)_{y}\right|^{2}|y|^{-a} \leq C \rho^{n+1-a+a}
$$

which implies that $T\left(y|y|^{a-1}\left(u_{\varepsilon}\right)_{y}\right)=0$ on $K \Subset B_{1}$. On the other hand, $u_{\varepsilon} \in C^{2} \cap \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$, which implies that $y|y|^{a-1}\left(u_{\varepsilon}\right)_{y}$ is continuous up to $y=0$, since we can explicitly write, for $y>0$, the symmetrized formula

$$
y^{a}\left(u_{\varepsilon}\right)_{y}(x, y)=\int_{\mathbb{R}^{n}} \frac{u_{\varepsilon}(x+z)+u_{\varepsilon}(x-z)-2 u_{\varepsilon}(x)}{|z|^{2}}|z|^{2} y^{a} \partial_{y} P(z, y) d z
$$

with locally integrable kernel $|z|^{2}\left|y^{a} \partial_{y} P(z, y)\right| \leq C /|z|^{n-1-a}$. Hence, we obtain that $\left(-\Delta_{x}\right)^{s} u_{\varepsilon}=$ $\partial_{y}^{a} u_{\varepsilon}=0$ on the ball $K \Subset B_{1}$. Then, passing to the limit as $\varepsilon \rightarrow 0$, this implies that $\left(-\Delta_{x}\right)^{s} u=0$ in $B_{1}$.

Case 2. $s=1 / 2($ or $a=0)$ and $\omega(r)=r^{\alpha}$. In this case, we have a bound

$$
\int_{\mathbb{B}_{\rho}}\left|u_{y}\right|^{2} \leq 4\left(\frac{\rho}{R}\right)^{n+3} \int_{\mathbb{B}_{R}}\left|u_{y}\right|^{2}+C\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1}\right)}^{2} R^{n-1+2 \sigma+\alpha} .
$$

Then, by Lemma 1.2.2, we have

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|u_{y}\right|^{2} & \leq C\left[\left(\frac{\rho}{R}\right)^{n-1+2 \sigma+\alpha} \int_{\mathbb{B}_{R}}\left|u_{y}\right|^{2}+\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1}\right)}^{2} \rho^{n-1+2 \sigma+\alpha}\right] \\
& \leq C\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1}\right)}^{2} \rho^{n+1+(\alpha-2+2 \sigma)} .
\end{aligned}
$$

Taking $1-\alpha / 4<\sigma<1$, we can guarantee that $\alpha-2+2 \sigma>\alpha / 2>0$, which implies that $T\left(y|y|^{-1} u_{y}\right)=0$ on $B_{1}$. Then, arguing as at the end of Case 1 , we conclude that $\left(-\Delta_{x}\right)^{1 / 2} u=0$ in $B_{1}$.

We finish this section with formal proof of Theorem J.

Proof of Theorem J. Parts (i), (ii), and (iii) are proved in Theorems 3.4.1, 3.4.2, and 3.4.4, respectively.

### 3.5 Almost minimizers for $s$-fractional obstacle problem

In this section we investigate the regularity of almost minimizers for the $s$-fractional obstacle problem with zero obstacle and give a proof of Theorem K. All results in this section are proved under the assumption $1 / 2 \leq s<1$, or $-1<a \leq 0$.

Theorem 3.5.1 (Almost Lipschitz regularity). Let u be an almost minimizer for s-fractional obstacle problem with zero obstacle in $B_{1}$, for $1 / 2 \leq s<1$. Then $u \in C^{0, \sigma}\left(B_{1}\right)$ for any $0<\sigma<1$ with

$$
\|u\|_{C^{0, \sigma}(K)} \leq C_{n, a, \omega, \sigma, K}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}
$$

for any $K \Subset B_{1}$.

Proof. Let $K \Subset B_{1}$ with $0 \in K$. Take $\delta=\delta_{n, a, \omega, \sigma, K}>0$ such that $\delta<\operatorname{dist}\left(K, \partial B_{1}\right)$ and $\omega(\delta) \leq \varepsilon$, where $\varepsilon=\varepsilon_{2, n+1+a, n-1+a+2 \sigma}$ as in Lemma 1.2.2. For $0<R<\delta$, let $v$ be the minimizer of

$$
\int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a}
$$

on $\mathfrak{K}_{0, u}\left(\mathbb{B}_{R},|y|^{a}\right)$. Then $v$ satisfies the variational inequality

$$
\int_{\mathbb{B}_{R}} \nabla v \nabla(w-v)|y|^{a} \geq 0
$$

for any $w \in \mathfrak{K}_{0, u}\left(\mathbb{B}_{R},|y|^{a}\right)$. Particularly, taking $w=u$, we have

$$
\int_{\mathbb{B}_{R}} \nabla v \nabla(u-v)|y|^{a} \geq 0
$$

As a consequence,

$$
\begin{aligned}
\int_{\mathbb{B}_{R}}|\nabla(u-v)|^{2}|y|^{a} & =\int_{\mathbb{B}_{R}}|\nabla u|^{2}|y|^{a}-\int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a}-2 \int_{\mathbb{B}_{R}} \nabla v \cdot \nabla(u-v)|y|^{a} \\
& \leq \omega(R) \int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a} .
\end{aligned}
$$

Next, we use (3.3.2) to derive a similar estimate for $u$. We have,

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}|\nabla u|^{2}|y|^{a} & \leq 2 \int_{\mathbb{B}_{\rho}}|\nabla v|^{2}|y|^{a}+2 \int_{\mathbb{B}_{\rho}}|\nabla(u-v)|^{2}|y|^{a} \\
& \leq 2\left(\frac{\rho}{R}\right)^{n+1+a} \int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a}+2 \omega(R) \int_{\mathbb{B}_{R}}|\nabla v|^{2}|y|^{a} \\
& \leq 2\left[\left(\frac{\rho}{R}\right)^{n+1+a}+\varepsilon\right] \int_{\mathbb{B}_{R}}|\nabla u|^{2}|y|^{a}
\end{aligned}
$$

Hence, by Lemma 1.2.2,

$$
\int_{\mathbb{B}_{\rho}}|\nabla u|^{2}|y|^{a} \leq C_{n, a, \sigma}\left(\frac{\rho}{R}\right)^{n-1+a+2 \sigma} \int_{\mathbb{B}_{R}}|\nabla u|^{2}|y|^{a} .
$$

As we have seen in Theorem 3.4.1, this implies

$$
\begin{equation*}
\int_{\mathbb{B}_{\rho}}|\nabla u|^{2}|y|^{a} \leq C_{n, a, \sigma, \delta}\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},\left.|y|\right|^{a}\right)}^{2} \rho^{n-1+a+2 \sigma} \tag{3.5.1}
\end{equation*}
$$

then

$$
\int_{\mathbb{B}_{\rho}}\left|u-\langle u\rangle_{\rho}\right|^{2}|y|^{a} \leq C_{n, a, \sigma, \delta}\|\nabla u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \rho^{n+1+a+2 \sigma}
$$

and ultimately

$$
\|u\|_{C^{0, \sigma}(K)} \leq C_{n, a, \omega, \sigma, K}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{\mid a}\right)} .
$$

Theorem 3.5.2 ( $C^{1, \beta}$ regularity). Let $u$ be an almost minimizer for the $s$-fractional obstacle problem with zero obstacle in $B_{1}, 1 / 2 \leq s<1$, and a gauge function $\omega(r)=r^{\alpha}$. Then $\nabla_{x} u \in C^{0, \beta}\left(B_{1}\right)$ for $\beta<\frac{\alpha s}{8(n+1+a+\alpha / 2)}$ and for any $K \Subset B_{1}$ there holds

$$
\left\|\nabla_{x} u\right\|_{C^{0, \beta}(K)} \leq C_{n, a, \alpha, \beta, K}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)} .
$$

Proof. Let $K$ be a thin ball centered at 0 such that $K \Subset \mathbb{B}_{1}$. Let $\varepsilon:=\frac{\alpha}{4(n+1+a+\alpha / 2)}$ and $\gamma:=1-\frac{s \varepsilon}{2(1-\varepsilon)}$. We fix $R_{0}=R_{0}(n, a, \alpha, K)>0$ small so that $R_{0}^{1-\varepsilon} \leq d / 2$, where $d:=$ $\operatorname{dist}\left(K, \partial B_{1}\right)$ and $R_{0}<\left(\frac{3}{16}\right)^{1 / \varepsilon}$. Then $\widetilde{K}:=\left\{x \in B_{1}: \operatorname{dist}(x, K) \leq R_{0}^{1-\varepsilon}\right\} \Subset B_{1}$. We claim that for $x_{0} \in K$ and $0<\rho<R<R_{0}$,

$$
\begin{gather*}
\int_{\mathbb{B}_{\rho}\left(x_{0}\right)}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{x_{0}, \rho}\right|^{2}|y|^{a} \leq C_{n, a}\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}\left(x_{0}\right)}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{x_{0}, R}\right|^{2}|y|^{a}  \tag{3.5.2}\\
+C_{n, a, \alpha, K}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} R^{n+1+a+s \varepsilon} .
\end{gather*}
$$

Note that once we have this bound, the proof will follow by the application of Lemma 1.2.2 and Theorem 3.A.1.

For simplicity we may assume $x_{0}=0$, and fix $0<R<R_{0}$. Let $\bar{R}:=R^{1-\varepsilon}$. Let $v$ be the minimizer of

$$
\int_{\mathbb{B}_{\bar{R}}}|\nabla v|^{2}|y|^{a}
$$

on $\mathfrak{K}_{0, u}\left(\mathbb{B}_{\bar{R}},|y|^{a}\right)$. Then by (3.3.2) and (3.5.1) with $\sigma=\gamma$, for $0<\rho \leq \bar{R}$

$$
\begin{align*}
\int_{\mathbb{B}_{\rho}}|\nabla v|^{2}|y|^{a} & \leq\left(\frac{\rho}{\bar{R}}\right)^{n+1+a} \int_{\mathbb{B}_{\bar{R}}}|\nabla v|^{2}|y|^{a} \leq\left(\frac{\rho}{\bar{R}}\right)^{n+1+a} \int_{\mathbb{B}_{\bar{R}}}|\nabla u|^{2}|y|^{a} \\
& \leq C_{n, a, \alpha, K}\left(\frac{\rho}{\bar{R}}\right)^{n+1+a}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \bar{R}^{n-1+a+2 \gamma}  \tag{3.5.3}\\
& \leq C_{n, a, \alpha, K}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \rho^{n-1+a+2 \gamma} .
\end{align*}
$$

This gives

$$
\begin{equation*}
f_{\mathbb{B}_{\rho}}\left|v-v_{\rho}\right|^{2}|y|^{a} \leq C_{1}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \rho^{2 \gamma}, \quad C_{1}=C_{n, a, \alpha, K} . \tag{3.5.4}
\end{equation*}
$$

Since this estimate holds for any $0<\rho<\bar{R}$, the standard dyadic argument gives

$$
\begin{equation*}
\left|v(0)-\langle v\rangle_{\bar{R}}\right| \leq C_{2}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},\left.|y|\right|^{a}\right)} \bar{R}^{\gamma}, \quad C_{2}=C_{n, a, \alpha, K} . \tag{3.5.5}
\end{equation*}
$$

Moreover, using (3.3.2) and (3.5.1) again, we have for any $x_{1} \in B_{\bar{R} / 2}, 0<\rho<\bar{R} / 2$,

$$
\begin{align*}
\int_{\mathbb{B}_{\rho}\left(x_{1}\right)}|\nabla v|^{2}|y|^{a} & \leq\left(\frac{2 \rho}{\bar{R}}\right)^{n+1+a} \int_{\mathbb{B}_{\bar{R} / 2}\left(x_{1}\right)}|\nabla v|^{2}|y|^{a} \leq\left(\frac{2 \rho}{\bar{R}}\right)^{n+1+a} \int_{\mathbb{B}_{\bar{R}}}|\nabla u|^{2}|y|^{a}  \tag{3.5.6}\\
& \leq C_{n, a, \alpha, K}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \rho^{n-1+a+2 \gamma},
\end{align*}
$$

which implies

$$
\begin{equation*}
[v]_{C^{0, \gamma}\left(\overline{B_{R / 2}}\right)} \leq C_{3}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}, \quad C_{3}=C_{n, a, \alpha, K} . \tag{3.5.7}
\end{equation*}
$$

Now we define

$$
C_{4}:=C_{1}+C_{2}^{2}+C_{3}^{2} .
$$

Our analysis then distinguishes the following two cases

$$
\left\langle v^{2}\right\rangle_{\bar{R}} \leq 6 C_{4}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \bar{R}^{2 \gamma} \quad \text { or } \quad\left\langle v^{2}\right\rangle_{\bar{R}}>6 C_{4}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \bar{R}^{2 \gamma} .
$$

Case 1. Suppose first that

$$
\left\langle v^{2}\right\rangle_{\bar{R}} \leq 6 C_{4}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \bar{R}^{2 \gamma}
$$

Note that $R_{0}<\left(\frac{3}{16}\right)^{1 / \varepsilon}$ implies $R<\frac{3}{16} \bar{R}$. Then, using Corollary 3.3.6, we see that for $0<\rho<R$,

$$
\begin{aligned}
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{\rho}\right|^{2}|y|^{a} \leq & 3 \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a}+6 \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} d x \\
\leq & C_{n, a}\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{R}\right|^{2}|y|^{a} \\
& +C_{n, a} \frac{R^{n+2}}{\bar{R}^{2+2 s}}\left\langle v^{2}\right\rangle_{\bar{R}}+6 \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} \\
\leq & C\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{R}\right|^{2}|y|^{a} \\
& +C \frac{R^{n+2}}{\bar{R}^{2+2 s}}\left\langle v^{2}\right\rangle_{\bar{R}}+C \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} .
\end{aligned}
$$

Note that for $\sigma:=1-\alpha / 4$

$$
\begin{aligned}
\int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} & \leq \int_{\mathbb{B}_{\bar{R}}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} \leq \bar{R}^{\alpha} \int_{\mathbb{B}_{\bar{R}}}|\nabla v|^{2}|y|^{a} \\
& \leq \bar{R}^{\alpha} \int_{\mathbb{B}_{\bar{R}}}|\nabla u|^{2}|y|^{a} \leq C_{n, a, \alpha, K} \bar{R}^{\alpha}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \bar{R}^{n-1+a+2 \sigma} \\
& =C\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} R^{n+1+a+\alpha / 4} .
\end{aligned}
$$

Moreover by the assumption

$$
C \frac{R^{n+2}}{\bar{R}^{2+2 s}}\left\langle v^{2}\right\rangle_{\bar{R}} \leq C_{n, a, \alpha, K}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} R^{n+2} \bar{R}^{2 \gamma-2-2 s}=C\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} R^{n+1+a+s \varepsilon} .
$$

Hence, we obtain (3.5.2) in this case.
Case 2. Now we assume

$$
\left\langle v^{2}\right\rangle_{\bar{R}}>6 C_{4}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \bar{R}^{2 \gamma}
$$

Then, by (3.5.4) and (3.5.5) we obtain

$$
f_{\mathbb{B}_{\bar{R}}}|v-v(0)|^{2}|y|^{a} \leq 2 f_{\mathbb{B}_{\bar{R}}}\left|v-v_{\bar{R}}\right|^{2}|y|^{a}+2 f_{\mathbb{B}_{\bar{R}}}\left|v_{\bar{R}}-v(0)\right|^{2}|y|^{a} \leq 2 C_{4}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \bar{R}^{2 \gamma} .
$$

Combining the latter bound and the assumption,

$$
\begin{aligned}
v(0)^{2} & =f_{\mathbb{B}_{\bar{R}}}|v(0)|^{2}|y|^{a} \geq \frac{1}{2} f_{\mathbb{B}_{\bar{R}}}|v(X)|^{2}|y|^{a}-f_{\mathbb{B}_{\bar{R}}}|v(X)-v(0)|^{2}|y|^{a} \\
& \geq C_{4}\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} \bar{R}^{2 \gamma} .
\end{aligned}
$$

Since $C_{4} \geq C_{3}^{2}$, we have $v>0$ on $B_{\bar{R} / 2}$ by (3.5.7). Thus, $L_{a} v=0$ in $\mathbb{B}_{\bar{R} / 2}$, and by Lemma 3.3.2 we have for $0<\rho<R$

$$
\int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a} \leq\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{R}\right|^{2}|y|^{a} .
$$

Thus,

$$
\begin{aligned}
& \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{\rho}\right|^{2}|y|^{a} \\
& \leq 3 \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{\rho}\right|^{2}|y|^{a}+6 \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} \\
& \leq 3\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} v-\left\langle\nabla_{x} v\right\rangle_{R}\right|^{2}|y|^{a}+6 \int_{\mathbb{B}_{\rho}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} \\
& \leq C\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{R}\right|^{2}|y|^{a}+C \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\nabla_{x} v\right|^{2}|y|^{a} \\
& \leq C\left(\frac{\rho}{R}\right)^{n+a+3} \int_{\mathbb{B}_{R}}\left|\nabla_{x} u-\left\langle\nabla_{x} u\right\rangle_{R}\right|^{2}|y|^{a}+C\|u\|_{W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right)}^{2} R^{n+1+a+\alpha / 4} .
\end{aligned}
$$

This implies (3.5.2) and completes the proof.
Proof of Theorem K. Parts (i) and (ii) are contained in Theorems 3.5.1 and 3.5.2, respectively.

## 3.A Morrey-Campanato-type Space

Theorem 3.A.1. Let $u \in L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)$ and $M$ be such that $\|u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)} \leq M$ and for some $\sigma \in(0,1)$

$$
\int_{\mathbb{B}_{r}(x)}\left|u-\langle u\rangle_{x, r}\right|^{2}|y|^{a} \leq M^{2} r^{n+1+a+2 \sigma}, \quad\langle u\rangle_{x, r}=\frac{1}{\omega_{n+1+a} r^{n+1+a}} \int_{\mathbb{B}_{r}(x)} u|y|^{a}
$$

for any ball $\mathbb{B}_{r}(x)$ centered at $x=(x, 0) \in B_{1 / 2}$ and radius $0<r<r_{0} \leq 1 / 2$. Then for any $x \in B_{1 / 2}$ there exists the limit of averages

$$
T u(x):=\lim _{r \rightarrow 0}\langle u\rangle_{x, r},
$$

which will also satisfy

$$
\int_{\mathbb{B}_{r}(x)}|u-T u(x)|^{2}|y|^{a} \leq C_{n, a, \sigma} M^{2} r^{n+1+a+2 \sigma}
$$

Moreover, $T u \in C^{0, \sigma}\left(B_{1 / 2}\right)$ with

$$
\|T u\|_{C^{0, \sigma}\left(B_{1 / 2}\right)} \leq C_{n, a, \sigma, r_{0}} M
$$

Remark 3.A.2. Note, we can redefine $u(x, 0)=T u(x)$ for any $x \in B_{1 / 2}$, making $(x, 0)$ a Lebesgue point for $u$.

Proof. Let $x, z \in B_{1 / 2}$ and $0<\rho<r<r_{0}$ be such that $\mathbb{B}_{\rho}(x) \subset \mathbb{B}_{r}(z)$. Then

$$
\begin{aligned}
\left|\langle u\rangle_{x, \rho}-\langle u\rangle_{z, r}\right| & \leq f_{\mathbb{B}_{\rho}(x)}\left|u-\langle u\rangle_{z, r}\right||y|^{a} \leq\left(\frac{r}{\rho}\right)^{n+1+a} f_{\mathbb{B}_{r}(z)}\left|u-\langle u\rangle_{z, r}\right||y|^{a} \\
& \leq\left(\frac{r}{\rho}\right)^{n+1+a}\left(f_{\mathbb{B}_{r}(z)}\left|u-\langle u\rangle_{z, r}\right|^{2}|y|^{a}\right)^{1 / 2}\left(f_{\mathbb{B}_{r}(z)}|y|^{a}\right)^{1 / 2} \\
& \leq C_{n, a}\left(\frac{r}{\rho}\right)^{n+1+a} M r^{\sigma} .
\end{aligned}
$$

Now, taking $x=z$ and using a dyadic argument, we can conclude that

$$
\left|\langle u\rangle_{x, \rho}-\langle u\rangle_{x, r}\right| \leq C_{n, a, \sigma} M r^{\sigma}, \quad \text { for any } 0<s=\rho<r<r_{0} .
$$

Indeed, let $k=0,1,2, \ldots$ be such that $r / 2^{k+1} \leq \rho<r / 2^{k}$. Then

$$
\left|\langle u\rangle_{x, \rho}-\langle u\rangle_{x, r}\right| \leq \sum_{\mathrm{j}=1}^{k}\left|\langle u\rangle_{x, r / 2^{\mathrm{j}-1}}-\langle u\rangle_{x, r / 2^{\mathrm{j}}}\right|+\left|\langle u\rangle_{x, r / 2^{k}}-\langle u\rangle_{x, \rho}\right|
$$

$$
\leq C_{n, a} M \sum_{\mathrm{j}=1}^{k+1}\left(r / 2^{\mathrm{j}-1}\right)^{\sigma} \leq C_{n, a, \sigma} M r^{\sigma} .
$$

This implies that the limit

$$
T u(x)=\lim _{r \rightarrow 0}\langle u\rangle_{x, r}
$$

exists and

$$
\left|T u(x)-\langle u\rangle_{x, r}\right| \leq C_{n, a, \sigma} M r^{\sigma} .
$$

Hence, we also have the Hölder integral bound

$$
\int_{\mathbb{B}_{r}(x)}|u-T u(x)|^{2}|y|^{a} \leq C_{n, a, \sigma} M^{2} r^{n+1+a+2 \sigma}
$$

Besides, we have

$$
|T u(x)| \leq\langle u\rangle_{x, r_{0}}+C_{n, a, \sigma} M r_{0}^{\sigma} \leq C_{n, a, \sigma, r_{0}} M .
$$

It remains to estimate the Hölder seminorm of $T u$ on $B_{1 / 2}$. Let $x, z \in B_{1 / 2}$ and consider two cases.

Case 1. If $|x-z|<r_{0} / 4$, let $r=2|x-z|$. Then note that $\mathbb{B}_{r / 2}(x) \subset \mathbb{B}_{r}(z)$ and therefore we can write

$$
\begin{aligned}
|T u(x)-T u(z)| & \leq\left|T u(x)-\langle u\rangle_{x, r / 2}\right|+\left|T u(z)-\langle u\rangle_{z, r}\right|+\left|\langle u\rangle_{x, r / 2}-\langle u\rangle_{z, r}\right| \\
& \leq C_{n, a, \sigma} M r^{\sigma}=C_{n, a, \sigma} M|x-z|^{\sigma}
\end{aligned}
$$

Case 2. If $|x-z| \geq r_{0} / 4$, then

$$
|T u(x)-T u(z)| \leq|T u(x)|+|T u(z)| \leq C_{n, a, \sigma, r_{0}} M \leq C_{n, a, \sigma, r_{0}} M|x-z|^{\sigma} .
$$

Thus, we conclude

$$
\|T u\|_{C^{0, \sigma}\left(B_{1 / 2}\right)} \leq C_{n, a, \sigma, r_{0}} M
$$

## 3.B Polynomial expansion for Caffarelli-Silvestre extension

Some of the results in Section 3.3 rely on polynomial expansion theorem for $L_{a}$-harmonic functions given below.

Theorem 3.B.1. Let $u \in W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right),-1<a<1$, be a weak solution of the equation $L_{a} u=0$ in $\mathbb{B}_{1}$, even in $y$. Then we have the following polynomial expansion:

$$
u(x, y)=\sum_{k=0}^{\infty} p_{k}(x, y)
$$

locally uniformly in $\mathbb{B}_{1}$, where $p_{k}(x, y)$ are $L_{a}$-harmonic polynomials, homogeneous of degree $k$ and even in $y$. Moreover, the polynomials $p_{k}$ above are orthogonal in $L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)$, i.e.,

$$
\int_{\partial \mathbb{B}_{1}} p_{k} p_{m}|y|^{a}=0, \quad k \neq m .
$$

In particular, $u$ is real analytic in $\mathbb{B}_{1}$.

This theorem has the following immediate corollaries, which are of independent interest and are likely known in the literature. We state them here for reader's convenience and for possible future reference.

Corollary 3.B.2. Let $u \in W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right),-1<a<1$, be a weak solution of the equation $L_{a} u=0$ in $\mathbb{B}_{1}$. Then, we have a representation

$$
u(x, y)=\varphi(x, y)+y|y|^{-a} \psi(x, y), \quad(x, y) \in \mathbb{B}_{1}
$$

where $\varphi(x, y)$ and $\psi(x, y)$ are real analytic functions, even in $y$.
Corollary 3.B.3. Let $u \in \mathcal{L}_{s}\left(\mathbb{R}^{n}\right)$ satisfies $(-\Delta)^{s} u=0$ in the unit ball $B_{1} \subset \mathbb{R}^{n}$. Then $u$ is real analytic in $B_{1}$.

Corollary 3.B.4. Let $u \in W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right),-1<a<1$, be a weak solution of the equation $L_{a} u=0$ in $\mathbb{B}_{1}$, even in $y$. If $u(\cdot, 0) \equiv 0$ in $B_{1}$, then $u \equiv 0$ in $\mathbb{B}_{1}$.

The proof of Theorem 3.B. 1 and subsequently those of Corollaries 3.B.2, 3.B.3, and 3.B.4 are based on the following lemmas. We follow the approach of [63] for harmonic functions.

Let $\mathcal{P}_{m}^{*}=\{p: p(x, y)$ polynomial of degree $\leq m$, even in $y\}$.
Lemma 3.B.5. Let $p \in \mathcal{P}_{m}^{*}$. Then there exists $\tilde{p} \in \mathcal{P}_{m}^{*}$ such that

$$
L_{a} \tilde{p}=0 \quad \text { in } \mathbb{B}_{1}, \quad \tilde{p}=p \quad \text { on } \partial \mathbb{B}_{1} .
$$

In other words, the solution of the Dirichlet problem for $L_{a}$ in $\mathbb{B}_{1}$ with boundary values in $\mathcal{P}_{m}^{*}$ on $\partial \mathbb{B}_{1}$ is itself in $\mathcal{P}_{m}^{*}$.

Proof. For $m=0,1$, we simply have $\tilde{p}=p$. For $m \geq 2$, we proceed as follows.
For $q \in \mathcal{P}_{m-2}^{*}$ define $T q \in \mathcal{P}_{m-2}^{*}$ by

$$
(T q)(x, y)=|y|^{-a} L_{a}\left(\left(1-x^{2}-y^{2}\right) q(x, y)\right) .
$$

(It is straightforward to verify that $T q$ is indeed in $\mathcal{P}_{m-2}^{*}$ ). We now claim that the mapping $T: \mathcal{P}_{m-2}^{*} \rightarrow \mathcal{P}_{m-2}^{*}$ is bijective. Since $T$ is clearly linear and $\mathcal{P}_{m-2}^{*}$ is finite dimensional it is equivalent to showing that $T$ is injective. To this end, suppose that $T q=0$ for some $q \in \mathcal{P}_{m-2}^{*}$. This means that $Q(x, y)=\left(1-x^{2}-y^{2}\right) q(x, y)$ is $L_{a}$-harmonic in $\mathbb{B}_{1}$ :

$$
L_{a} Q=0 \quad \text { in } \mathbb{B}_{1} .
$$

On the other hand $Q=0$ on $\partial \mathbb{B}_{1}$ and therefore, by the maximum principle $Q=0$ in $\mathbb{B}_{1}$. But this implies that $q=0$ in $\mathbb{B}_{1}$, or that $q \equiv 0$. Hence, the mapping $T$ is injective, and consequently bijective. It is now easy to see that

$$
\tilde{p}=p-\left(1-x^{2}-y^{2}\right) T^{-1}\left(|y|^{-a} L_{a}(p)\right) \in \mathcal{P}_{m}^{*}
$$

satisfies the required properties.
Lemma 3.B.6. Polynomials, even in $y$, are dense in the subspace of functions in $L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)$, even in $y$.

Proof. Polynomials, even in $y$ are dense in the space of continuous functions in $C\left(\partial \mathbb{B}_{1}\right)$, even in $y$, with the uniform norm. The claim now follows from the observation that the embedding $C\left(\partial \mathbb{B}_{1}\right) \hookrightarrow L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)$ is continuous:

$$
\|v\|_{L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)} \leq\|v\|_{L^{\infty}\left(\partial \mathbb{B}_{1}\right)}\left(\int_{\partial \mathbb{B}_{1}}|y|^{a}\right)^{1 / 2} \leq C\|v\|_{L^{\infty}\left(\partial \mathbb{B}_{1}\right)} .
$$

Lemma 3.B.7. The subspace of functions in $L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)$, even in $y$, has an orthonormal basis $\left\{p_{k}\right\}_{k=0}^{\infty}$ consisting of homogeneous $L_{a}$-harmonic polynomials $p_{k}$, even in $y$.

Proof. If $p$ is a polynomial, even in $y$, then restricted to $\partial \mathbb{B}_{1}$ it can be replaced with an $L_{a}$-harmonic polynomial $\tilde{p}$. On the other hand, if we decompose

$$
\tilde{p}=\sum_{i=0}^{m} q_{i}
$$

where $q_{i}$ is a homogeneous polynomial of order i , even in $y$, then

$$
|y|^{-a} L_{a} \tilde{p}=\sum_{i=2}^{m}|y|^{-a} L_{a} q_{i}
$$

where $|y|^{-a} L_{a} q_{i}$ is a homogeneous polynomial of order $i-2, i=2, \ldots, m$. Hence, $L_{a} \tilde{p}=0$ iff $L_{a} q_{i}=0$, for all $i=0, \ldots, m$ (for $i=0,1$ this holds automatically).

We further note that if $q_{i}$ and $q_{j}$ are two homogeneous $L_{a}$-harmonic polynomials of degrees $i \neq j$, then they are orthogonal in $L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)$. Indeed,

$$
0=\int_{\mathbb{B}_{1}} q_{i} \operatorname{div}\left(|y|^{a} \nabla q_{j}\right)-\operatorname{div}\left(|y|^{a} \nabla q_{i}\right) q_{j}=\int_{\partial \mathbb{B}_{1}}\left(q_{i} \partial_{\nu} q_{j}-q_{j} \partial_{\nu} q_{i}\right)|y|^{a}=(j-i) \int_{\partial \mathbb{B}_{1}} q_{i} q_{j}|y|^{a} .
$$

Using this and following the standard orthogonalization process, we can construct a basis consisting of homogeneous $L_{a}$-harmonic polynomials.

Lemma 3.B.8. Let $u \in W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right) \cap C\left(\overline{\mathbb{B}_{1}}\right)$ is a weak solution of $L_{a} u=0$ in $\mathbb{B}_{1}$. Then

$$
\|u\|_{L^{\infty}(K)} \leq C_{n, a, K}\|u\|_{L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)} .
$$

for any $K \Subset \mathbb{B}_{1}$.

Proof. First, we note that by [64]

$$
\|u\|_{L^{\infty}(K)} \leq C_{n, a, K}\|u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)} .
$$

So we just need to show that

$$
\|u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)} \leq C_{n, a}\|u\|_{L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)} .
$$

This follows from the fact that $u^{2}$ is a subsolution: $L_{a}\left(u^{2}\right) \geq 0$ in $\mathbb{B}_{1}$ and therefore the weighted spherical averages

$$
r \mapsto \frac{1}{\omega_{n, a} r^{n+a}} \int_{\partial \mathbb{B}_{r}} u^{2}|y|^{a}, \quad 0<r<1
$$

are increasing. Integrating, we easily obtain that

$$
\|u\|_{L^{2}\left(\mathbb{B}_{1},|y|^{a}\right)} \leq C_{n, a}\|u\|_{L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)} .
$$

We are now ready to prove Theorem 3.B.1.
Proof of Theorem 3.B.1. Without loss of generality we may assume $u \in W^{1,2}\left(\mathbb{B}_{1},|y|^{a}\right) \cap$ $C\left(\overline{\mathbb{B}_{1}}\right)$, otherwise we can consider a slightly smaller ball. Now, using the orthonormal basis $\left\{p_{k}\right\}_{k=0}^{\infty}$ from Lemma 3.B. 7 we represent

$$
u=\sum_{k=0}^{\infty} a_{k} p_{k} \quad \text { in } L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right) .
$$

We then claim that

$$
u(x, y)=\sum_{k=0}^{\infty} a_{k} p_{k}(x, y) \quad \text { uniformly on any } K \Subset \mathbb{B}_{1} .
$$

Indeed, if $u_{m}(x, y)=\sum_{k=0}^{m} a_{k} p_{k}(x, y)$, then $\left\|u-u_{m}\right\|_{L^{2}\left(\partial \mathbb{B}_{1},|y|^{2}\right)} \rightarrow 0$ as $m \rightarrow \infty$ and therefore by Lemma 3.B. 8

$$
\left\|u-u_{m}\right\|_{L^{\infty}(K)} \leq C_{n, a, K}\left\|u-u_{m}\right\|_{L^{2}\left(\partial \mathbb{B}_{1},|y|^{a}\right)} \rightarrow 0 .
$$

We now give the proofs of the corollaries.

Proof of Corollary 3.B.2. Write $u(x, y)$ in the form

$$
u(x, y)=u_{\mathrm{even}}(x, y)+u_{\mathrm{odd}}(x, y)
$$

where $u_{\text {even }}$ and $u_{\text {odd }}$ are even and odd in $y$, respectively. Clearly, both functions are $L_{a^{-}}$ harmonic. Moreover, by Theorem 3.B.1, $u_{\text {even }}$ is real analytic and we take $\varphi=u_{\text {even }}$. On the other hand, consider

$$
v(x, y)=|y|^{a} \partial_{y} u_{\mathrm{odd}}(x, y)
$$

Then, $v$ is $L_{-a}$-harmonic in $\mathbb{B}_{1}$ and again by Theorem 3.B.1, $v$ is real analytic. We can now represent

$$
u_{\mathrm{odd}}(x, y)=y|y|^{-a} \psi(x, y), \quad \psi(x, y)=y^{-1}|y|^{a} \int_{0}^{y}|s|^{-a} v(x, s) d s
$$

It is not hard to see that $\psi(x, y)$ is real analytic, which completes our proof.
Proof of Corollary 3.B.3. The proof follows immediately from Theorem 3.B. 1 by considering the Caffarelli-Silvestre extension

$$
u(x, y)=u * P(\cdot, y)=\int_{\mathbb{R}^{n}} P(x-z, y) u(z) d z, \quad(x, y) \in \mathbb{R}^{n} \times \mathbb{R}_{+}
$$

where $P(x, y)=C_{n, a} \frac{y^{1-a}}{\left(|x|^{2}+y^{2}\right)^{(n+1-a) / 2}}$ is the Poisson kernel for $L_{a}$, and noting that its extension to $\mathbb{R}^{n+1}$ by even symmetry in $y$ (still denoted $u$ ) satisfies $L_{a} u=0$ in $\mathbb{B}_{1}$.

Proof of Corollary 3.B.4. Represent $u(x, y)$ as a locally uniformly convergent in $\mathbb{B}_{1}$ series

$$
u(x, y)=\sum_{k=0}^{\infty} q_{k}(x, y)
$$

where $q_{k}(x, y)$ is a homogeneous of degree $k L_{a}$-harmonic polynomial, even in $y$. We have

$$
u(x, 0)=\sum_{k=0}^{\infty} q_{k}(x, 0) \equiv 0
$$

from which we conclude that $q_{k}(x, 0) \equiv 0$. We now want to show that $q_{k} \equiv 0$. To this end represent

$$
q_{k}(x)=\sum_{j=0}^{[k / 2]} p_{k-2 j}(x) y^{2 j}
$$

where $p_{k-2 j}(x)$ is a homogeneous polynomial of order $k-2 j$ in $x$. Clearly $p_{k}(x) \equiv 0$. Taking partial derivatives $\partial_{x}^{\alpha} q_{k}(x)$ of order $|\alpha|=k-2$, we see that

$$
\partial_{x}^{\alpha} q_{k}(x)=c_{\alpha} y^{2}, \quad c_{\alpha}=\partial_{x}^{\alpha} p_{k-2}
$$

is $L_{a}$-harmonic, which can happen only when $c_{\alpha}=0$. Hence $D_{x}^{k-2} p_{k-2}(x) \equiv 0$ and therefore $p_{k-2} \equiv 0$. Then taking consequently derivatives of orders $k-2 j, j=2, \ldots$, we conclude that $p_{k-2 \mathrm{j}}(x) \equiv 0$ for all $j=0, \ldots,[k / 2]$ and hence $q_{k}(x, y) \equiv 0$.

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[^0]:    ${ }^{1} \uparrow$ This can be seen on the explicit example $u(x)=\operatorname{Re}\left(x_{1}+i\left|x_{n}\right|\right)^{3 / 2}$, which is a solution of the obstacle problem with $\psi=0$ on $\mathcal{M}=\left\{x_{n}=0\right\}$.

[^1]:    ${ }^{3} \uparrow$ We use the superscript $A$ to distinguish this rescaling from the other rescalings, namely, homogeneous and almost homogeneous rescalings that we consider later.

